

**The 1st Setouchi Neutrino Workshop**

**Teshima, Japan**

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# **Exploring the Fate of Massive Stars with Diffuse Supernova Neutrino Background**

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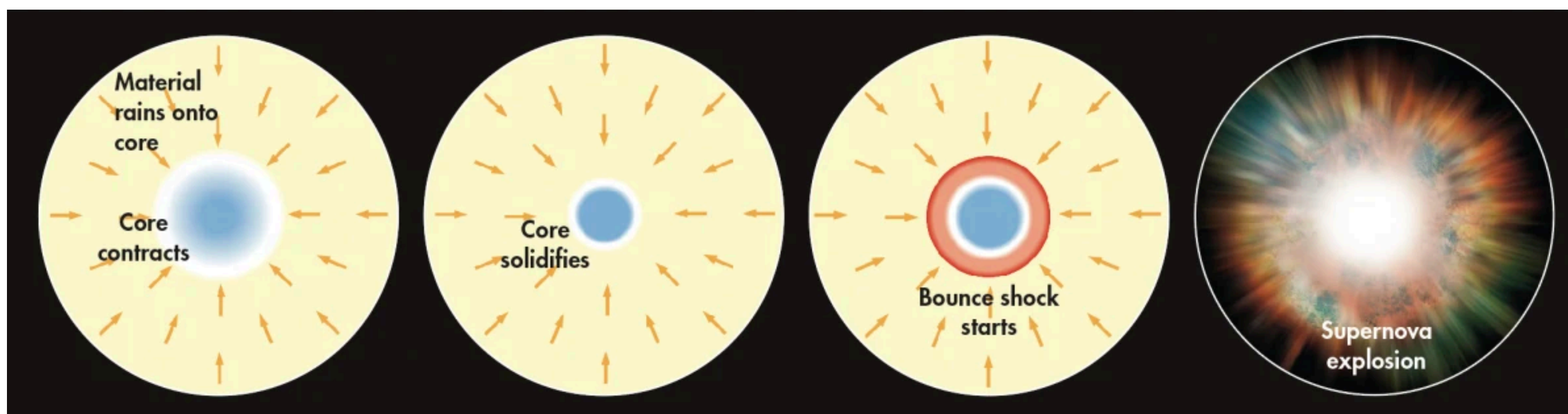
<sup>3</sup>National Astronomical Observatory of Japan



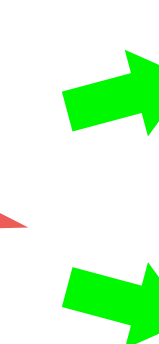
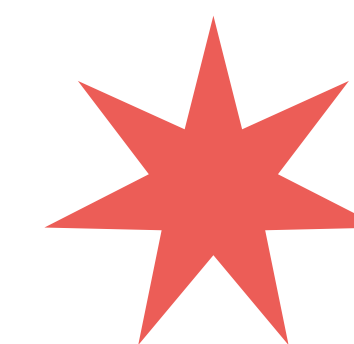
# Introduction

# Stellar Core Collapse

- Massive stars ( $>8M_{\text{Sun}}$ ) cause **gravitational collapse of the core** at the end of their life.
  - One of the most dynamic phenomena in the Universe (**energy  $\sim 10^{53}$  erg**).
  - Play an important role in cosmic chemical evolution.
- Multiple patterns are possible after core collapse.
  - **In a successful case, a neutron star is produced afterwards, and the explosion is optically observable.**
  - Otherwise, **a black hole is produced and not possible to be optically observed (failed SN).**
  - **Neutrinos** are emitted from both cases and play an important role in the mechanism.



core collapse



neutron star

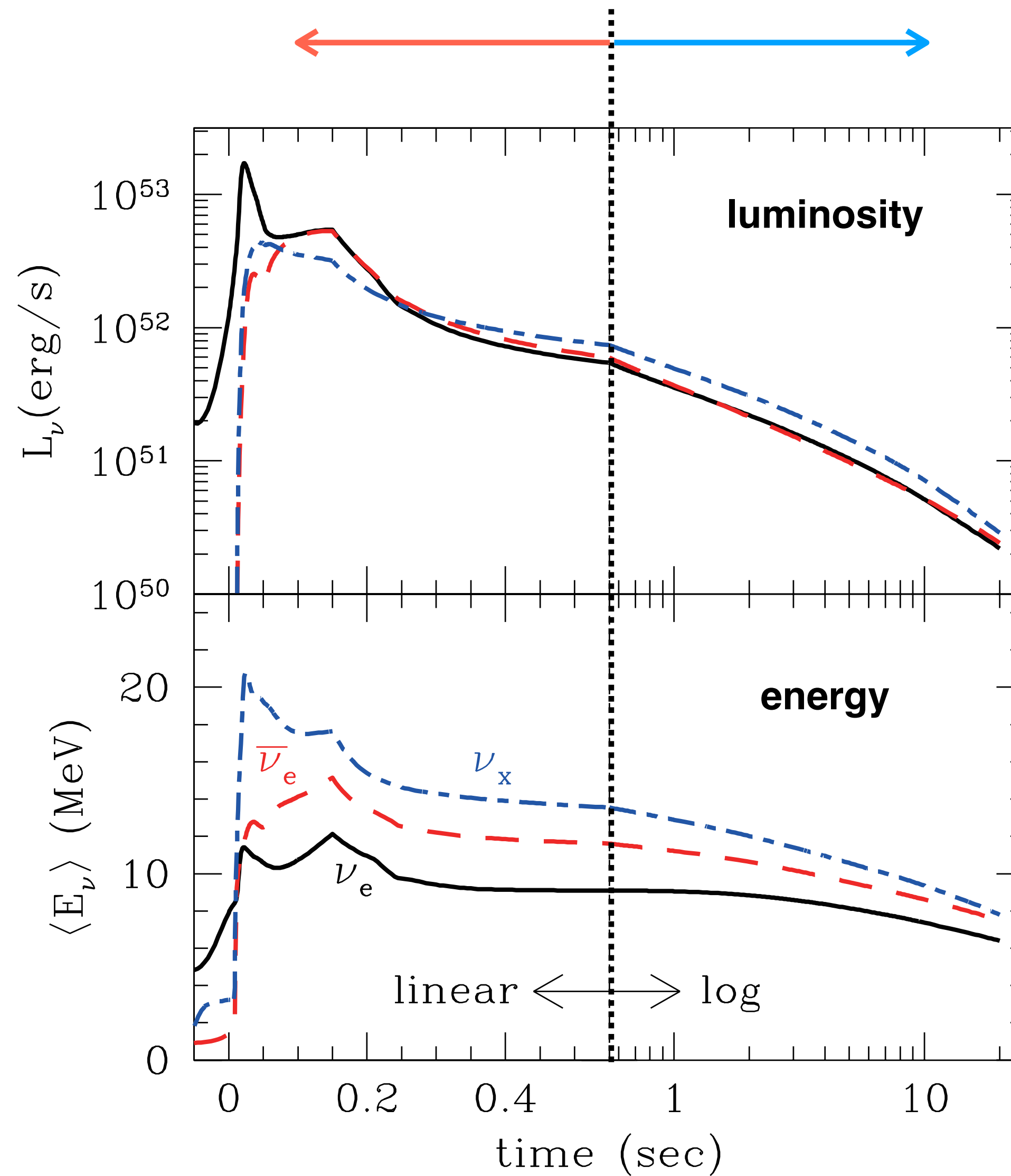
black hole (failed SN)

## accretion phase

- Both **NS** and **BH** cases
- Longer accretion for BH case
- Mainly  $\nu_e$  and anti- $\nu_e$
- Higher energy

## cooling phase

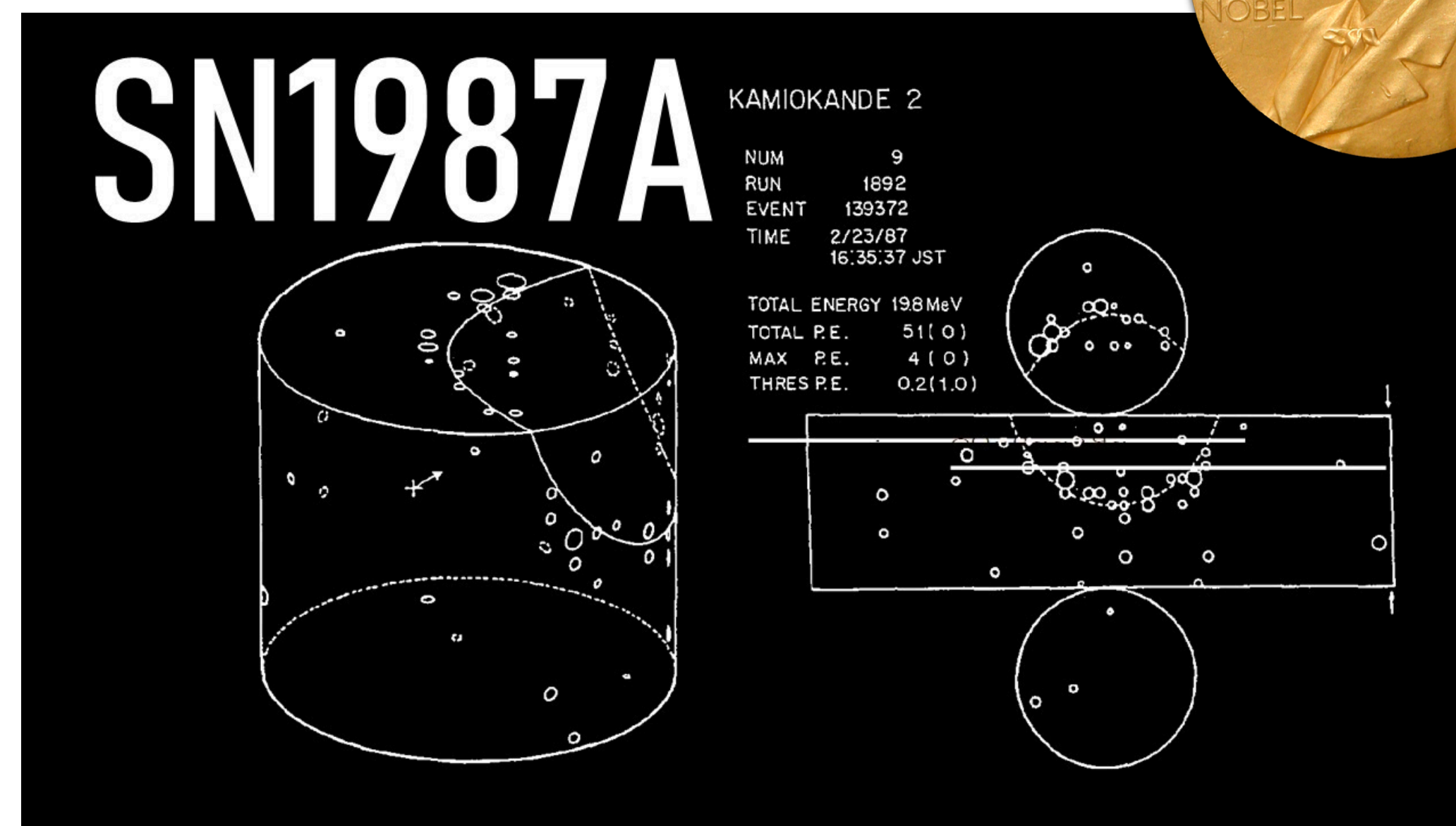
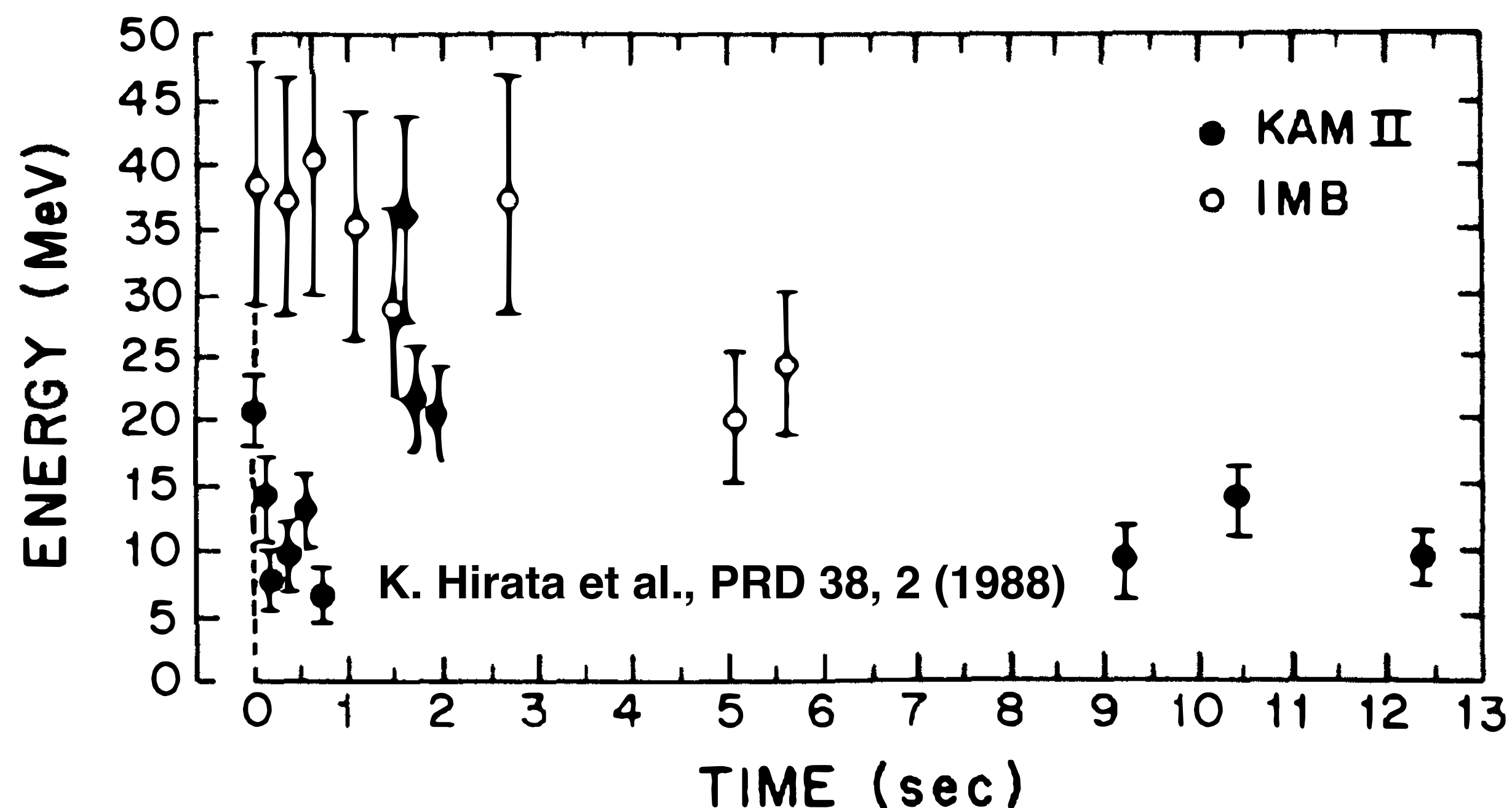
- Only **NS** case
- Large integrated flux
- All flavors
- Lower energy



# Observation of Supernova Neutrinos

So far **only one observation of supernova neutrinos (SN1987A)**.

- Low supernova rate (~a few times/century/galaxy)
- Small neutrino cross section
- ~10 neutrino events expected, for a *Mpc* far SN, even at a *Mton*-size detector.



# Diffuse Supernova Neutrino Background

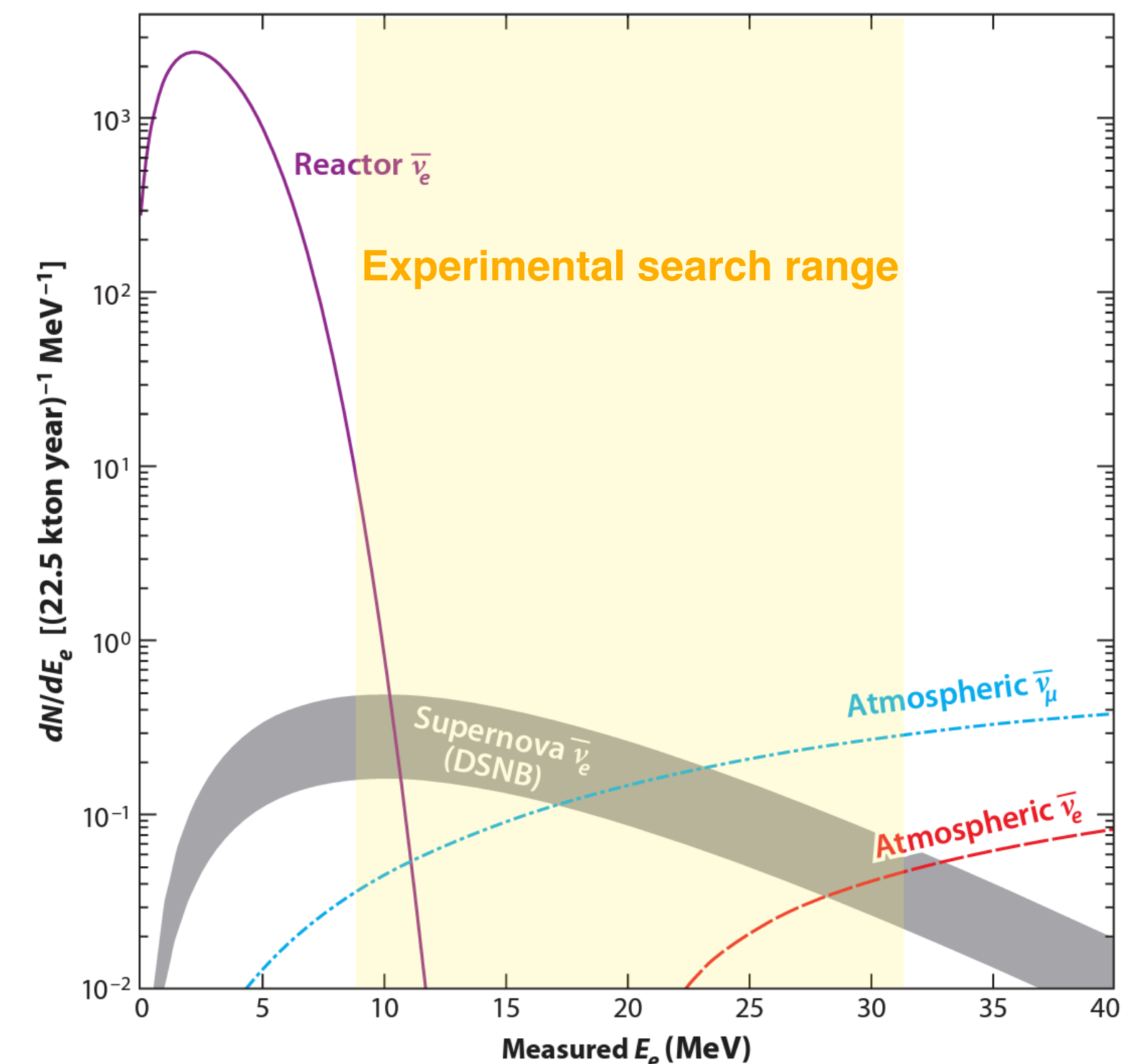
The accumulated flux of neutrinos from all past core collapses over the cosmic history

= **Diffuse Supernova Neutrino Background (DSNB)**

- Many factors affecting DSNB (SFR, nuclear EOS, BH formation, neutrino oscillation, etc).
- Experimental detection is challenging because of small flux & huge backgrounds.

$$\Phi = \int [\nu \text{ emission}] \otimes [\text{Star formation}] \otimes [\text{Universe expansion}]$$

J. F. Beacom, ARNPS 60, 439 (2010)

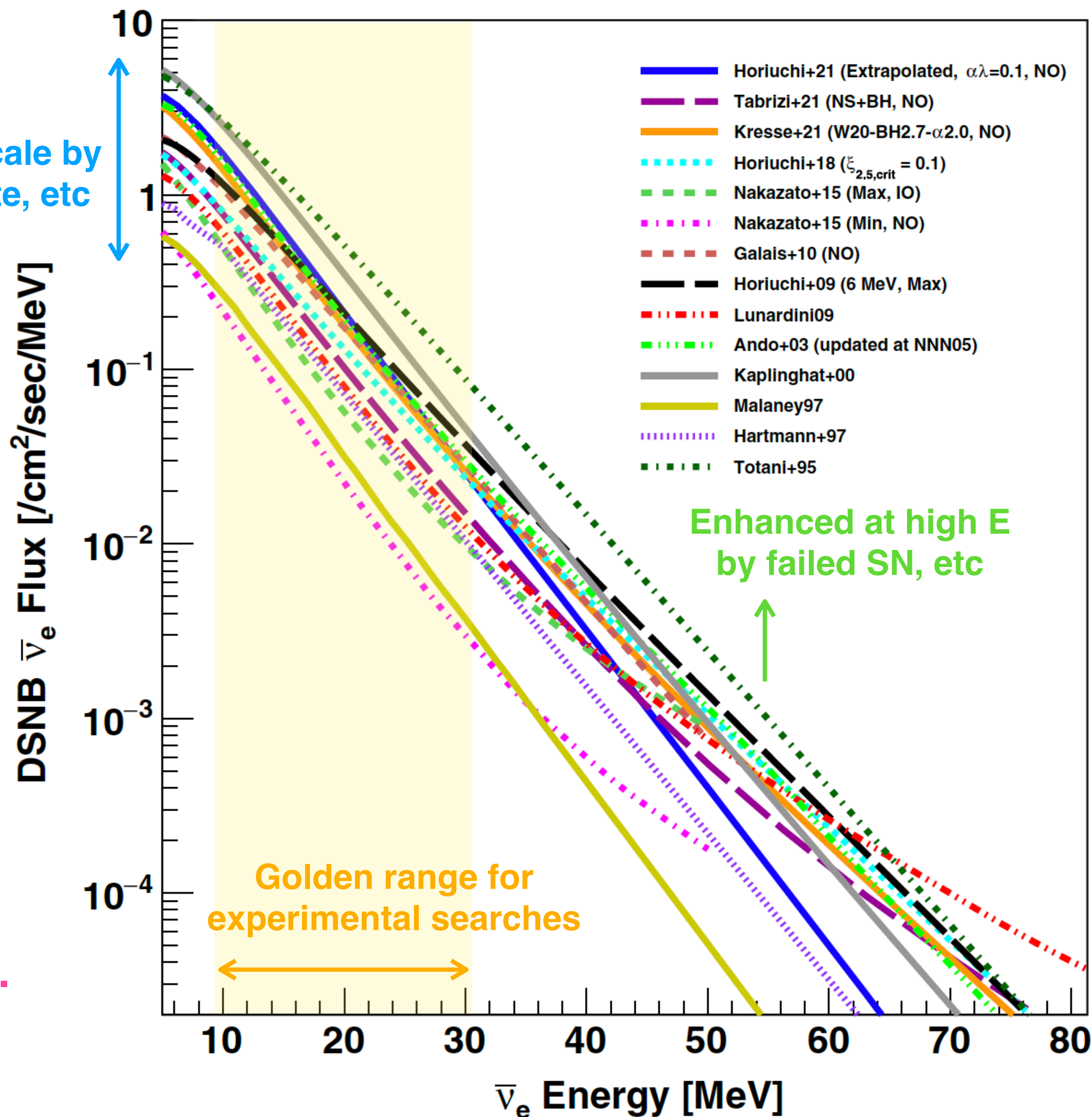


# DSNB Flux Predictions

Most theoretical predictions exist within  
~1 order of magnitude at 10~30 MeV.

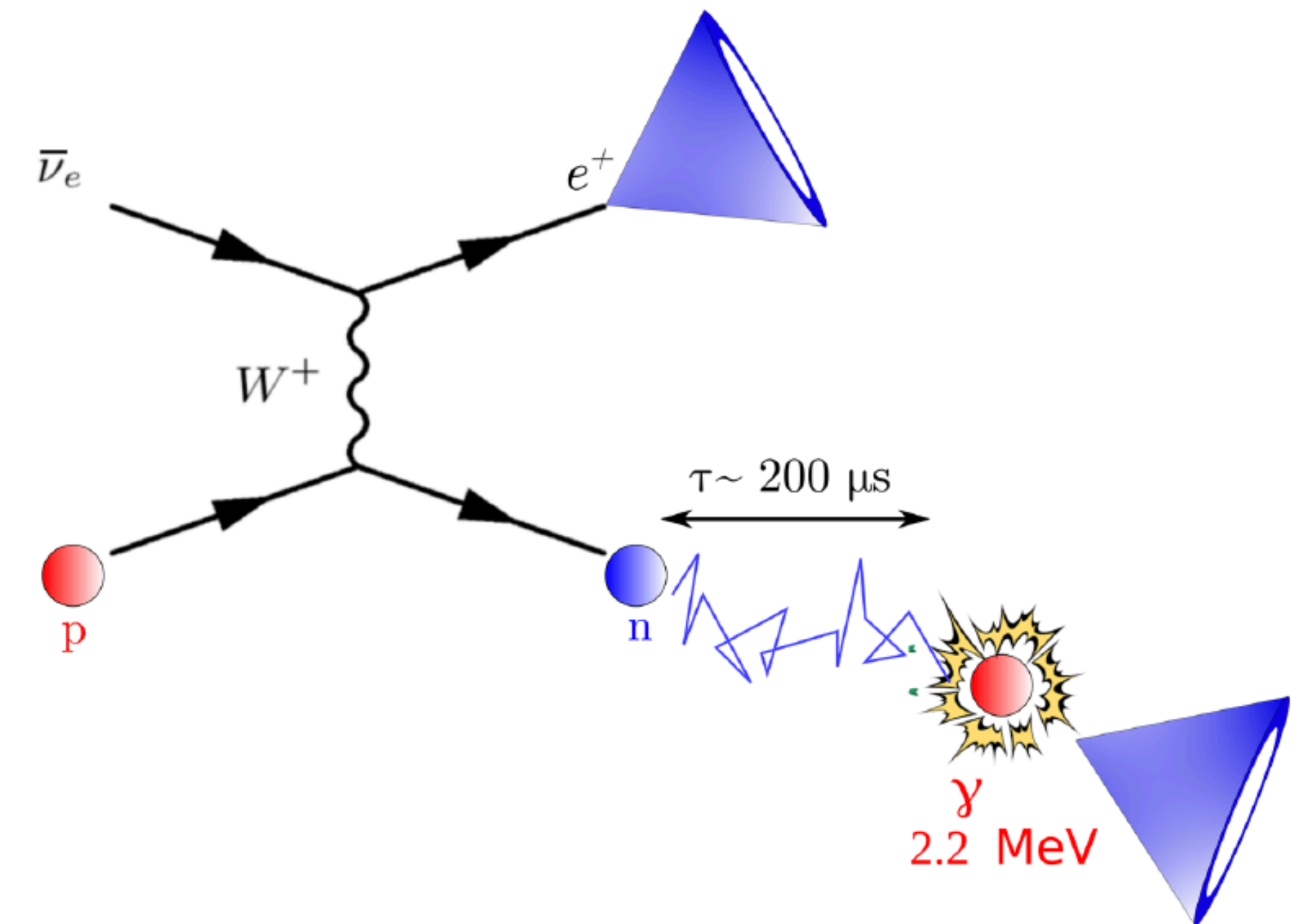
Not experimentally discovered so far...  
(best limits by Super-Kamiokande)

Overall scale by  
CCSN rate, etc



# Experimental Search at Water Cherenkov Detectors

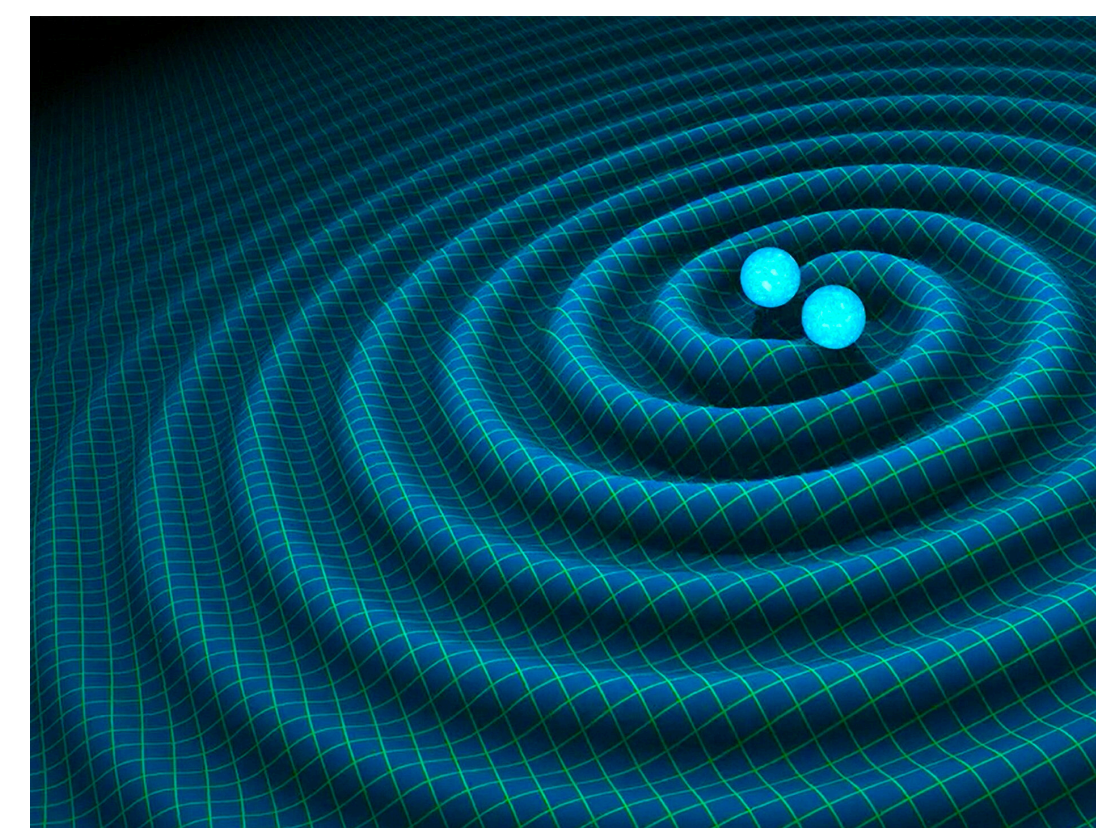
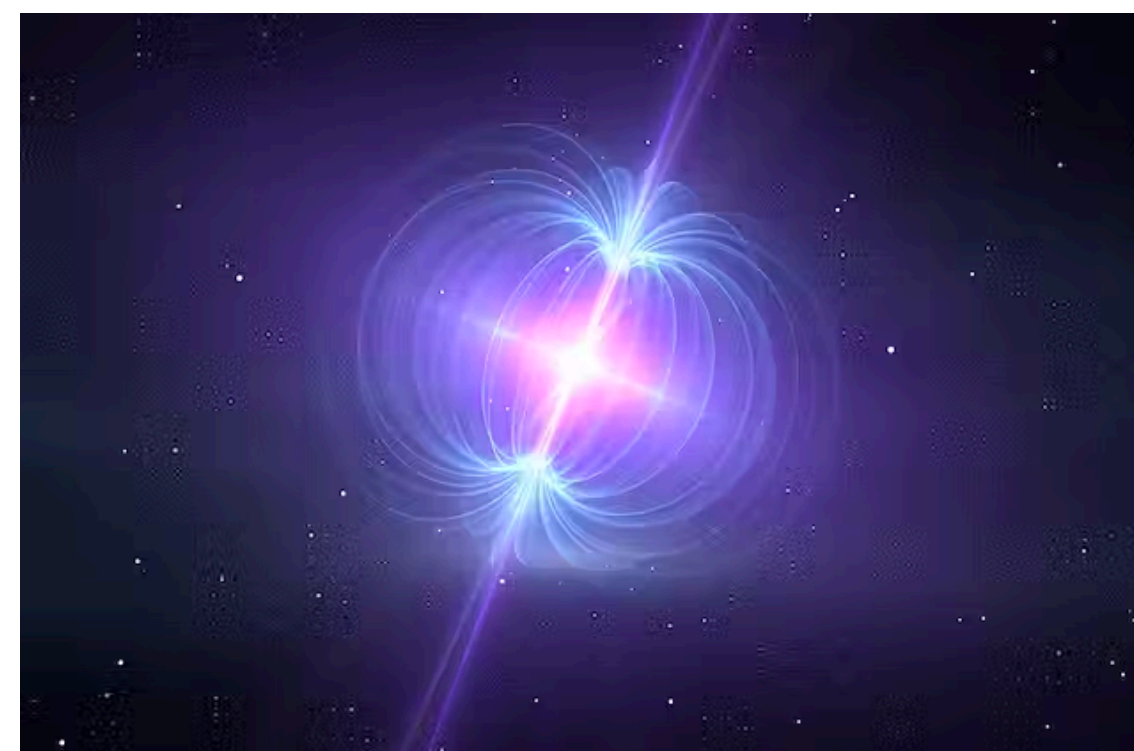
- Signal = **inverse beta decay (IBD)**,  $\bar{\nu}_e + p \rightarrow e^+ + n$  (largest cross section)
  - $e^+$  = “*prompt*” signal (main signal range: **10~30 MeV**)
  - $n$  = “*delayed*” signal via  **$\gamma$ -ray(s)** from thermal capture on hydrogen or gadolinium
- Many types of backgrounds mimicking this signature.
  - Atmospheric neutrinos
  - Radioactive isotopes produced by atmospheric muons
  - Solar neutrinos
  - Reactor neutrinos



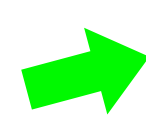


# Focus in This Study: “Fate” of Stellar Collapse

- A lot of DSNB models are proposed for recent years.
  - Many of them **end up serving flux**.
  - Most **lacking precise bkg estimation** in their sensitivity discussion.
- Focus on **stellar core collapse fate** (NS or BH), which is accessible by other observations, making **multi-messenger studies** possible.
  - Pulsars
  - Monitoring luminous stars for failed SNe
  - Gravitational wave observations for binaries



core collapse



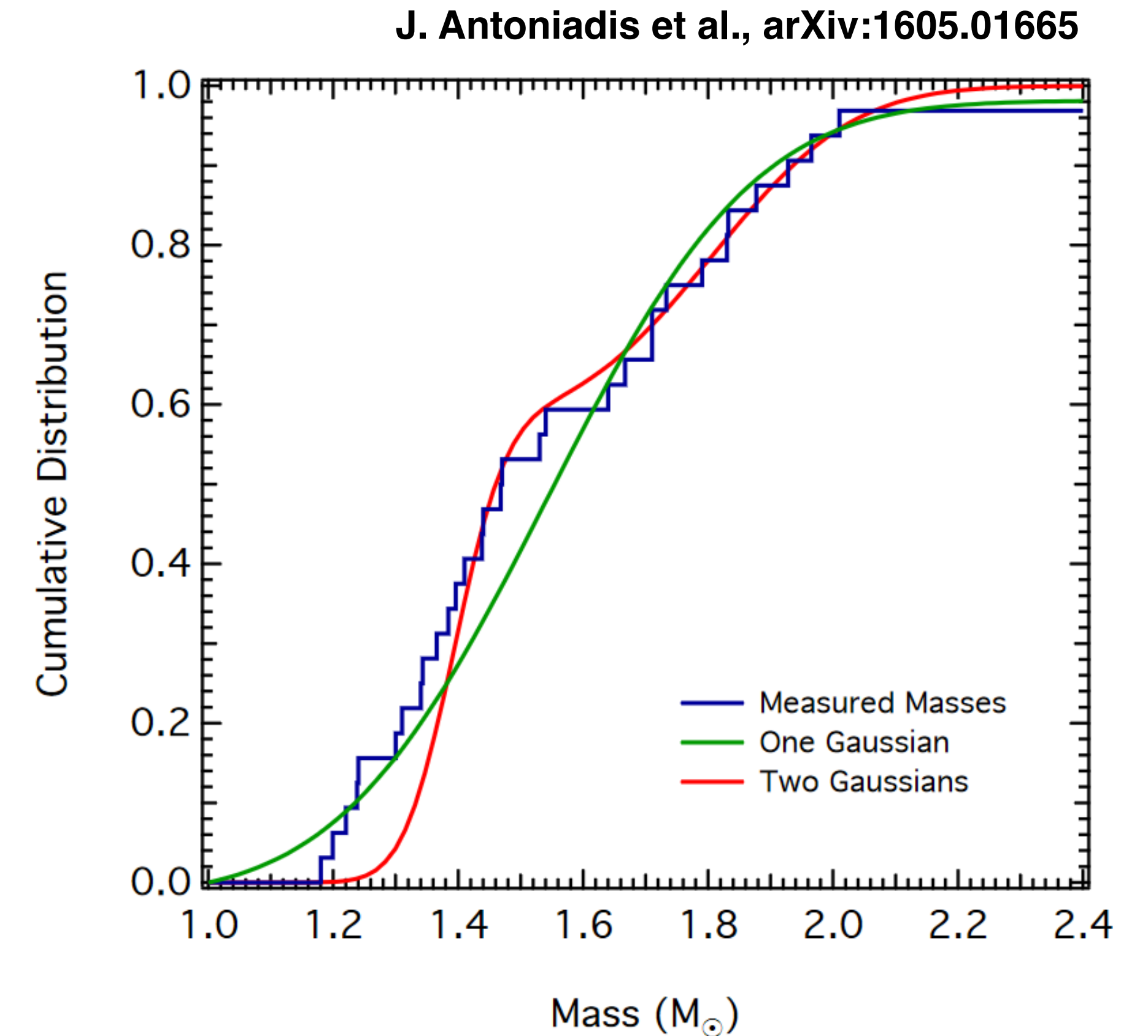
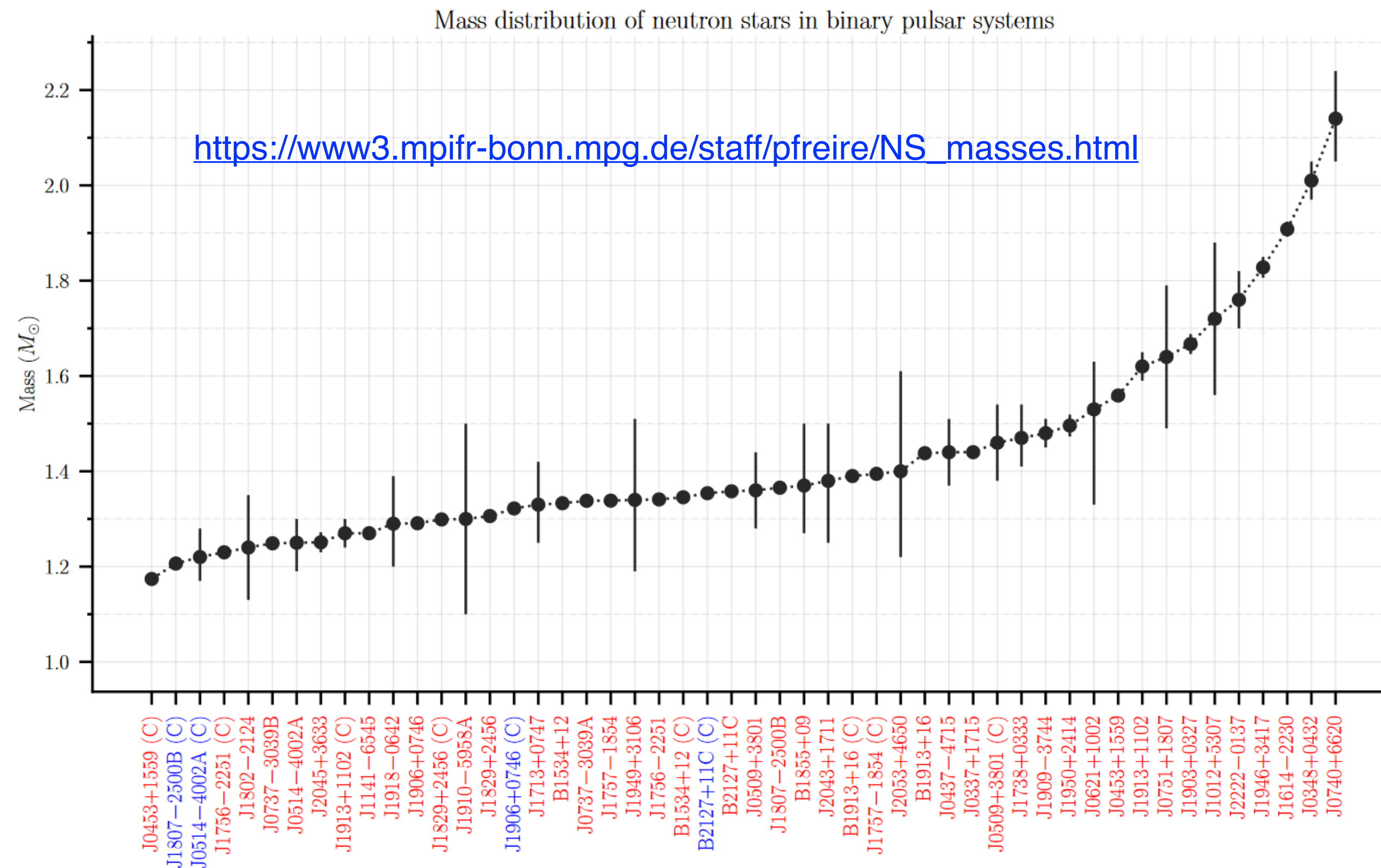
neutron star



black hole (failed SN)

# Ref: Observed Neutron Star Mass

- Neutron mass distribution from optical observations of the binary system shows a peak and higher tail.
- Natural born heavy, *or* gained mass through accretion from companion stars.



# Ref: Failed SN Fraction

- Monitoring luminous stars gave constraints on a failed SN fraction.
- 2 failed SN candidates (N6946-BH1, M101-OC1) out of 8 SNe.
  - Failed SN fraction  $\sim 4\text{--}39\%$  (90% C.L.), assuming  $N_{\text{FSN}} = 1$  and  $N_{\text{SN}} = 8$ .

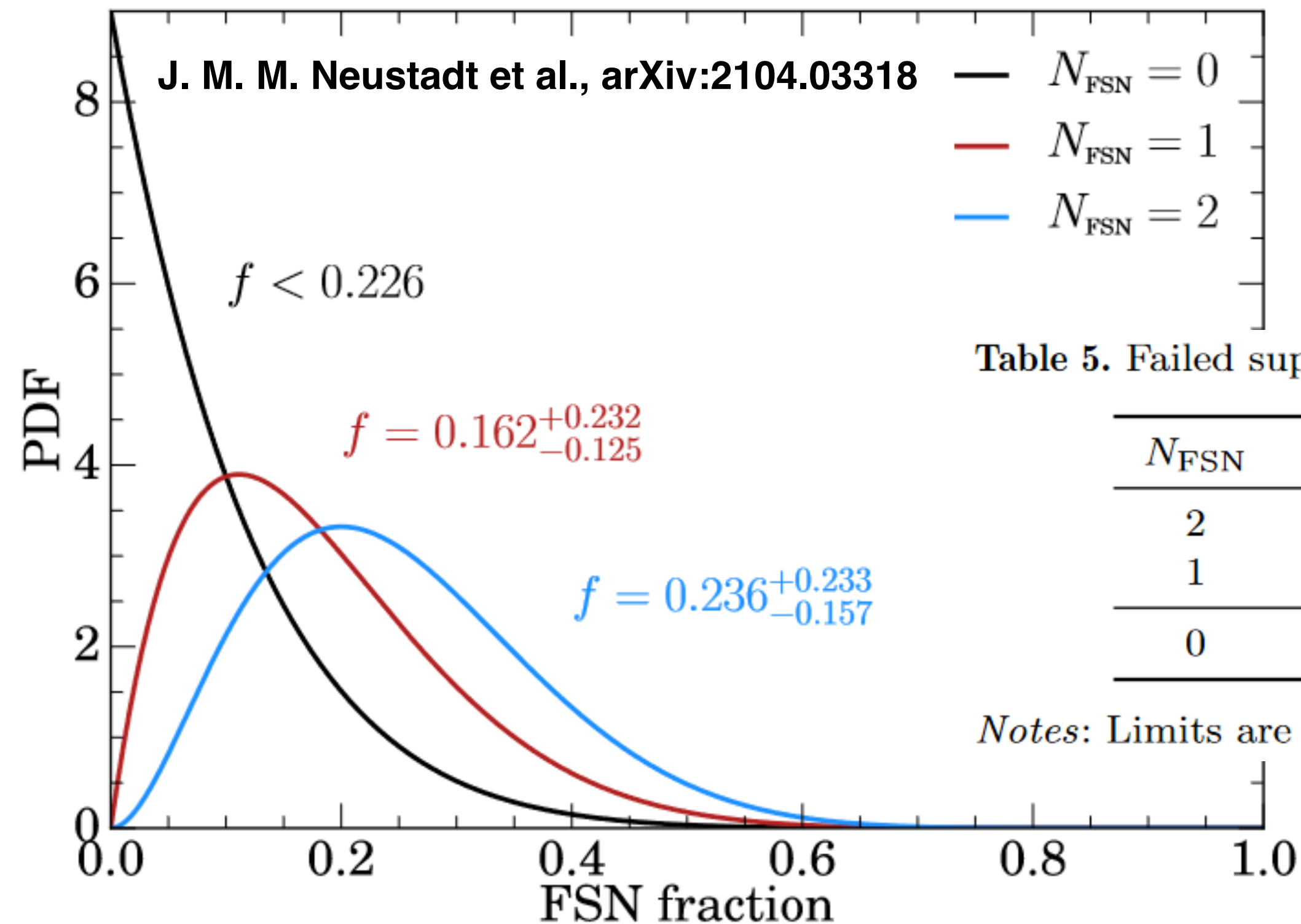
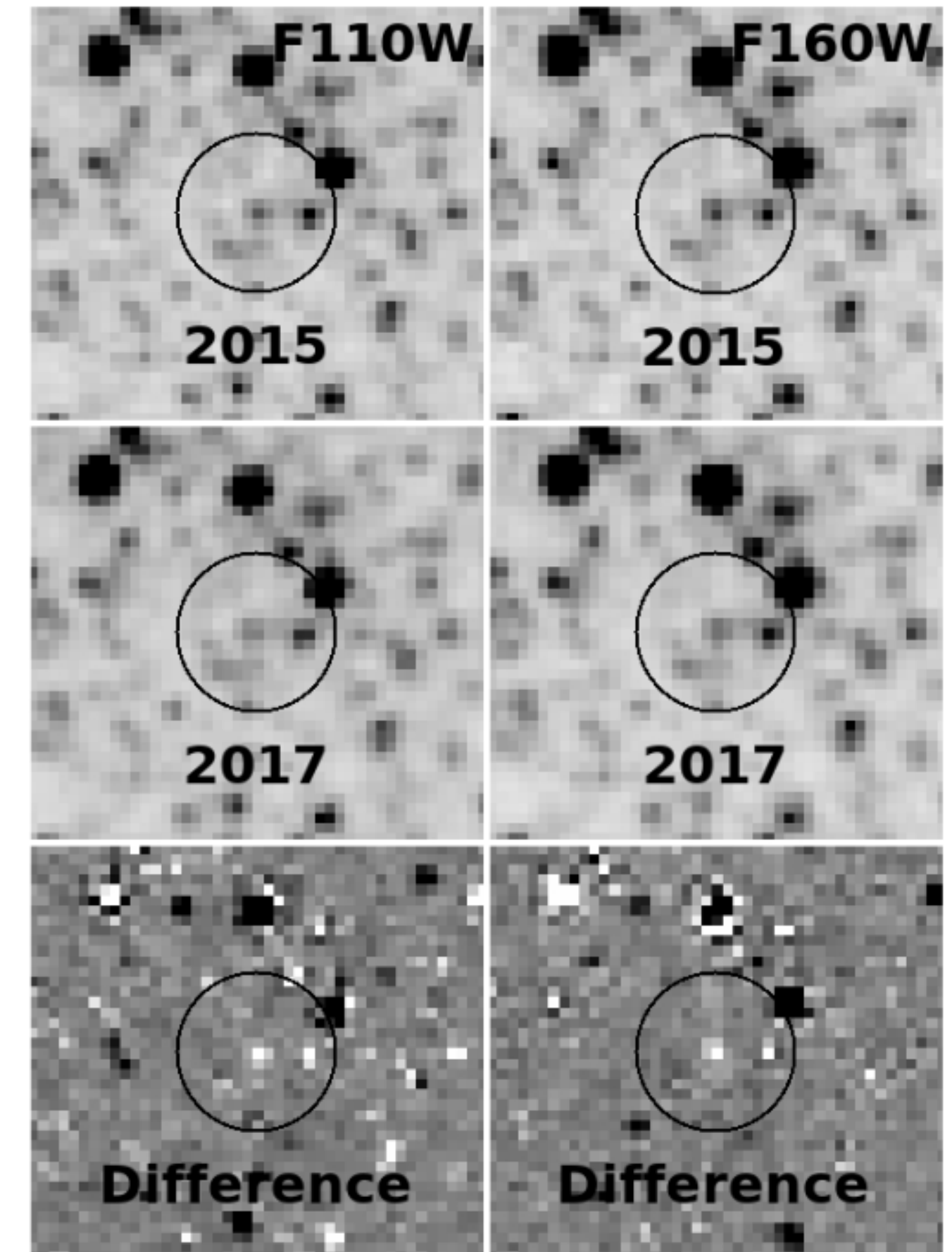


Table 5. Failed supernova/core-collapse fraction

$N_{\text{FSN}}$	Lower limit	Median	Upper limit
2	0.079	0.236	0.470
1	0.037	0.162	0.394
0	—	—	0.226

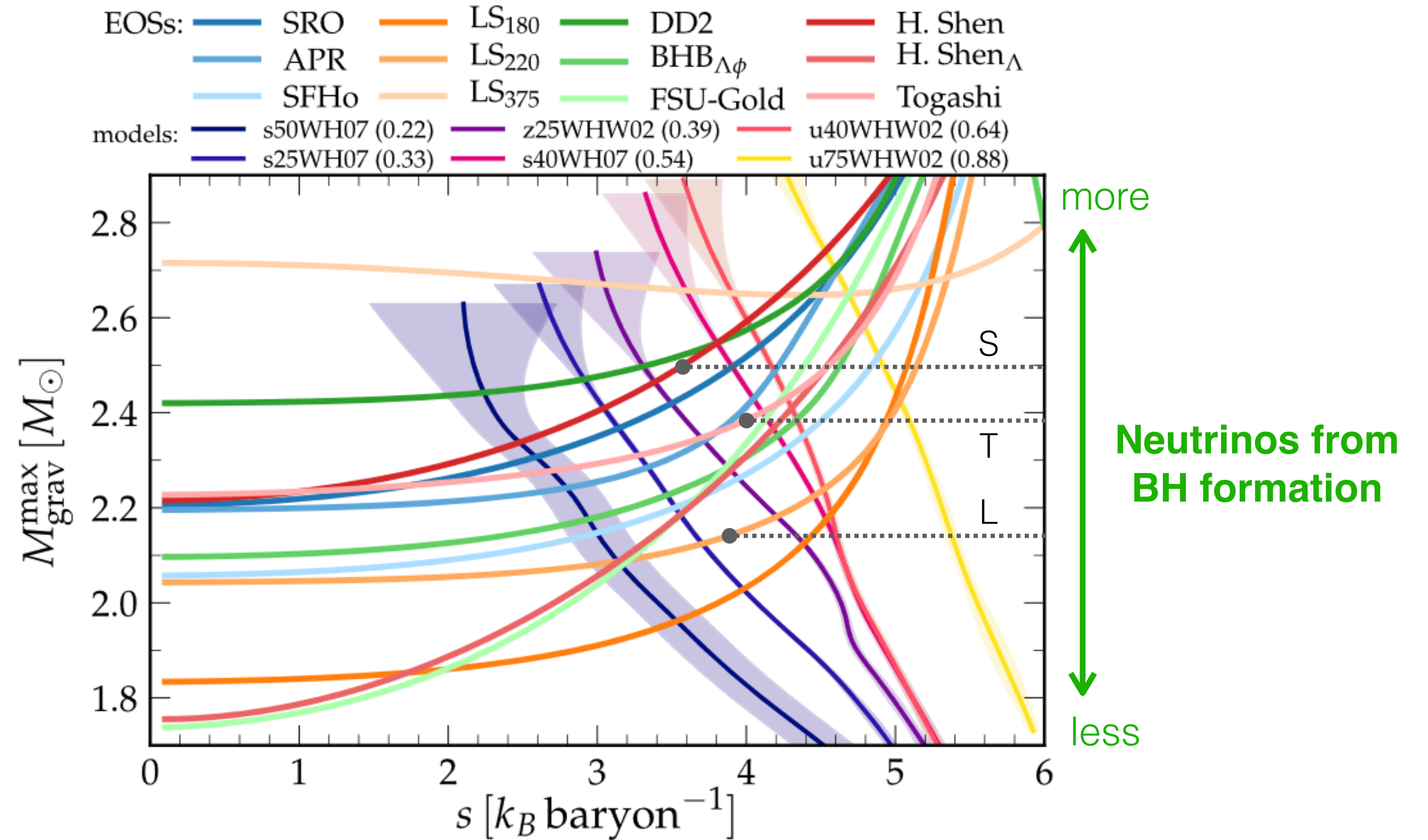
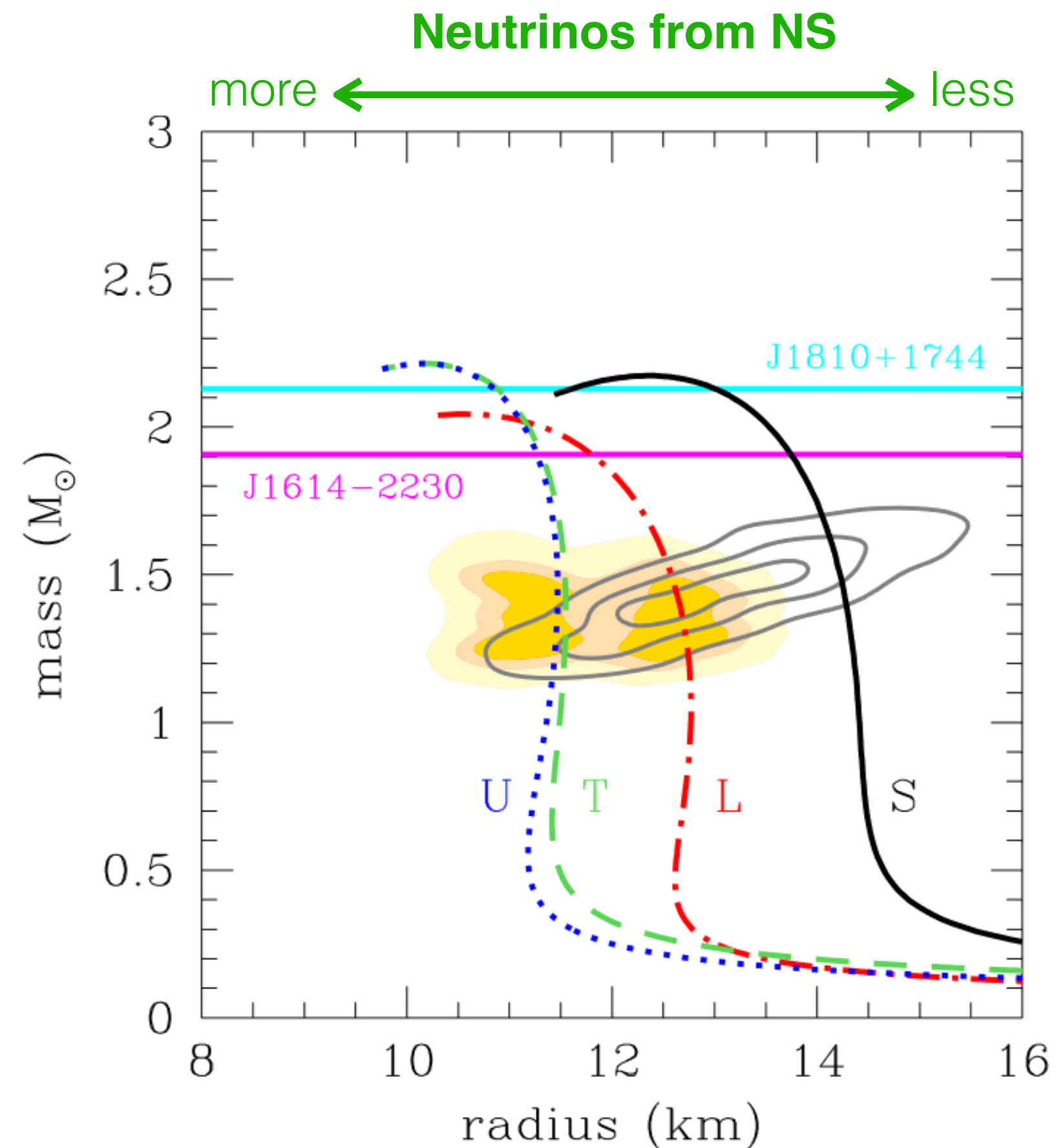
Notes: Limits are presented at the 90 per cent confidence level.

C. M. Basinger et al., arXiv:2007.15658



# Ref: Nuclear Equation-of-State Impact

- In the NS case, neutrino emission amount depends on **radius** of proto-NS.
- In the BH case, neutrino emission amount depends on **maximum mass** of proto-NS.





# Exploring the Fate of Stellar Core Collapse with Supernova Relic Neutrinos

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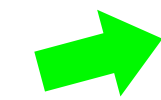
## Abstract

Core collapse of massive stars leads to different fates for various physical factors, which gives different spectra of the emitted neutrinos. We focus on the supernova relic neutrinos (SRNs) as a probe to investigate the stellar collapse fate. We present the SRN fluxes and event rate spectra at a detector for three resultant states after stellar core collapse, the typical mass neutron star, the higher mass neutron star, or the failed supernova forming a black hole, based on different nuclear equations of state. Then possible SRN fluxes are formed as mixtures of the three components. We also show the expected sensitivities at the next-generation water-based Cherenkov detectors, SK-Gd and Hyper-Kamiokande, as constraining the mixture fractions. This study provides a practical example of extracting astrophysical constraints through SRN measurement.

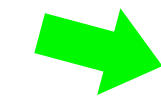
*Unified Astronomy Thesaurus concepts:* [Neutrino astronomy \(1100\)](#); [Supernova neutrinos \(1666\)](#); [Core-collapse supernovae \(304\)](#); [Massive stars \(732\)](#); [Neutron stars \(1108\)](#); [Black holes \(162\)](#)

- Emitted neutrino spectrum is expected to depend on the remnant after core collapse (“*fate*”).
- Consider three major cases as a fate.
  - Canonical mass neutron stars ( $1.47 M_{\text{Sun}}$ )
  - High mass neutron stars ( $1.86 M_{\text{Sun}}$ )
  - Black holes (failed SNe)
- Use numerical simulations results for neutrino spectrum for each case from K. Nakazato et al., ApJ Suppl. 205, 2 (2013) & K. Nakazato et al., ApJ 925, 98 (2022).
- Calculate DSNB flux, following the techniques in K. Nakazato et al., ApJ 804, 75 (2015).

core collapse



neutron star (canonical or heavy)



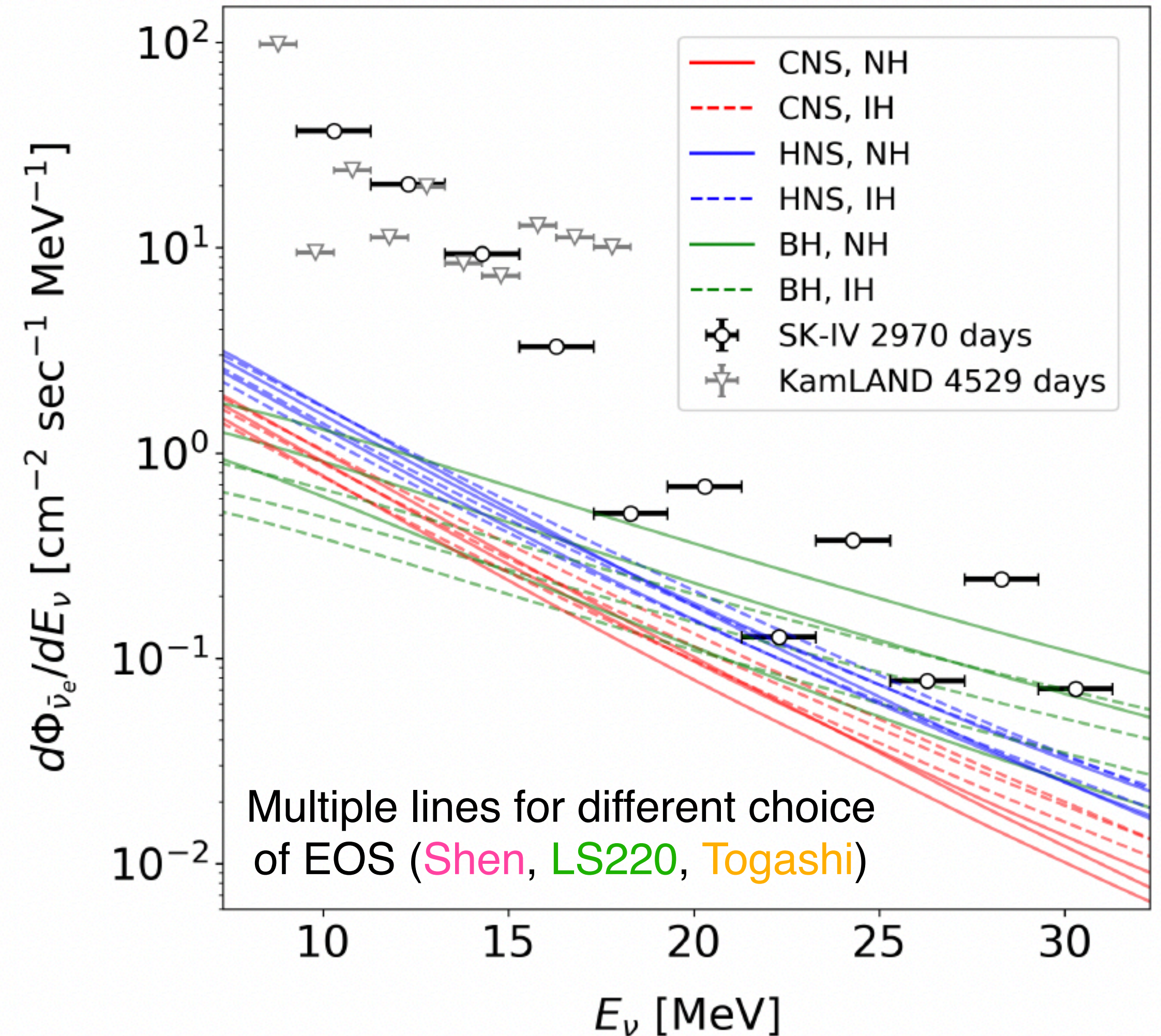
black hole (failed SN)

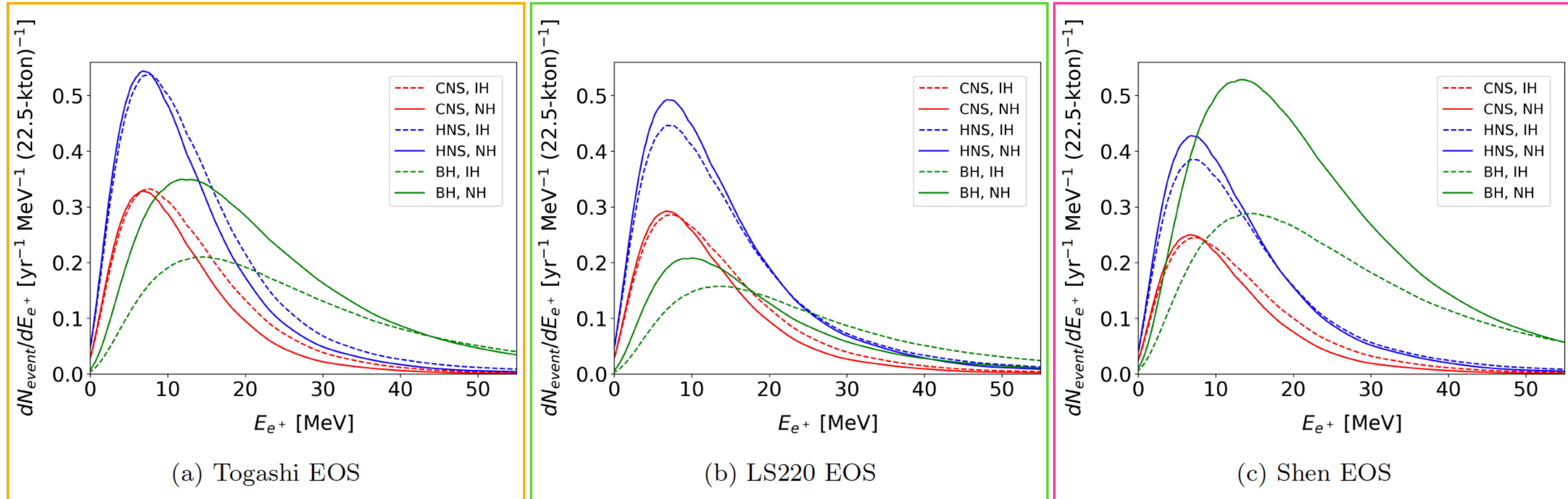
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- Calculate DSNB flux, following the techniques in K. Nakazato et al. (2013)

core collapse



- neutron star (canonical or heavy)
- black hole (failed SN)



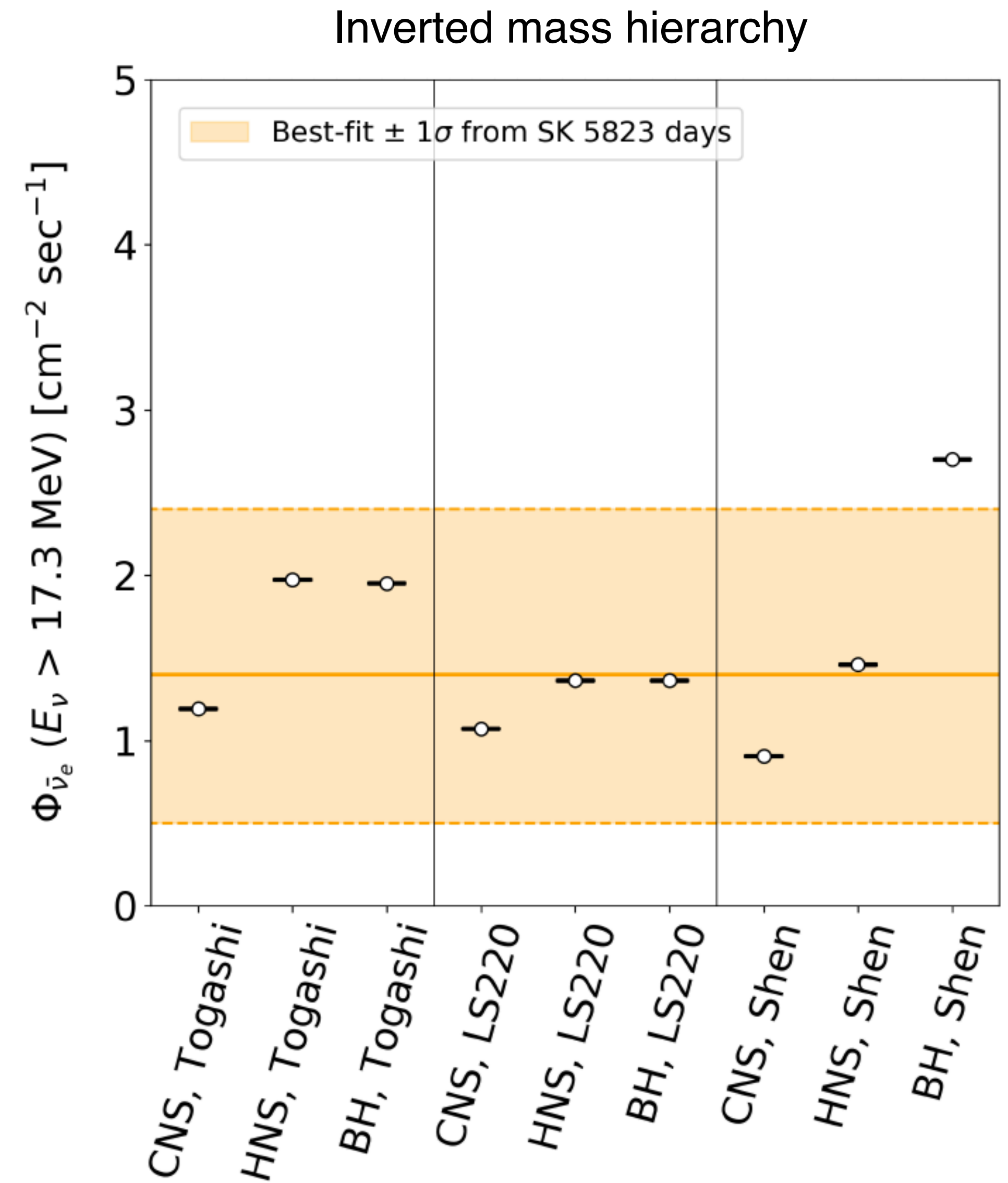
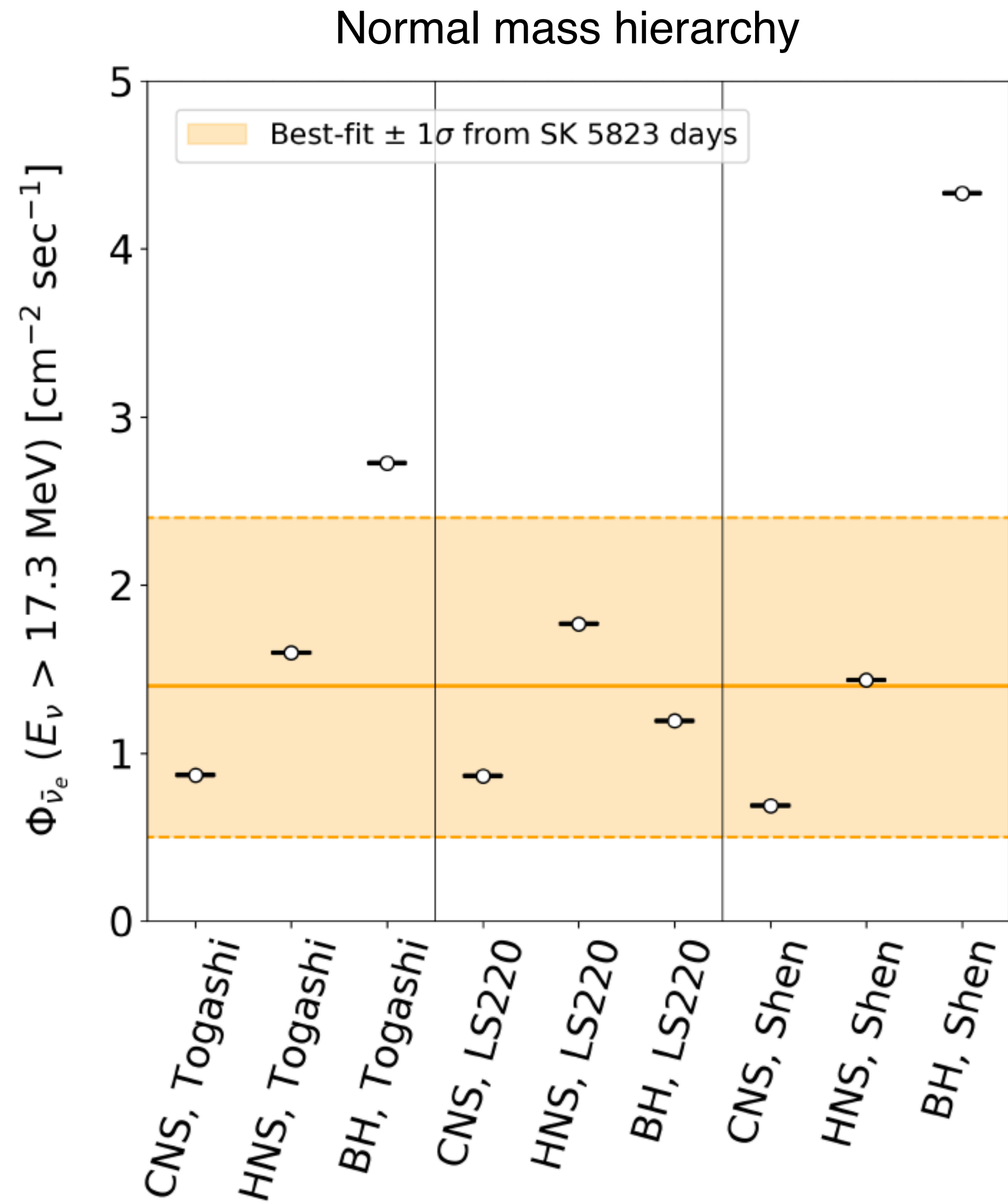


**canonical mass NS, high mass NS, BH**

- HNS is generally scale up of CNS. BH serves higher energies than NS.
- EOS impact is larger for BH.
- Neutrino mass hierarchy affects differently for NS and BH (also depends on EOS type).



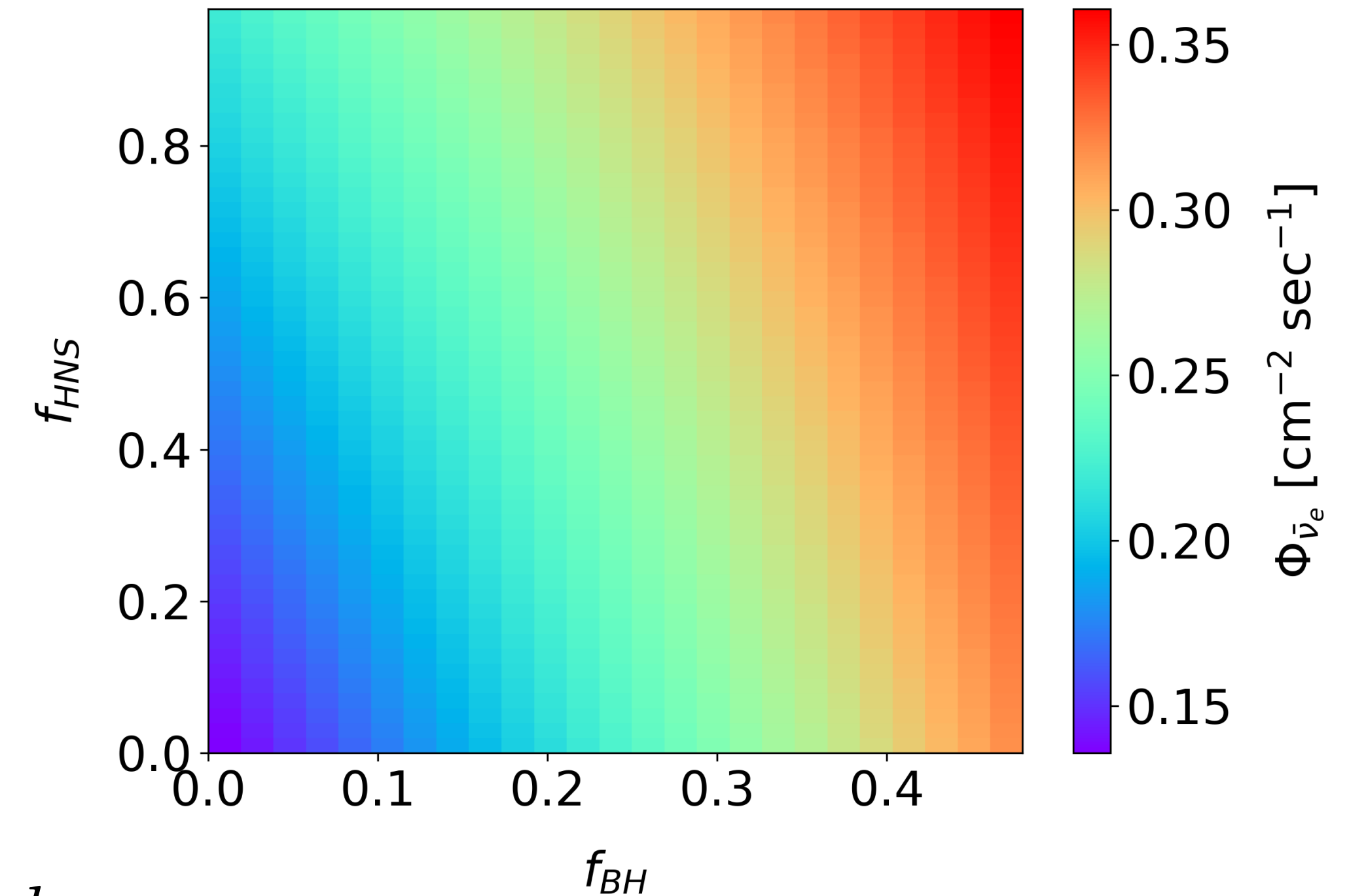
# Integrated Flux ( $>17.3$ MeV)



# Mixture of Three Fluxes

- Mix three cases by fractional parameters,  $f_{\text{HNS}}$  and  $f_{\text{BH}}$ , to form a DSNB flux.
  - $f_{\text{BH}}$ : fraction of BH to total core-collapse
  - $f_{\text{HNS}}$ : fraction of high mass NS to total NS
- Integrated DSNB flux differs for  $\{f_{\text{HNS}}, f_{\text{BH}}\}$  combinations.

(Example) Integrated DSNB flux for a certain energy range for different  $\{f_{\text{HNS}}, f_{\text{BH}}\}$  combinations



DSNB flux at Earth

Neutrino spectrum from a core-collapse

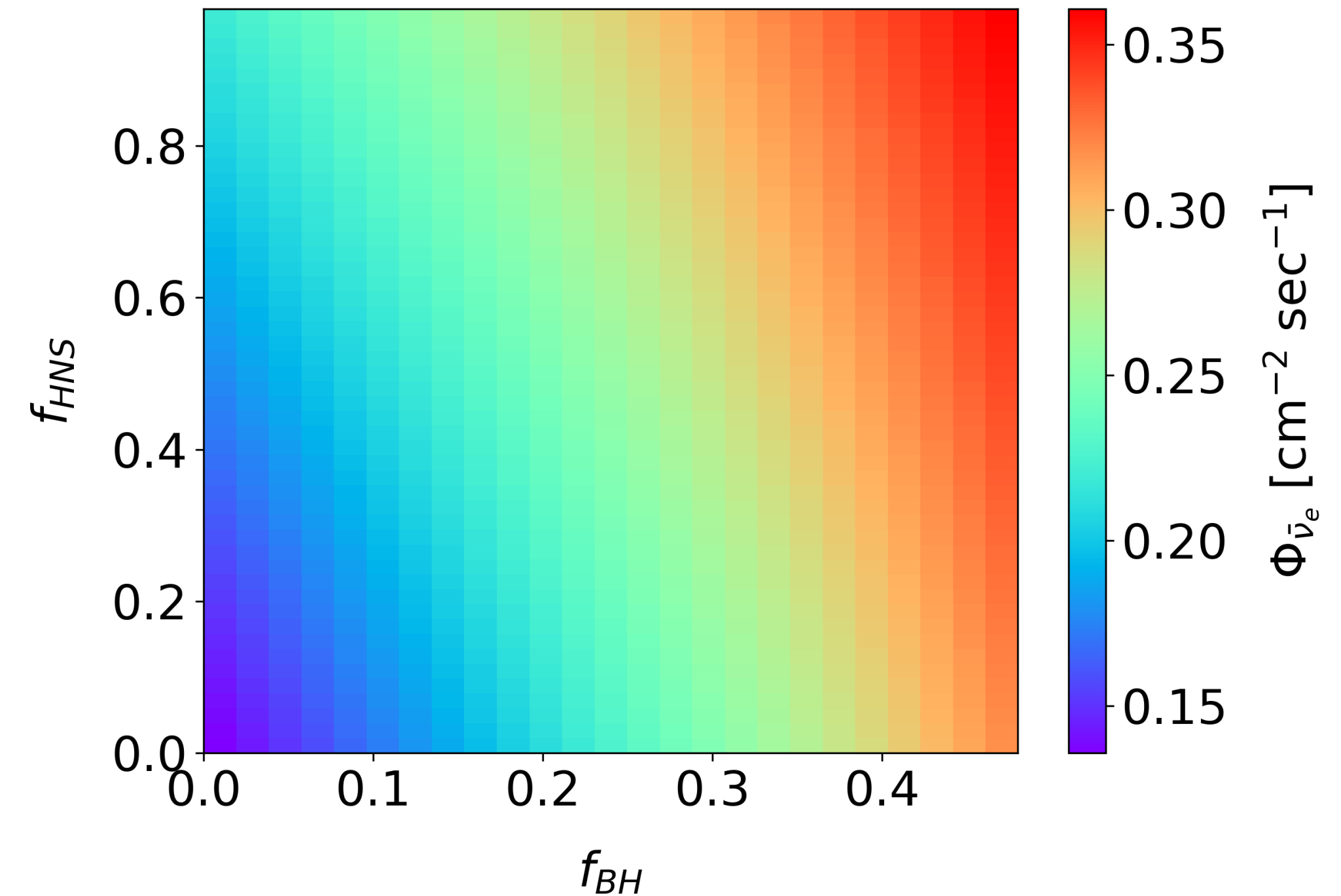
$$\frac{d\Phi(E_\nu)}{dE_\nu} = c \int_0^{z_{\text{max}}} R_{\text{CC}}(z) \left\langle \frac{dN(E'_\nu)}{dE'_\nu} \right\rangle \times \frac{dz}{H_0 \sqrt{\Omega_m (1+z)^3 + \Omega_\Lambda}}$$

$$\left\langle \frac{dN(E'_\nu)}{dE'_\nu} \right\rangle = f_{\text{BH}} \frac{dN_{\text{BH}}(E'_\nu)}{dE'_\nu} + (1 - f_{\text{BH}}) \times \left[ f_{\text{HNS}} \frac{dN_{\text{HNS}}(E'_\nu)}{dE'_\nu} + (1 - f_{\text{HNS}}) \frac{dN_{\text{CNS}}(E'_\nu)}{dE'_\nu} \right]$$

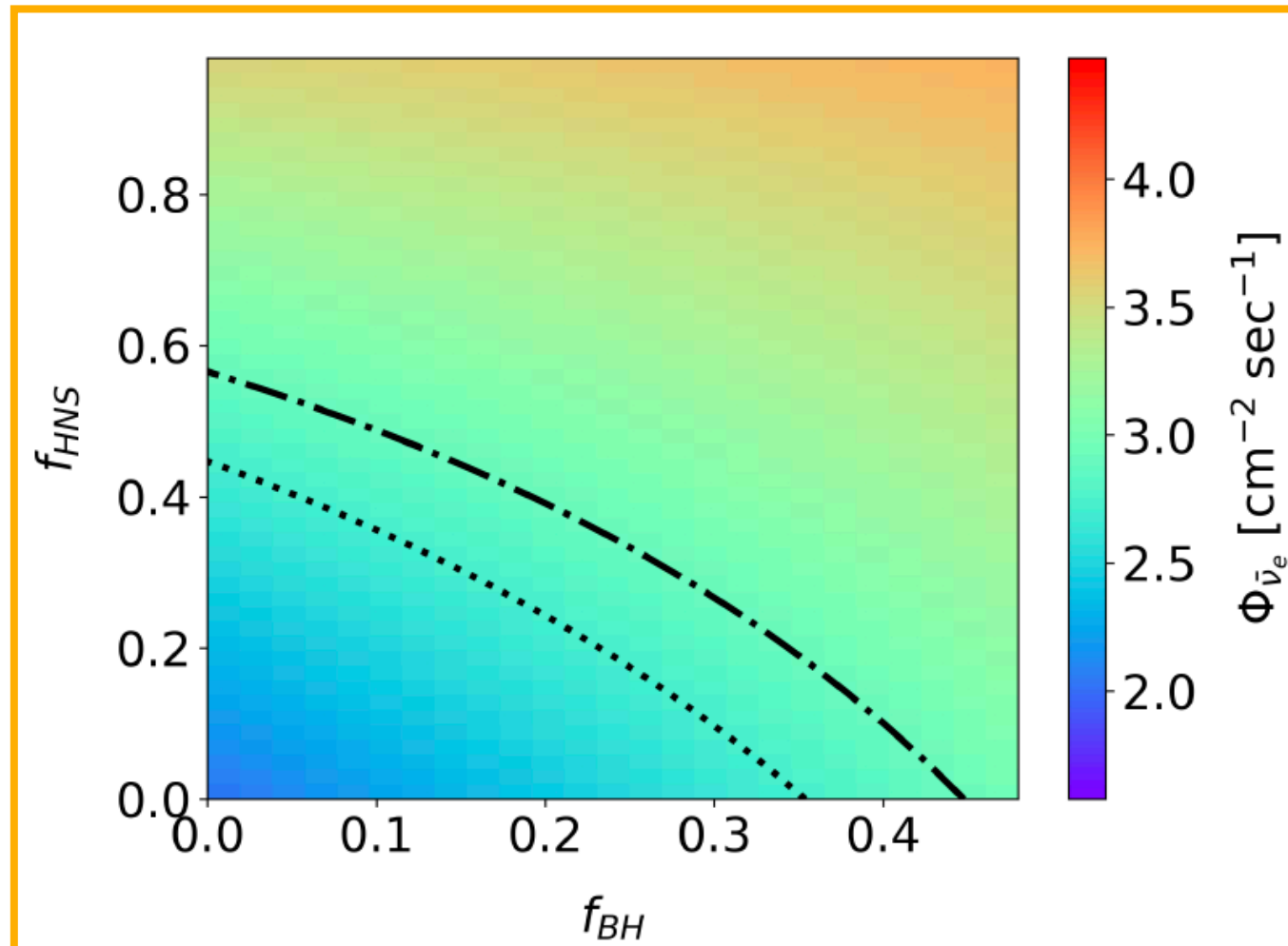
# Sensitivity on Fractional Parameters

- Mix three cases by fractional parameters,  $f_{\text{HNS}}$  and  $f_{\text{BH}}$ , to form a DSNB flux.
  - $f_{\text{BH}}$ : fraction of BH to total core-collapse
  - $f_{\text{HNS}}$ : fraction of high mass NS to total NS
- Perform experimental sensitivity study as **extrapolated from the SK-IV analysis about bkg**, drawing contours in  $\{f_{\text{HNS}}, f_{\text{BH}}\}$  parameter range, based on integrated flux.
- Consider two next-generation detectors.
  - **SK-Gd**:  $\times 1/10$  accidental bkg, 70% ntag efficiency
  - **Hyper-Kamiokande**:  $\times 8.4$  detector mass, same ntag efficiency

(Example) Integrated DSNB flux for a certain energy range for different  $\{f_{\text{HNS}}, f_{\text{BH}}\}$  combinations



Normal mass ordering



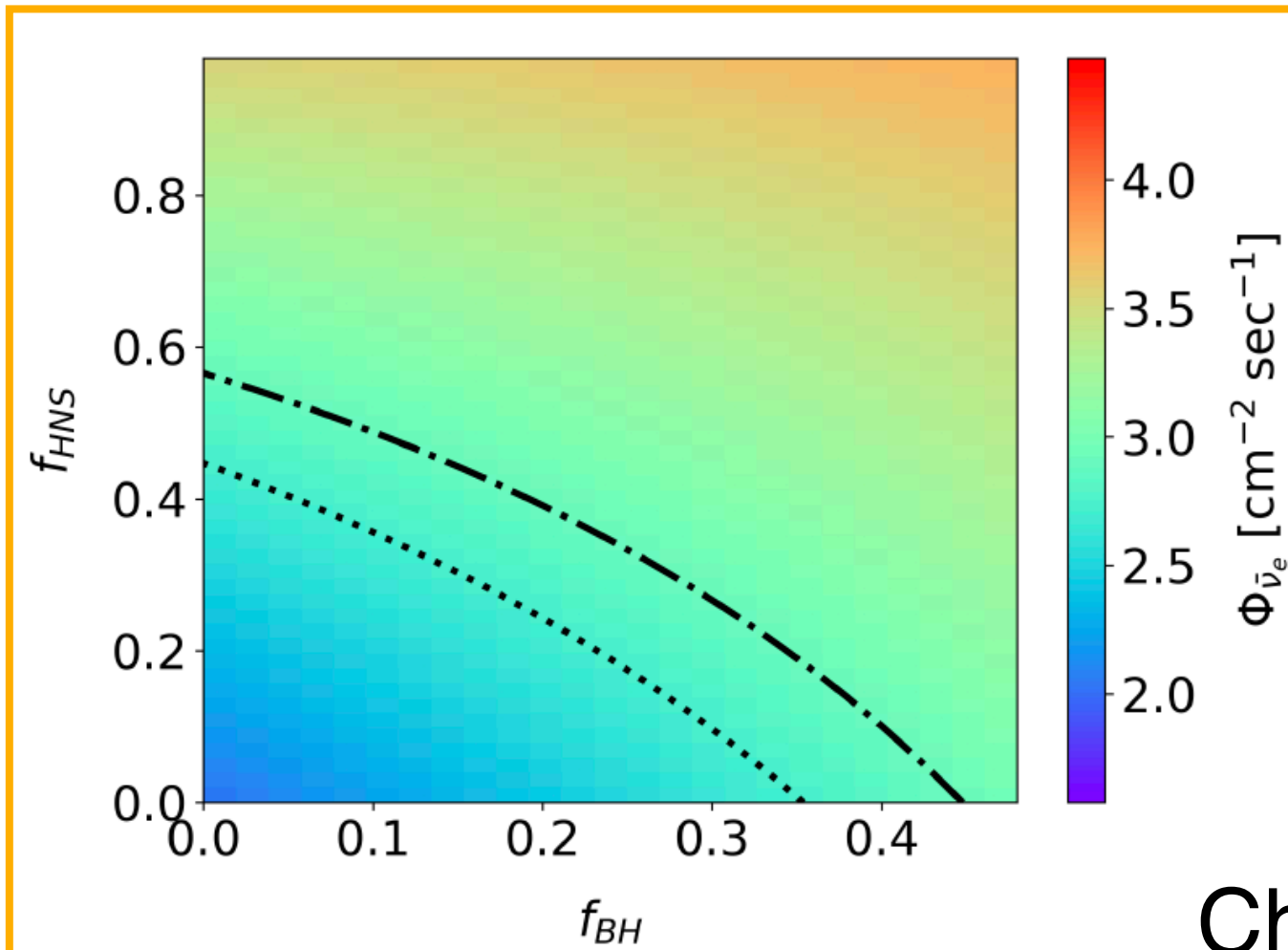
Integrated flux for  $13.3 < E_\nu < 31.3$  MeV

(1a) Togashi EOS,  
 $13.3 < E_\nu < 31.3$  MeV

Togashi

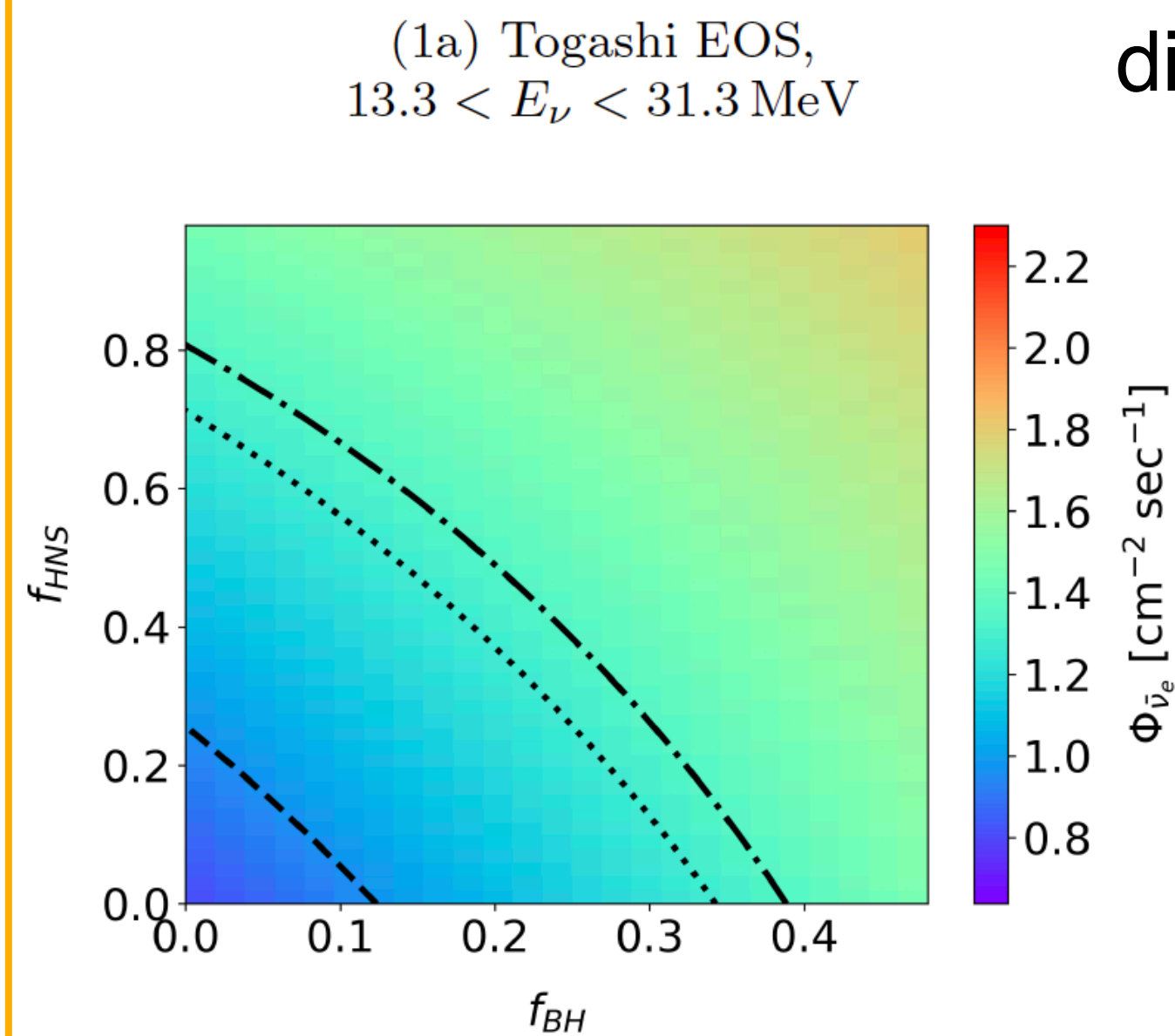
- · — · — SK-Gd 10yr
- ..... Hyper-K 3yr
- Hyper-K 5yr

Normal mass ordering



Integrated flux for  $13.3 < E_\nu < 31.3$  MeV

Choice of integration range serves different contours in some cases.



Integrated flux for  $17.3 < E_\nu < 31.3$  MeV

- · — · SK-Gd 10yr
- ..... Hyper-K 3yr
- Hyper-K 5yr

**Togashi** (2a) Togashi EOS,  $17.3 < E_\nu < 31.3$  MeV

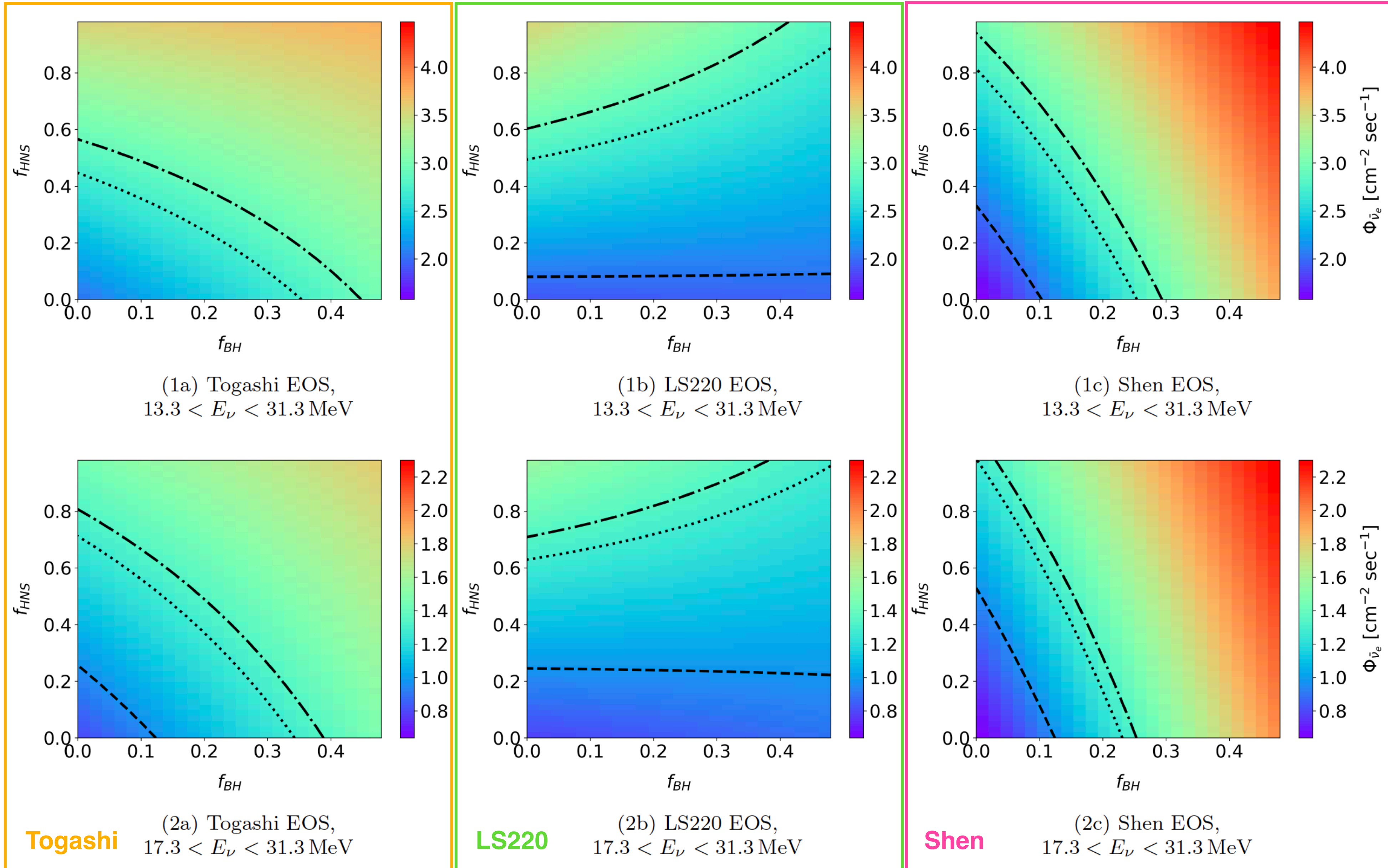
# Experimental Sensitivity (2 $\sigma$ C.L.)

Detectable above lines

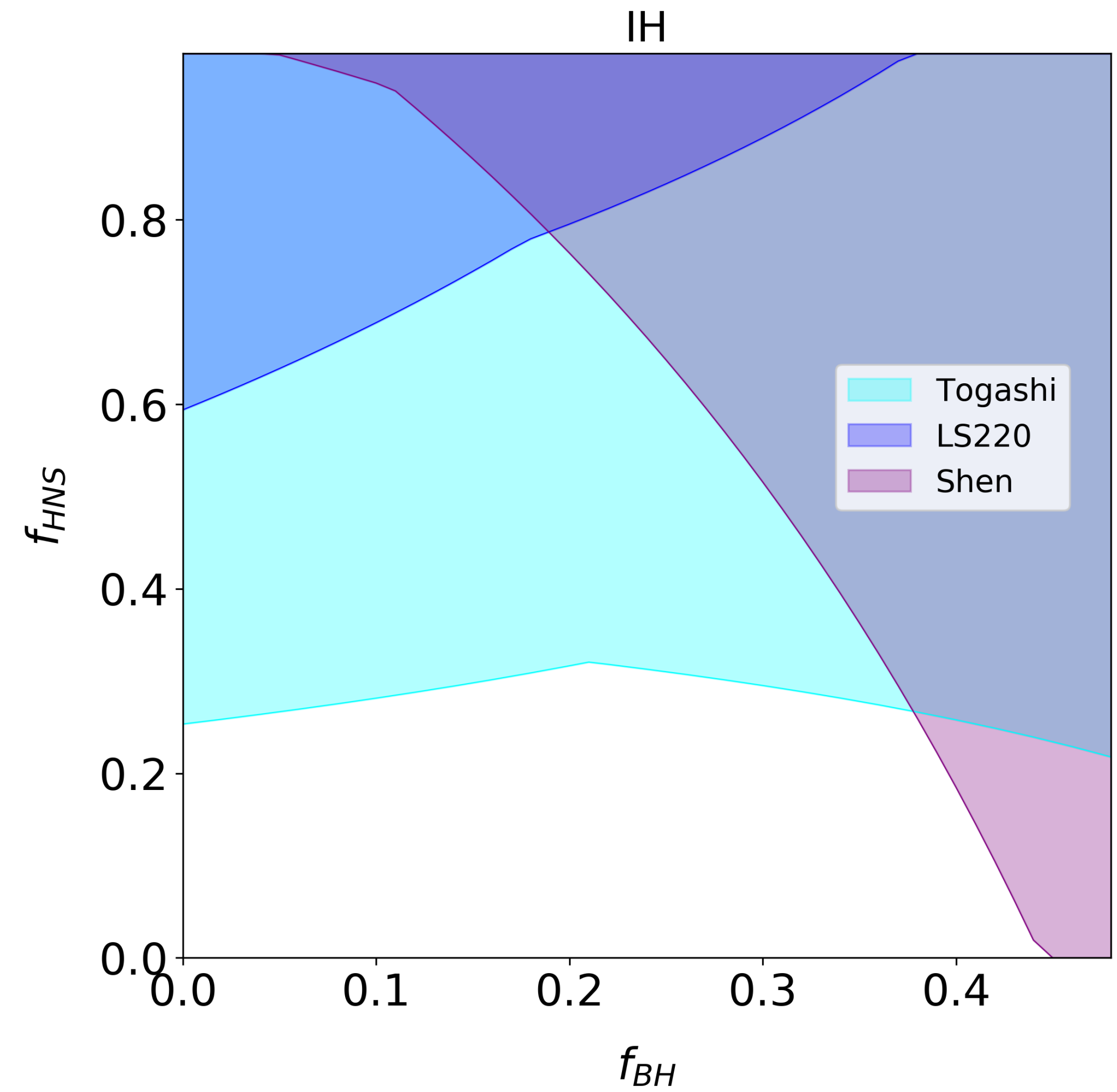
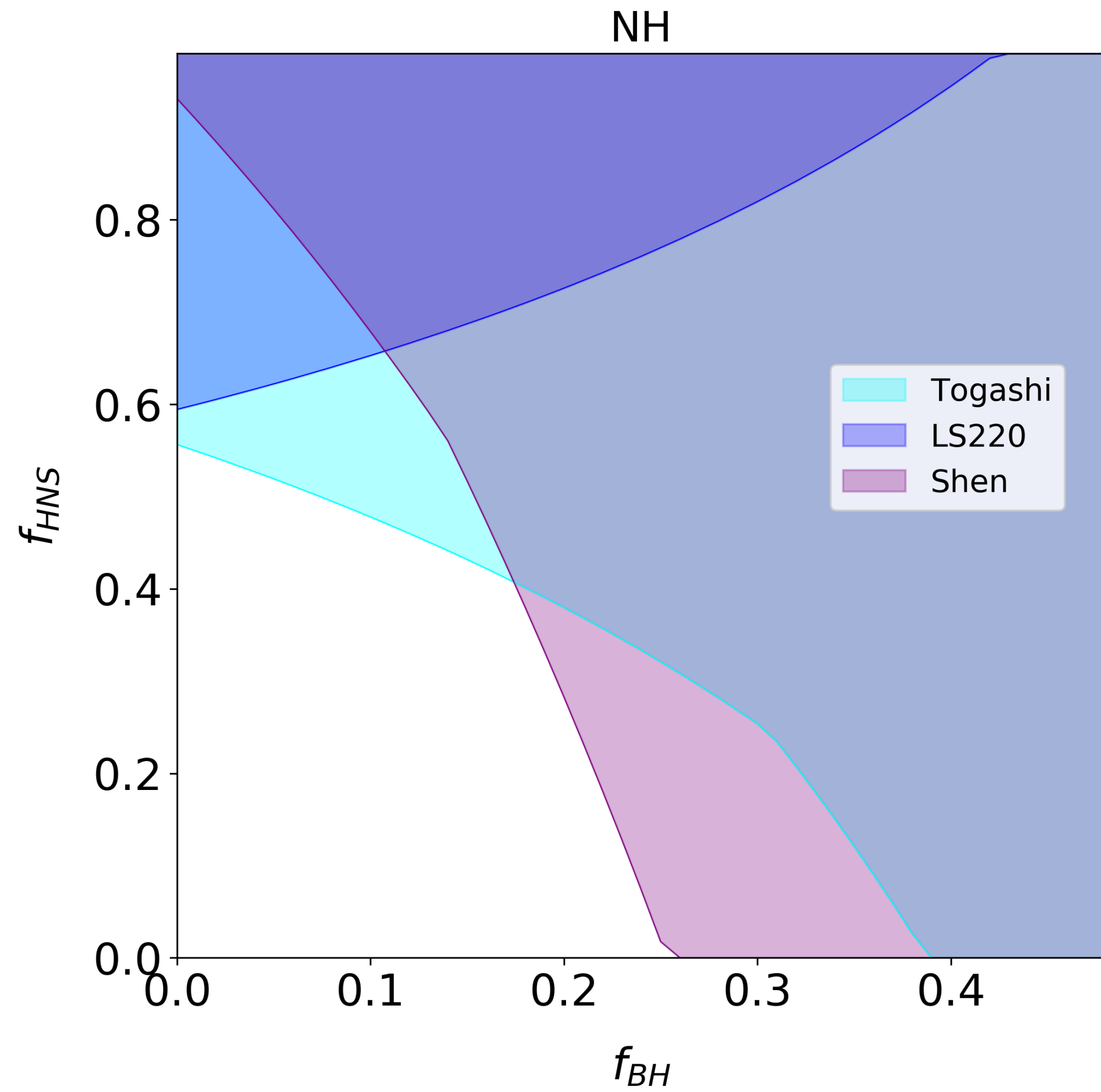
Normal mass ordering

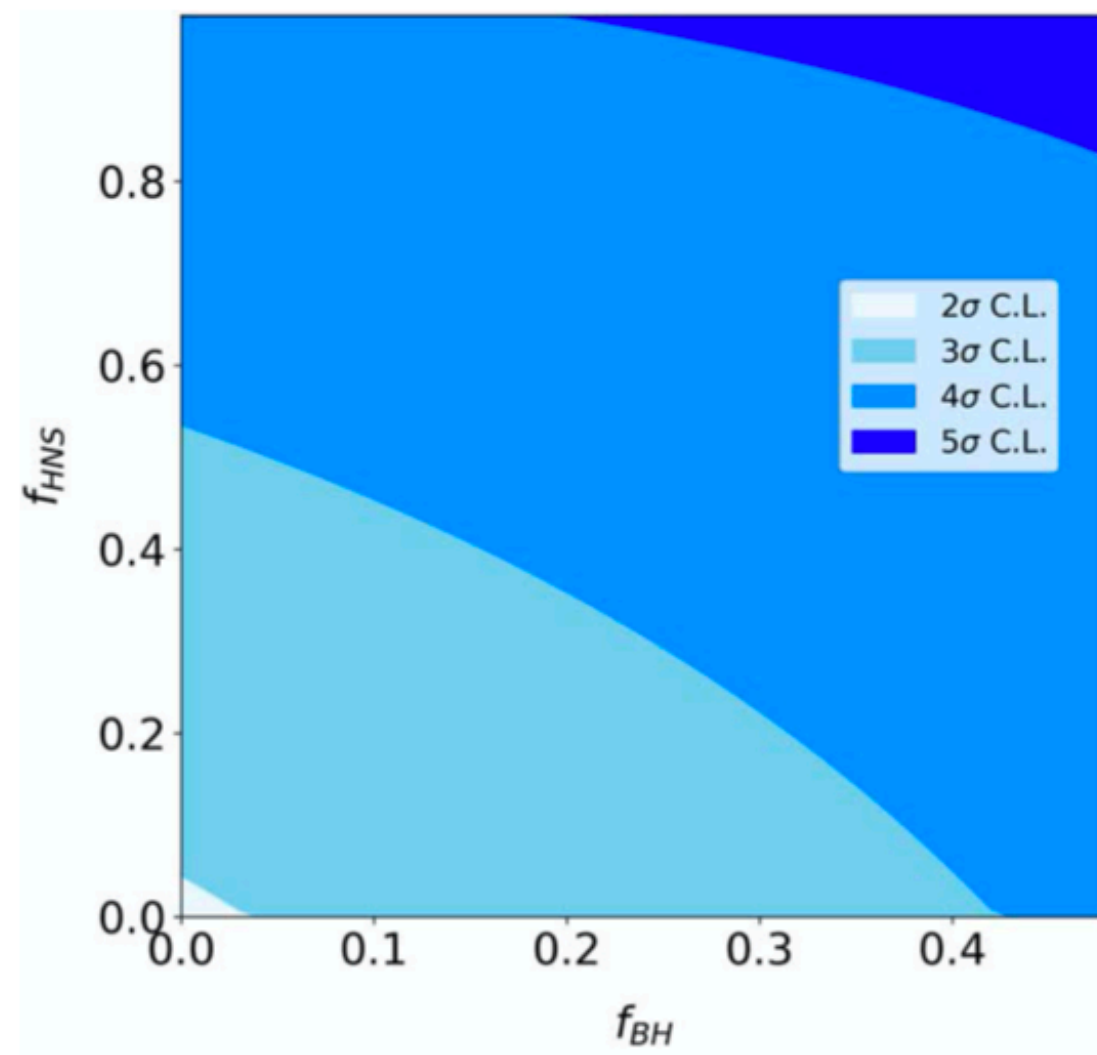
Contours change for different EOSs.

- SK-Gd 10yr
- ..... Hyper-K 3yr
- Hyper-K 5yr

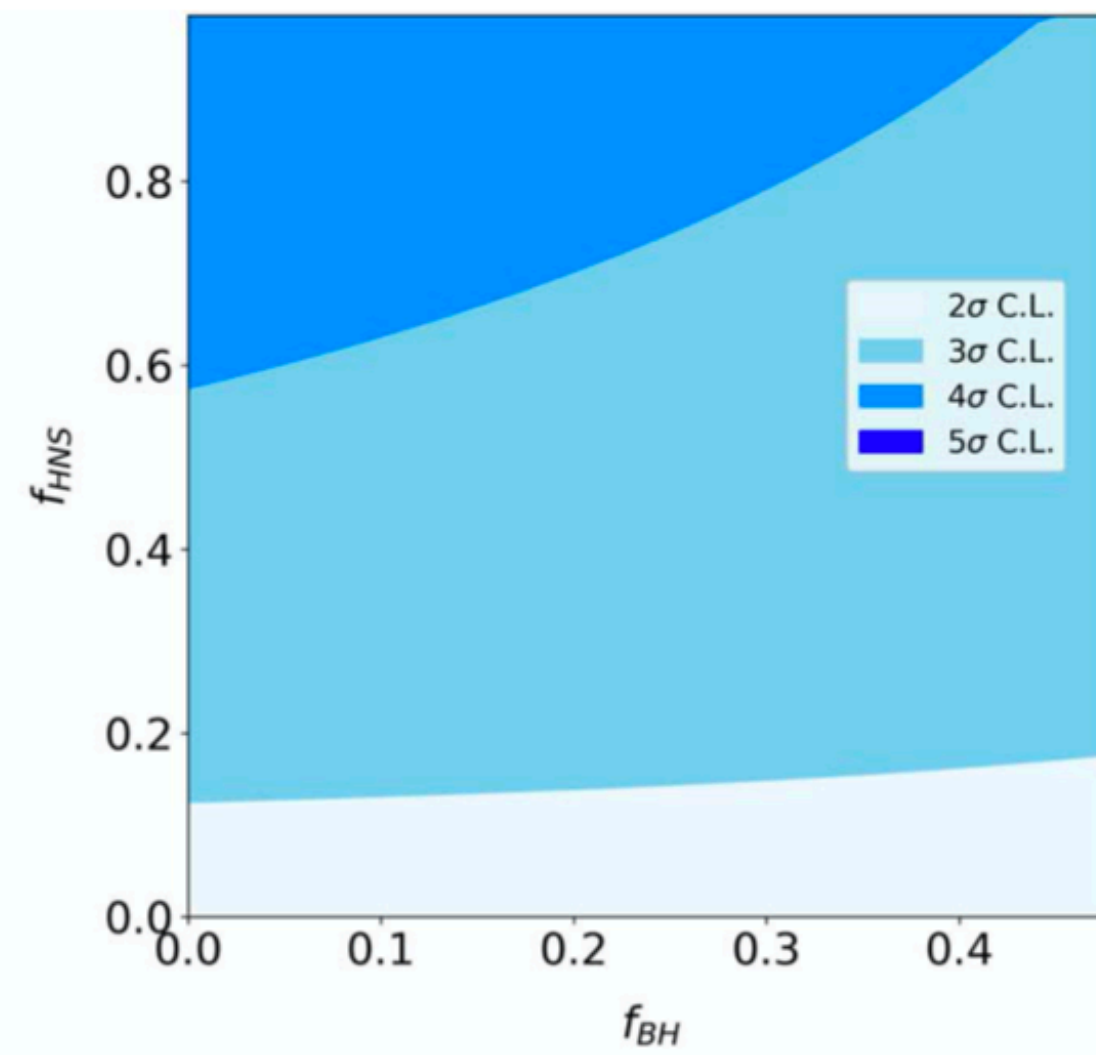


# Nuclear EOS Dependence

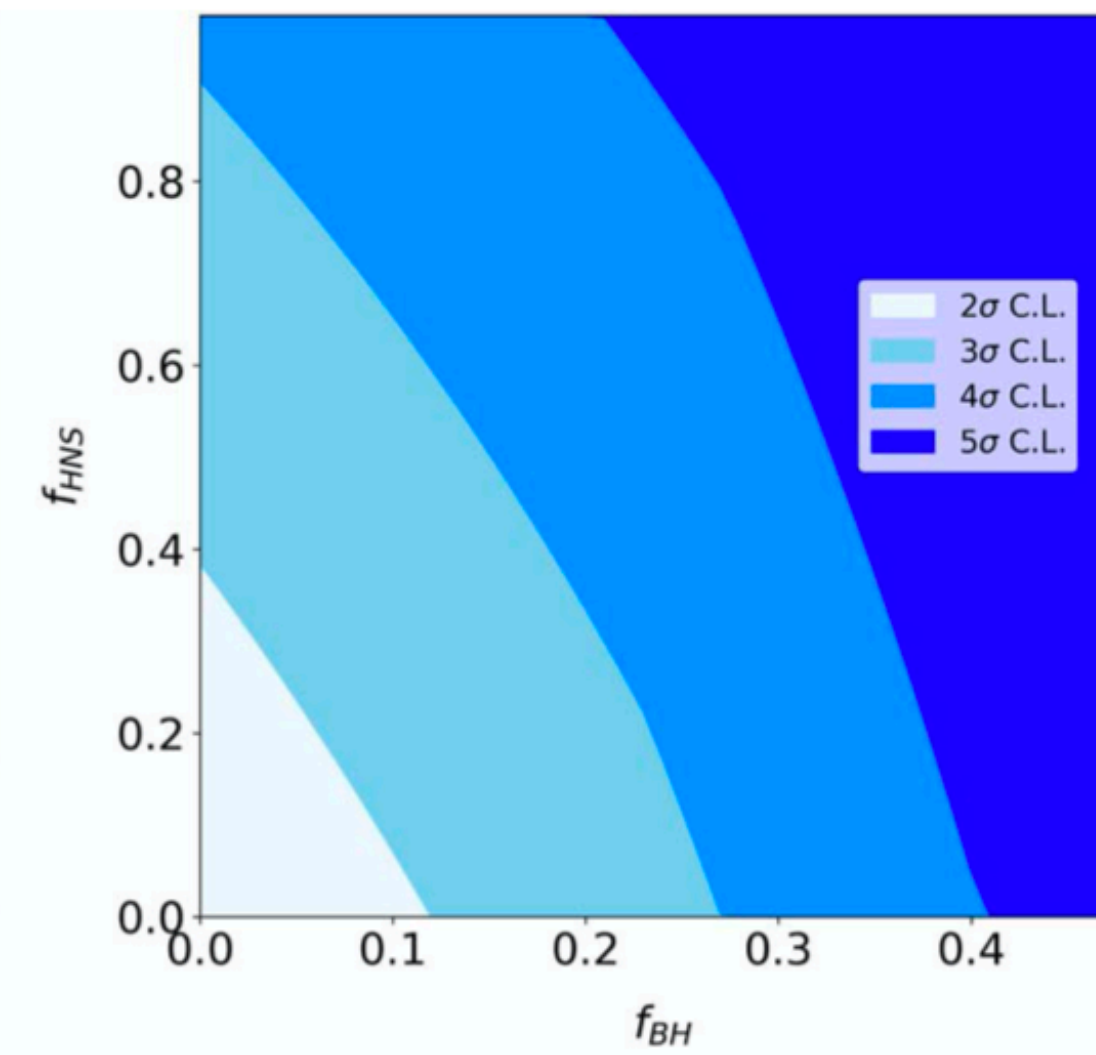




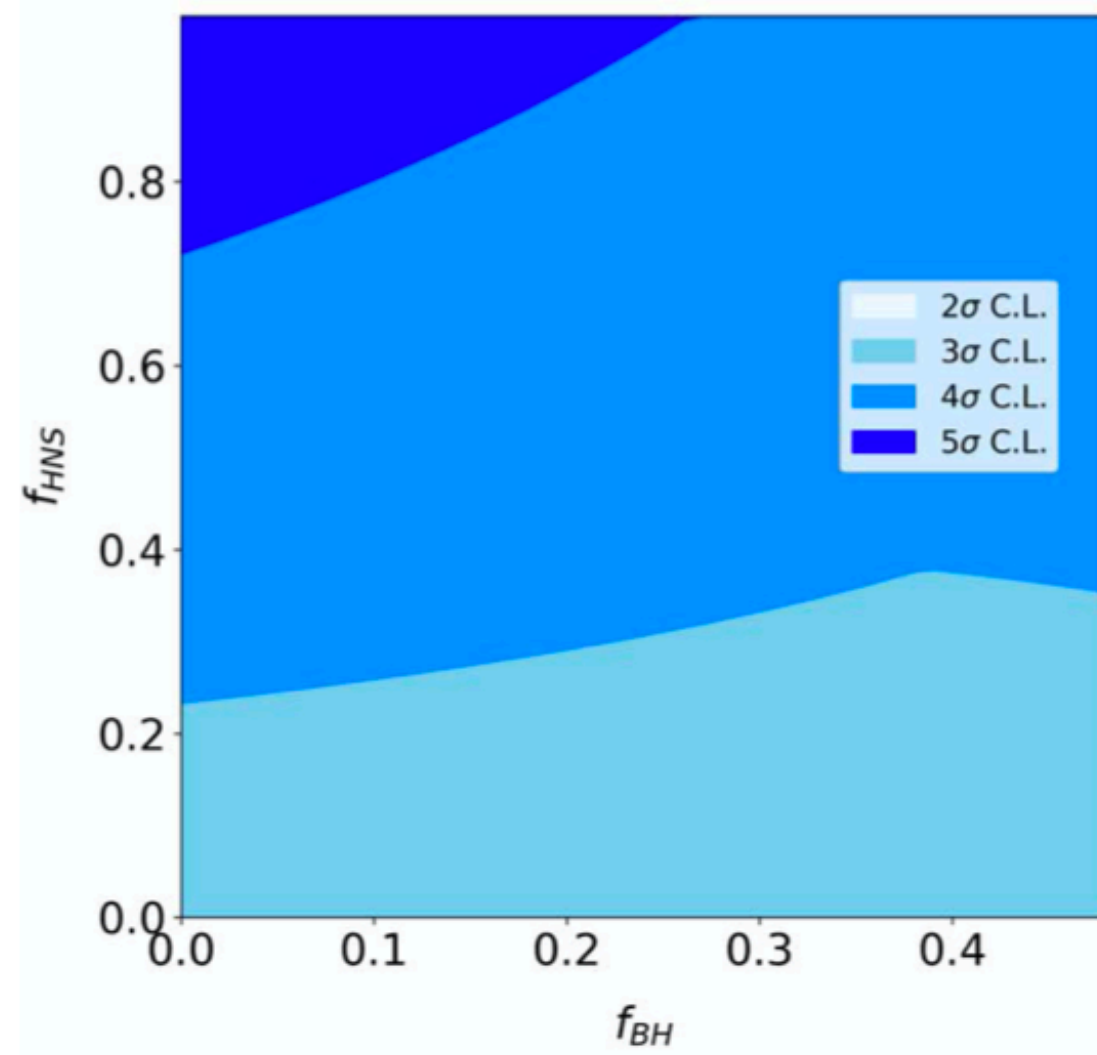
(1a) Togashi EOS, NH



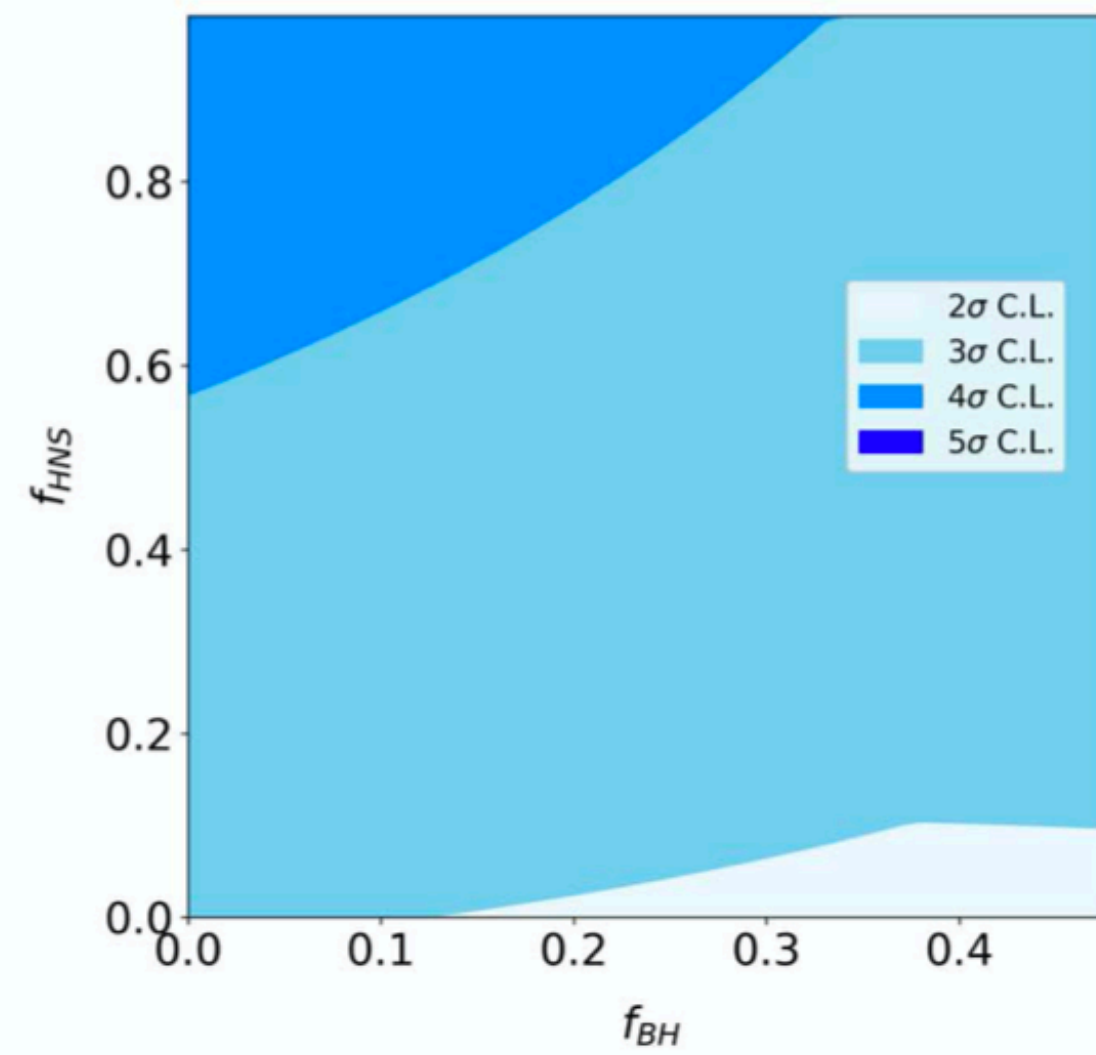
(1b) LS220 EOS, NH



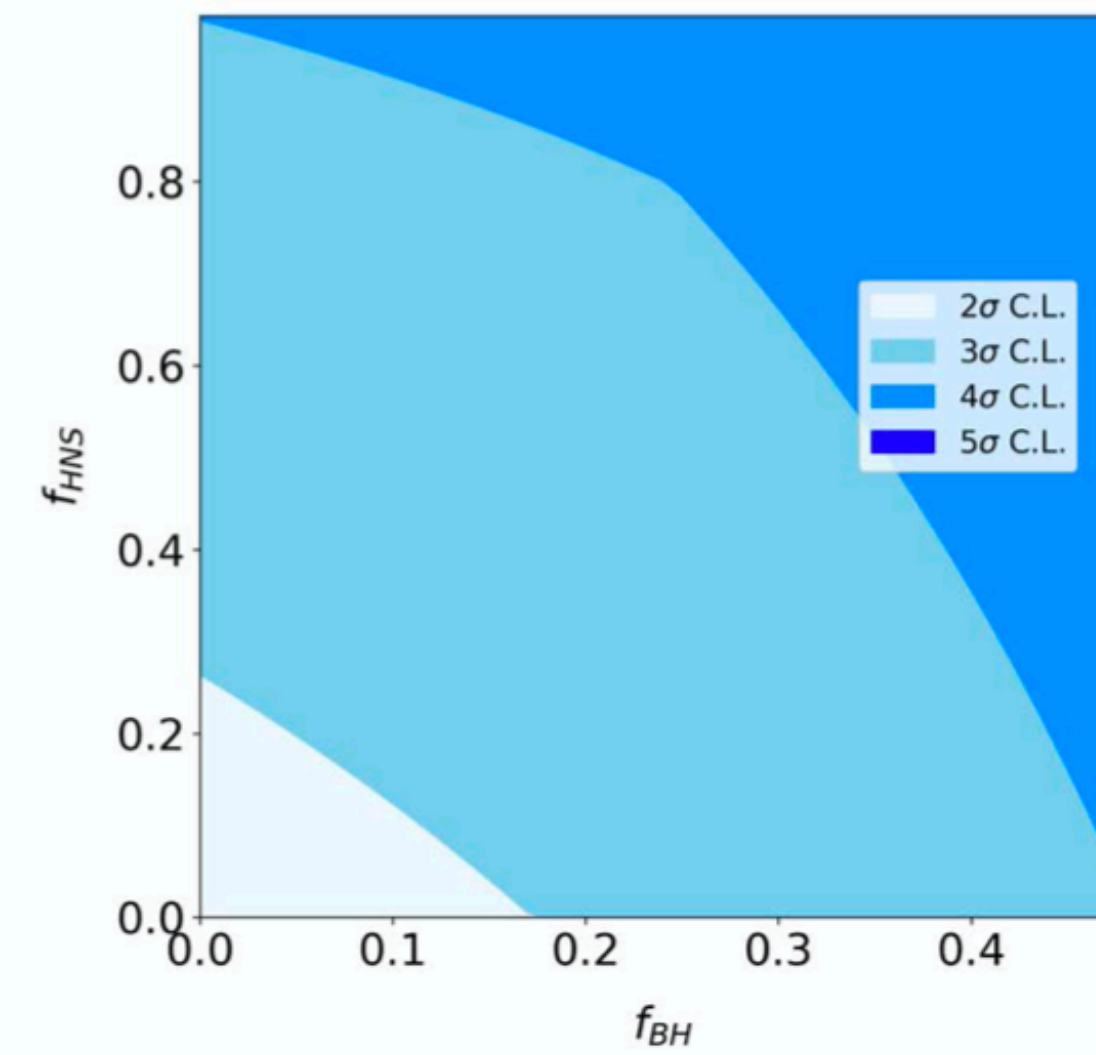
(1c) Shen EOS, NH



(2a) Togashi EOS, IH



(2b) LS220 EOS, IH



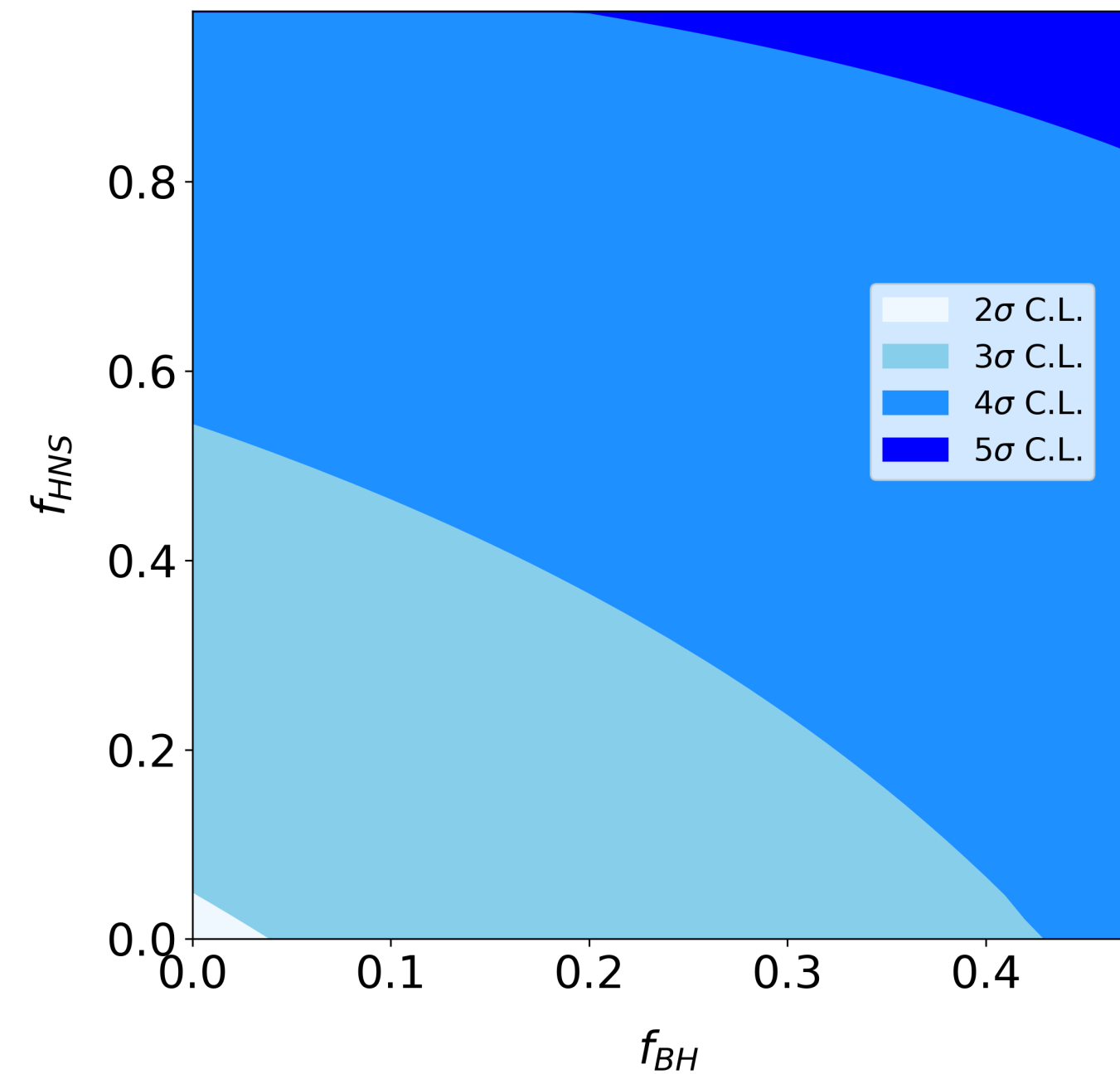
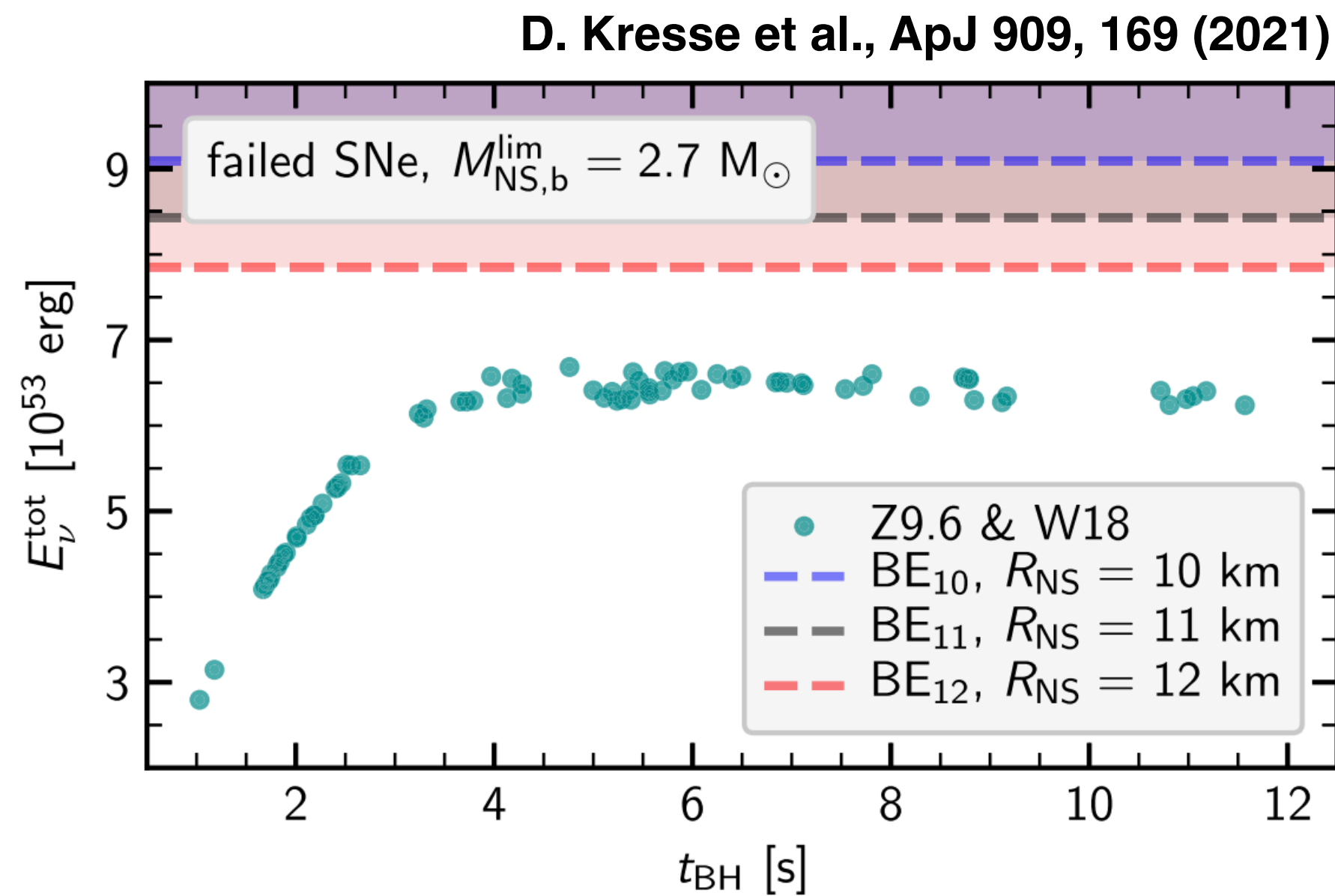
(2c) Shen EOS, IH



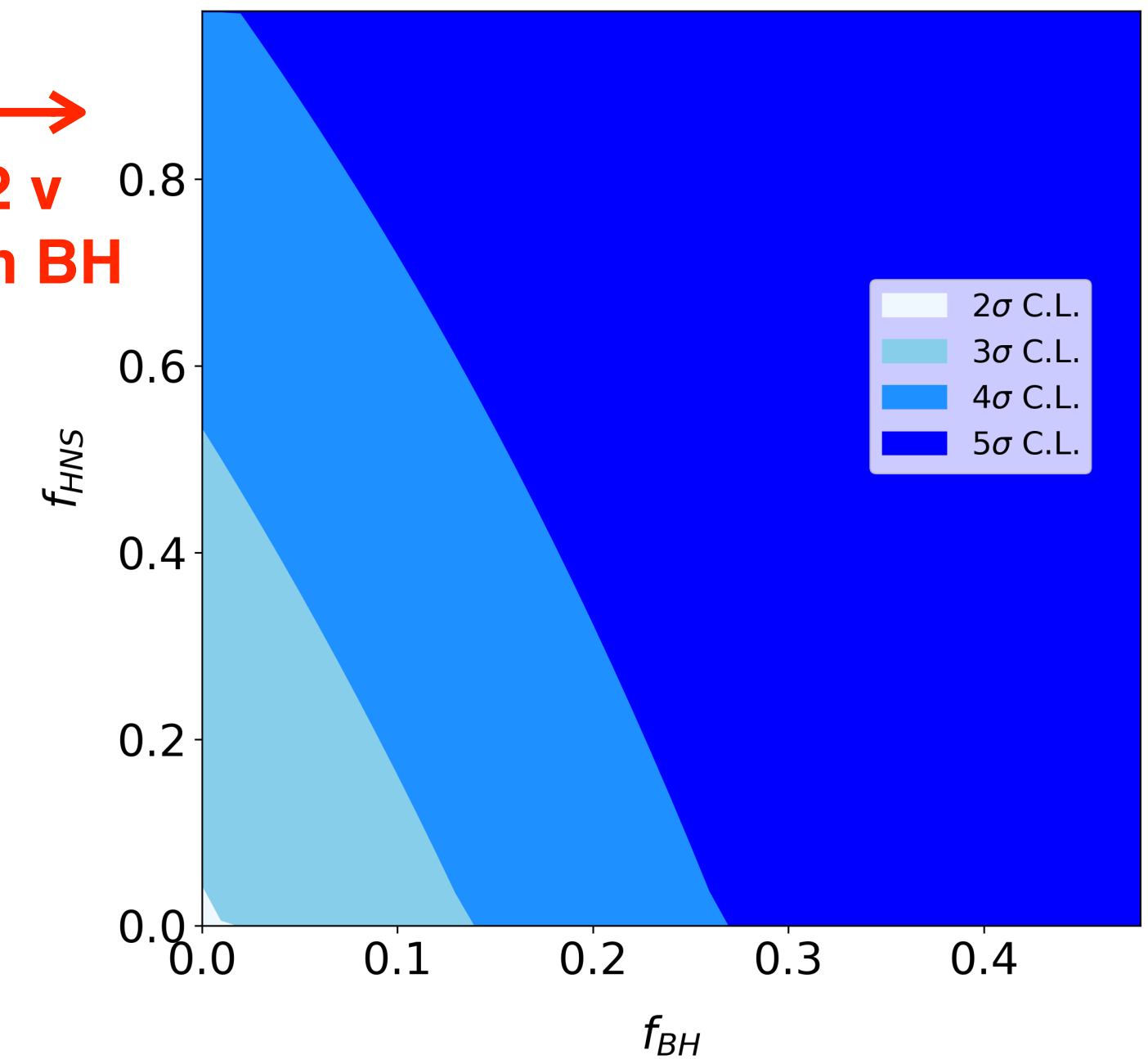
# Late Time BH Formation

- Some models predict BH formation with a significant delay ( $\gg 1$  sec), predicting up to  $\times \sim 2$  larger released neutrino energy because of more mass accretion.
- This leads to stricter constraints on  $f_{\text{BH}}$  (most obvious in Shen EOS case).

## Hyper-K 10yr (Togashi EOS)



$\rightarrow$   
 $\times 2 \nu$   
 from BH





## Diffuse Neutrino Flux Based on the Rates of Core-collapse Supernovae and Black Hole Formation Deduced from a Novel Galactic Chemical Evolution Model

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Received 2023 May 22; revised 2023 July 1; accepted 2023 July 2; published 2023 August 14

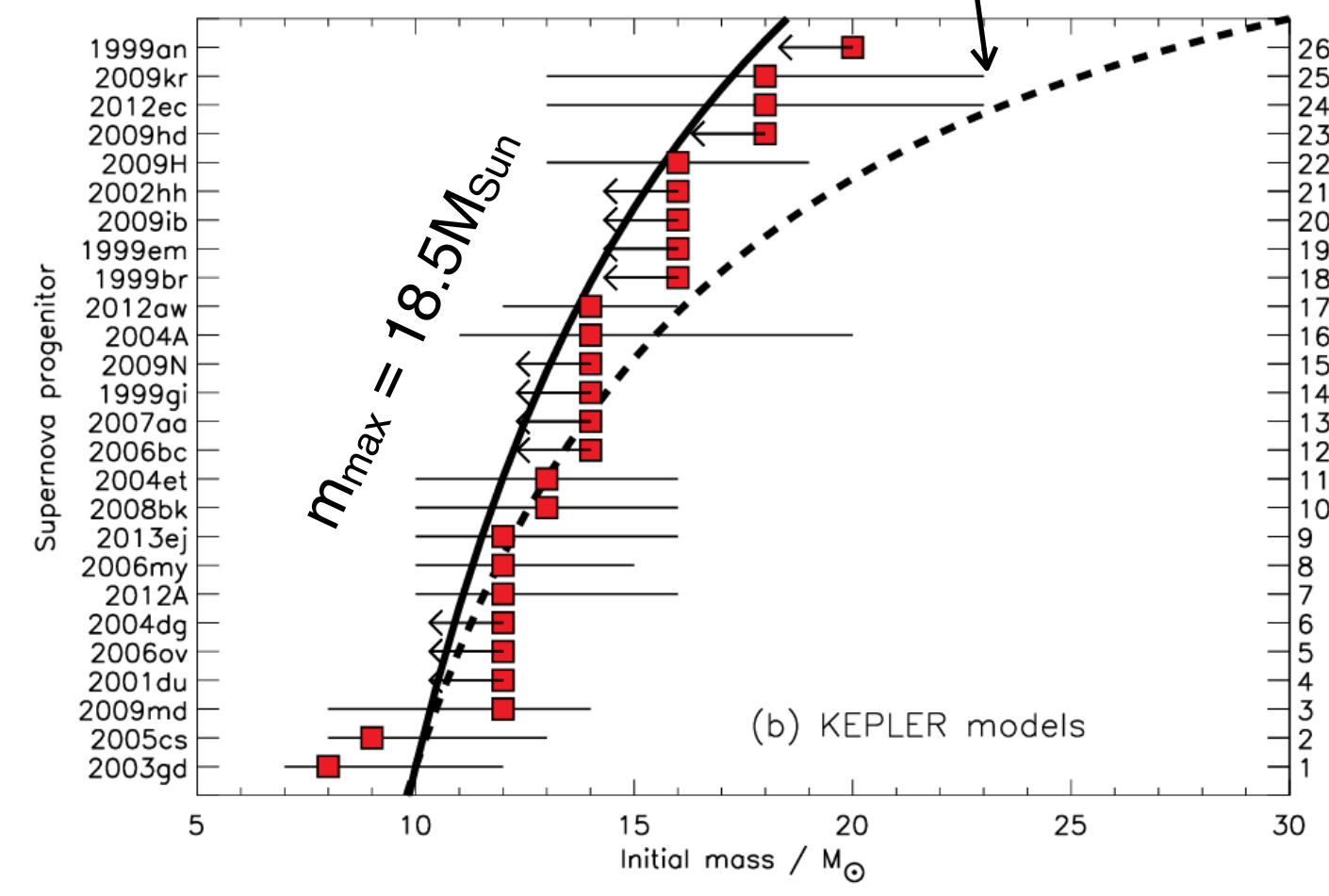
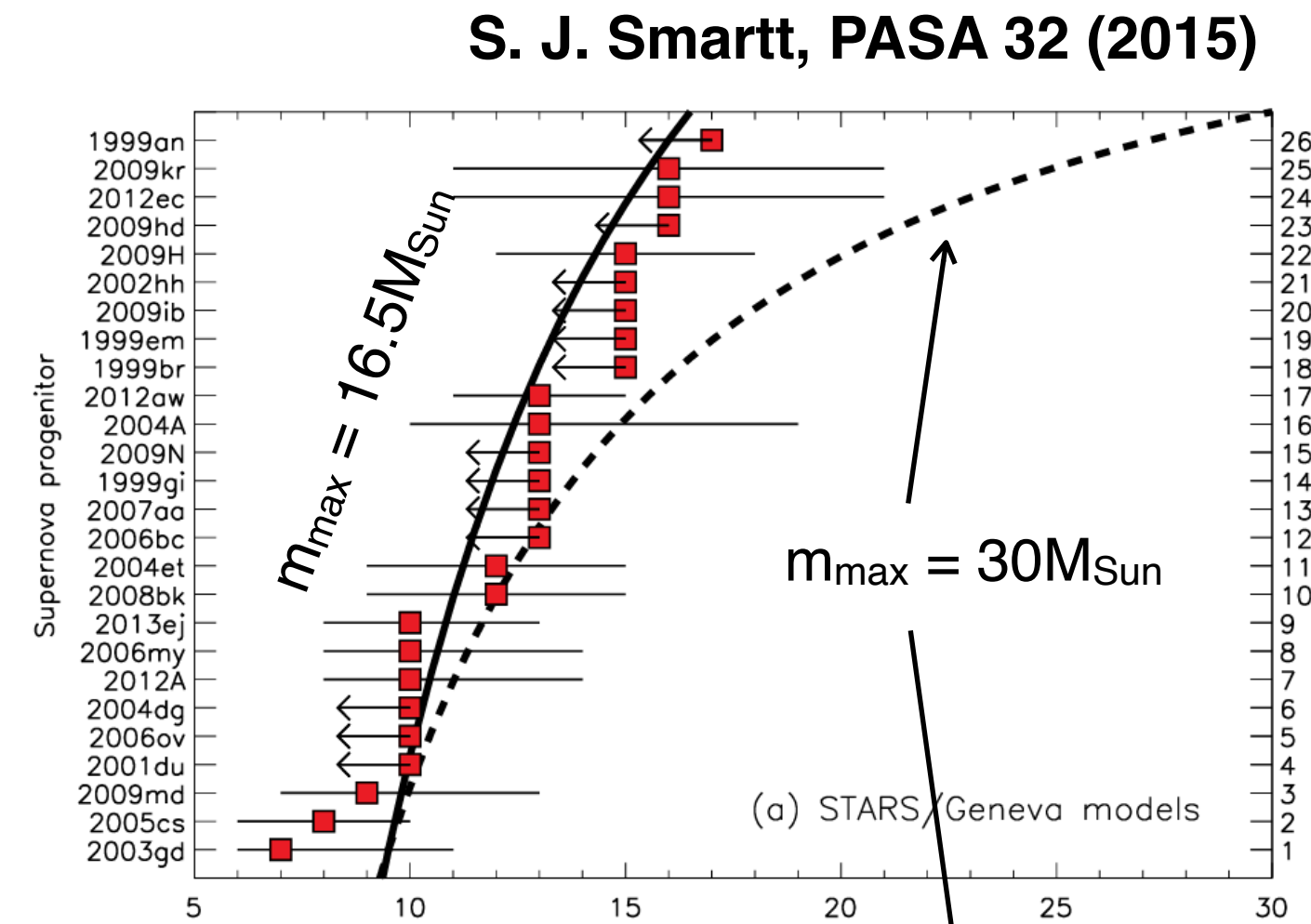
### Abstract

Fluxes of the diffuse supernova neutrino background (DSNB) are calculated based on a new modeling of galactic chemical evolution, where a variable stellar initial mass function (IMF), depending on the galaxy type, is introduced and black hole (BH) formation from the failed supernova is considered for progenitors heavier than  $18M_{\odot}$ . The flux calculations are performed for different combinations of the star formation rate, nuclear equation of state, and neutrino mass hierarchy, to examine the systematic effects from these factors. In any case, our new model predicts the enhanced DSNB  $\bar{\nu}_e$  flux at  $E_{\nu} \gtrsim 30$  MeV and  $E_{\nu} \lesssim 10$  MeV, due to more frequent BH formation and a larger core-collapse rate at high redshifts in early-type galaxies, respectively. Event rate spectra of the DSNB  $\bar{\nu}_e$  at a detector from the new model are shown, and the detectability at water-based Cherenkov detectors, Super-Kamiokande with a gadolinium dissolution and Hyper-Kamiokande, is discussed. In order to investigate the impacts of the assumptions in the new model, we prepare alternative models, based on different IMF forms and treatments of BH formation, and estimate the discrimination capabilities between the new and alternative models at these detectors.

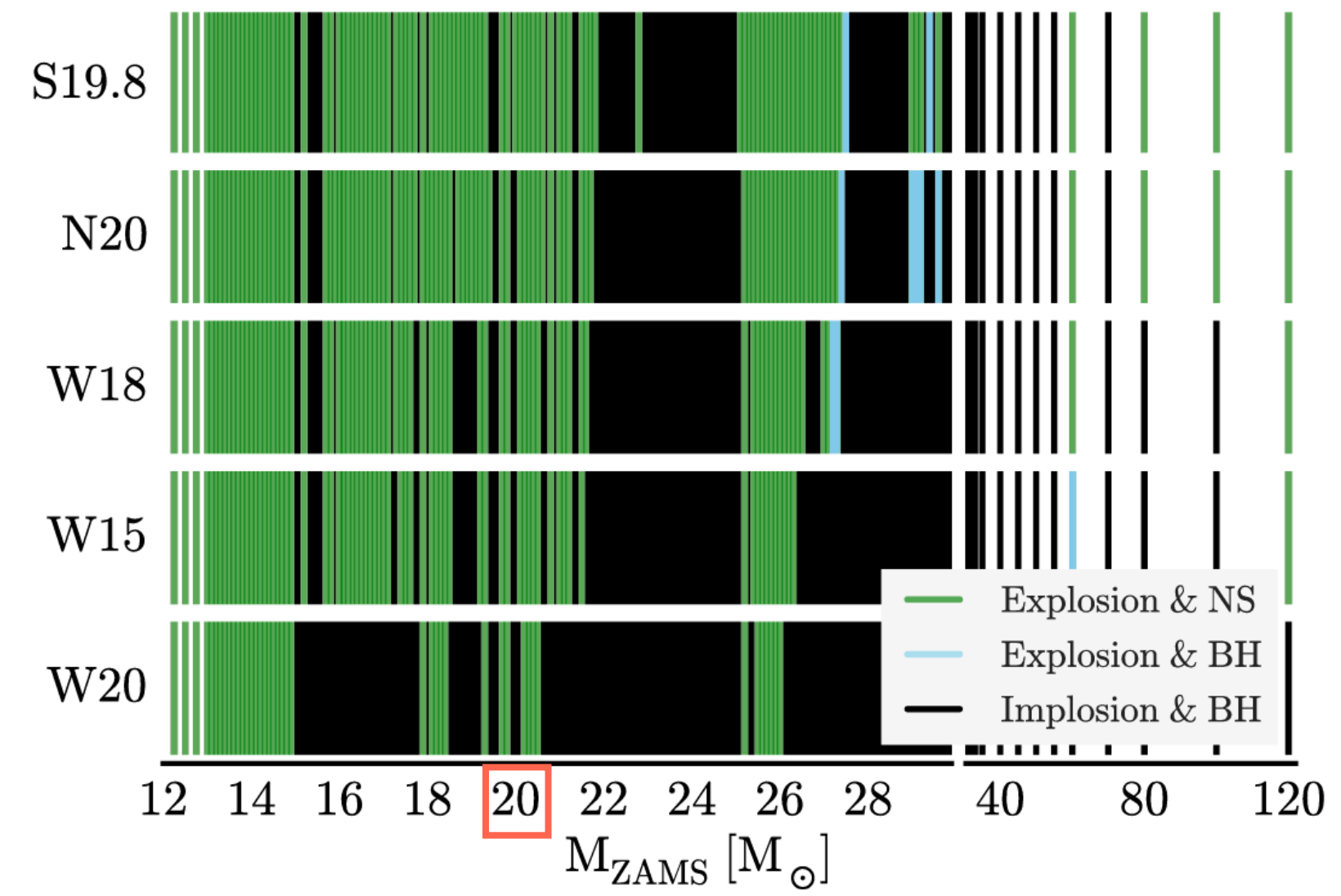
*Unified Astronomy Thesaurus concepts:* [Neutrino astronomy \(1100\)](#); [Supernova neutrinos \(1666\)](#); [Core-collapse supernovae \(304\)](#); [Massive stars \(732\)](#); [Neutron stars \(1108\)](#); [Black holes \(162\)](#); [Galaxy chemical evolution \(580\)](#); [Star formation \(1569\)](#); [Initial mass function \(796\)](#)

# Constructing New Framework: CCSN Mass Limit

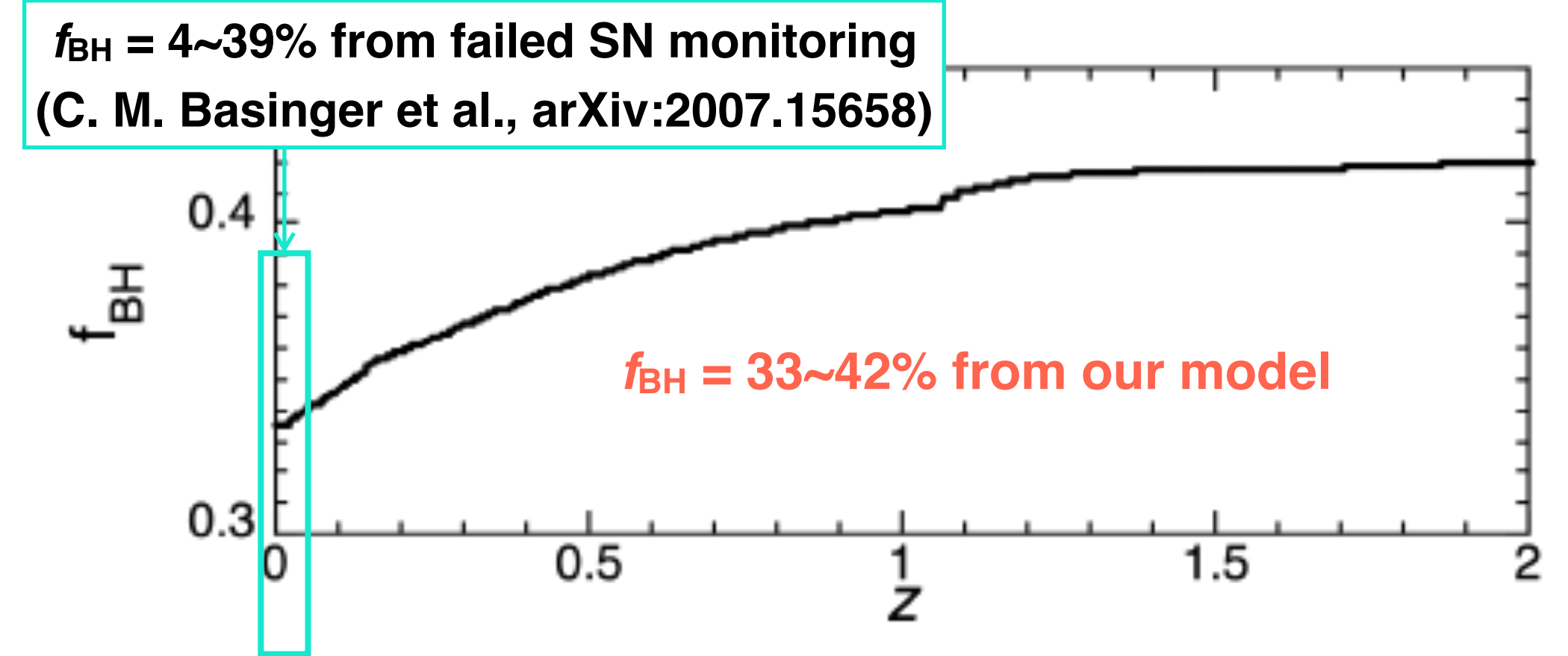
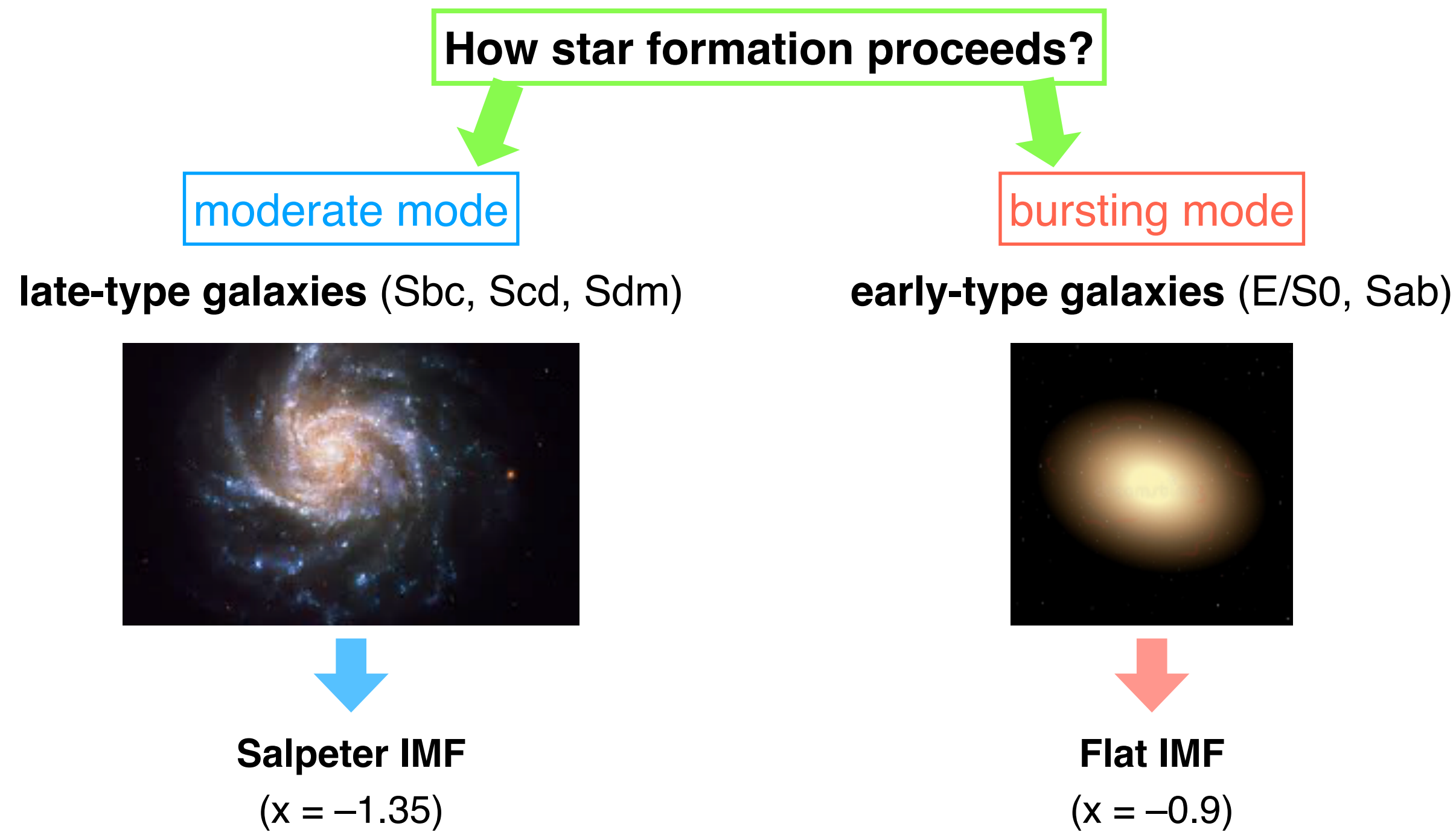
- We set the maximum mass of progenitors for successful explosions to  $18M_{\text{Sun}}$ .
  - Observationally,  $m_{\text{min}} \sim 8M_{\text{Sun}}$  and  $m_{\text{max}} \sim 18M_{\text{Sun}}$  are supported.
  - There is a theoretical work that implies failed SNe above  $\sim 20M_{\text{Sun}}$ .
- Many galactic chemical evolution schemes adopt a high  $m_{\text{max}}$  ( $50\sim 100M_{\text{Sun}}$ ).
  - Our  $m_{\text{max}} = 18M_{\text{Sun}}$  assumption reduces the number of CCSNe to  $\sim 70\%$ .
  - Accordingly, the total amount of heavy elements is reduced to  $\sim 50\%$ .



T. Sukhbold et al., ApJ 821, 38 (2016)



- In order to achieve consistency with observed chemical abundance, we propose a new evolution model.
- We categorize galaxies into five and assume different initial mass functions (IMF) depending types.
- The fraction for BH formation from this model is **33~42%** (higher rate than many other DSNB models).

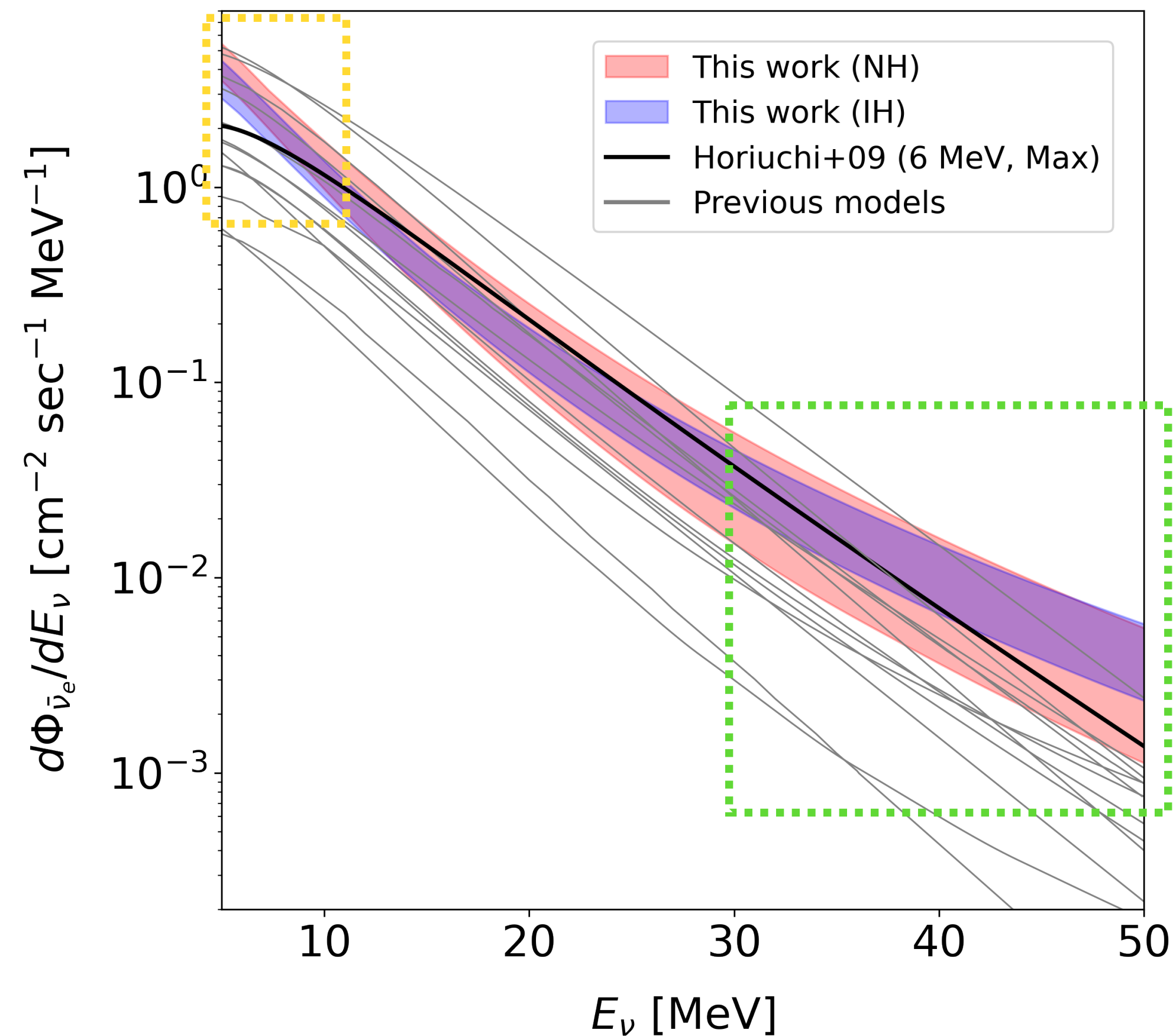


Our model nomenclature

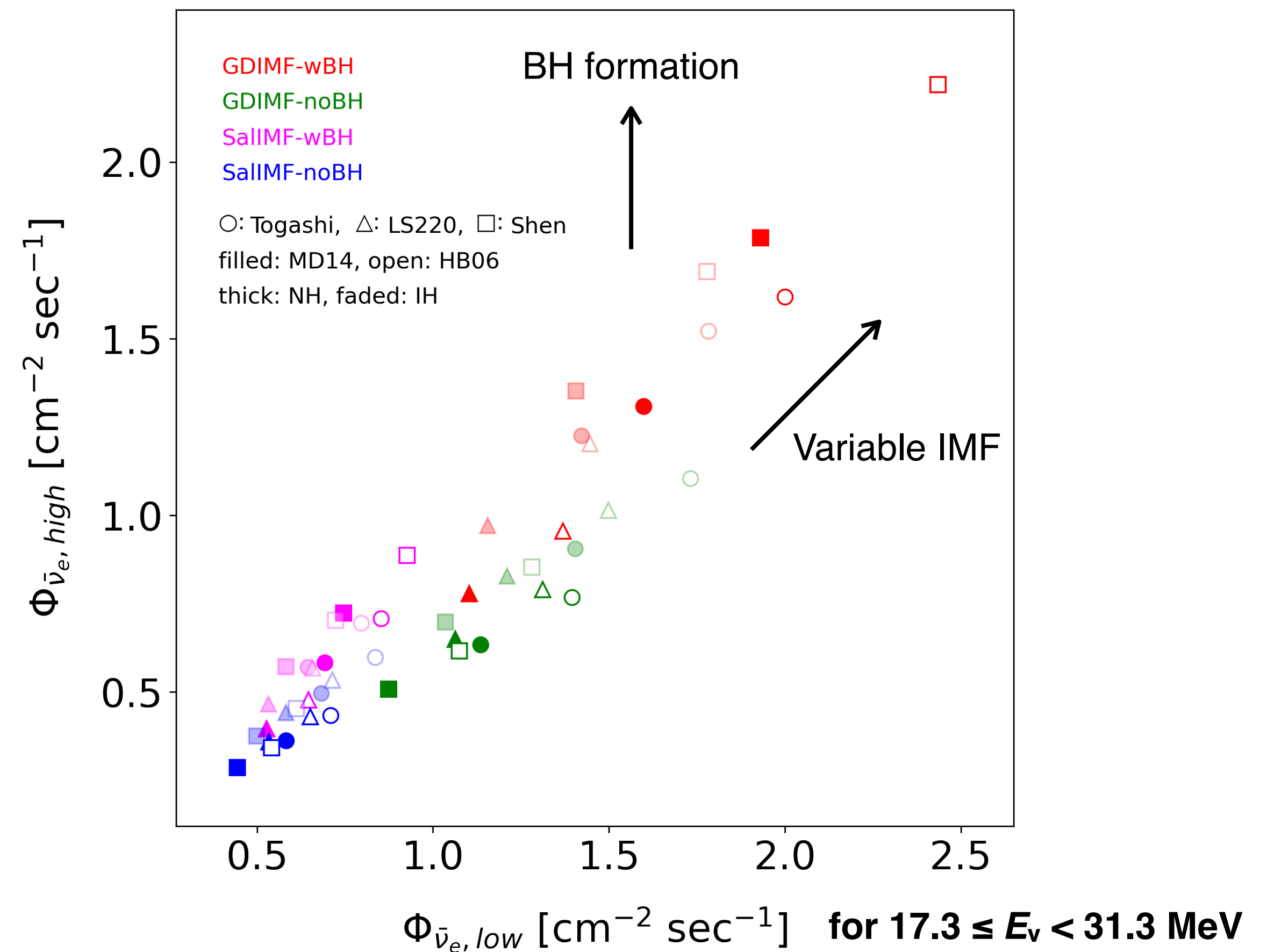
Name	IMF form	BH treatment
GDIMF-wBH (ref.)	Variable	BHs for $18-100M_{\odot}$
GDIMF-noBH	Variable	No BH
SalIMF-wBH	Salpeter	BHs for $18-100M_{\odot}$
SalIMF-noBH	Salpeter	No BH

# Calculated DSNB Flux

- Our model shows DSNB flux enhancements at high and low energies.
- High energy ( $>30$  MeV): Large contribution from BH formation.
- Low energy ( $<10$  MeV): Redshifted neutrinos from early-type galaxies with large CCSN rates.

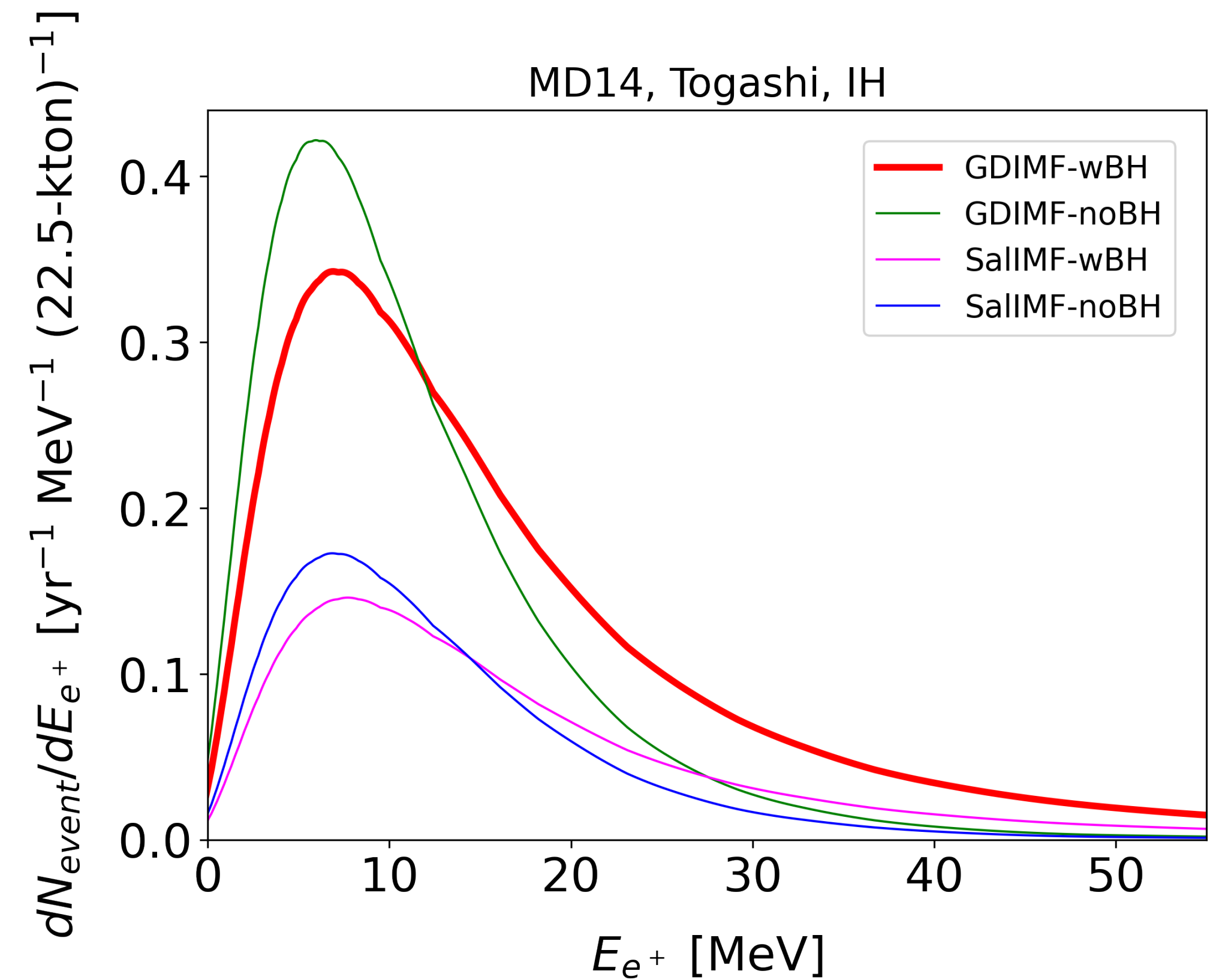
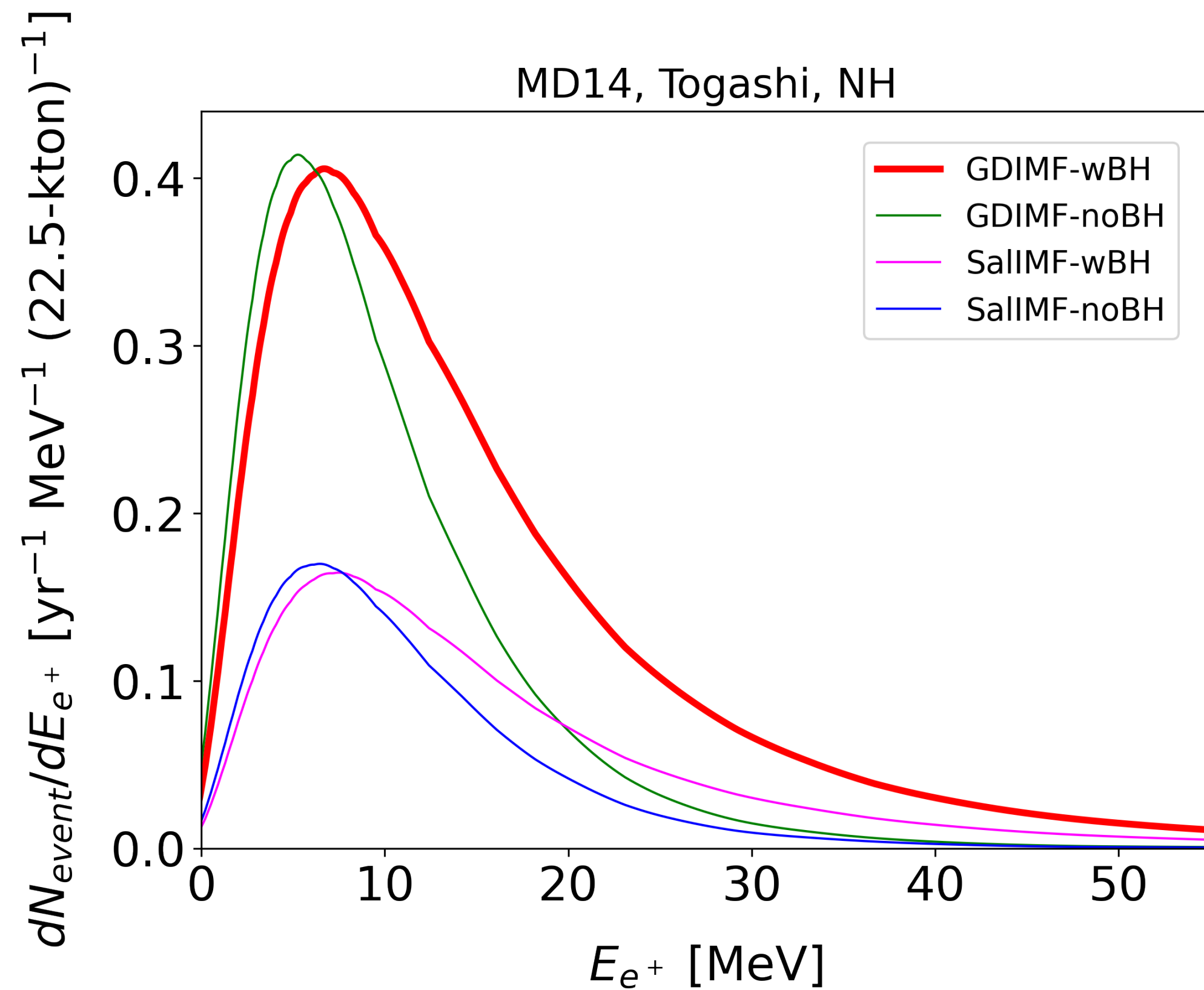


for  $13.3 \leq E_\nu < 17.3$  MeV



# Event Rate

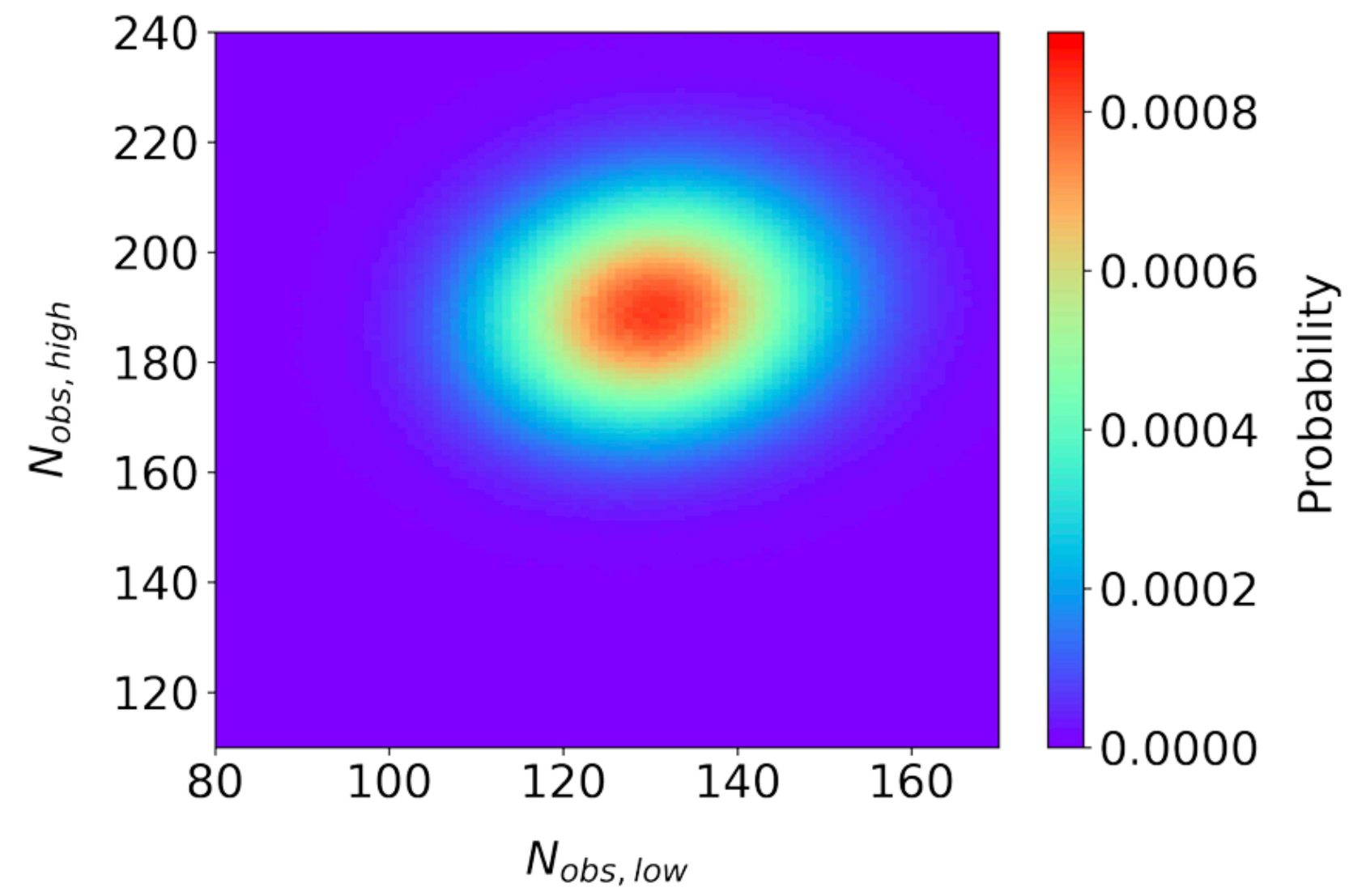
- Event rate spectrum at a water volume is calculated with IBD cross section.
- SK-Gd w/ 70% ntag efficiency over 10yr = 15~18 signals
- Hyper-K w/ Super-K ntag efficiency over 10yr = 50~60 signals



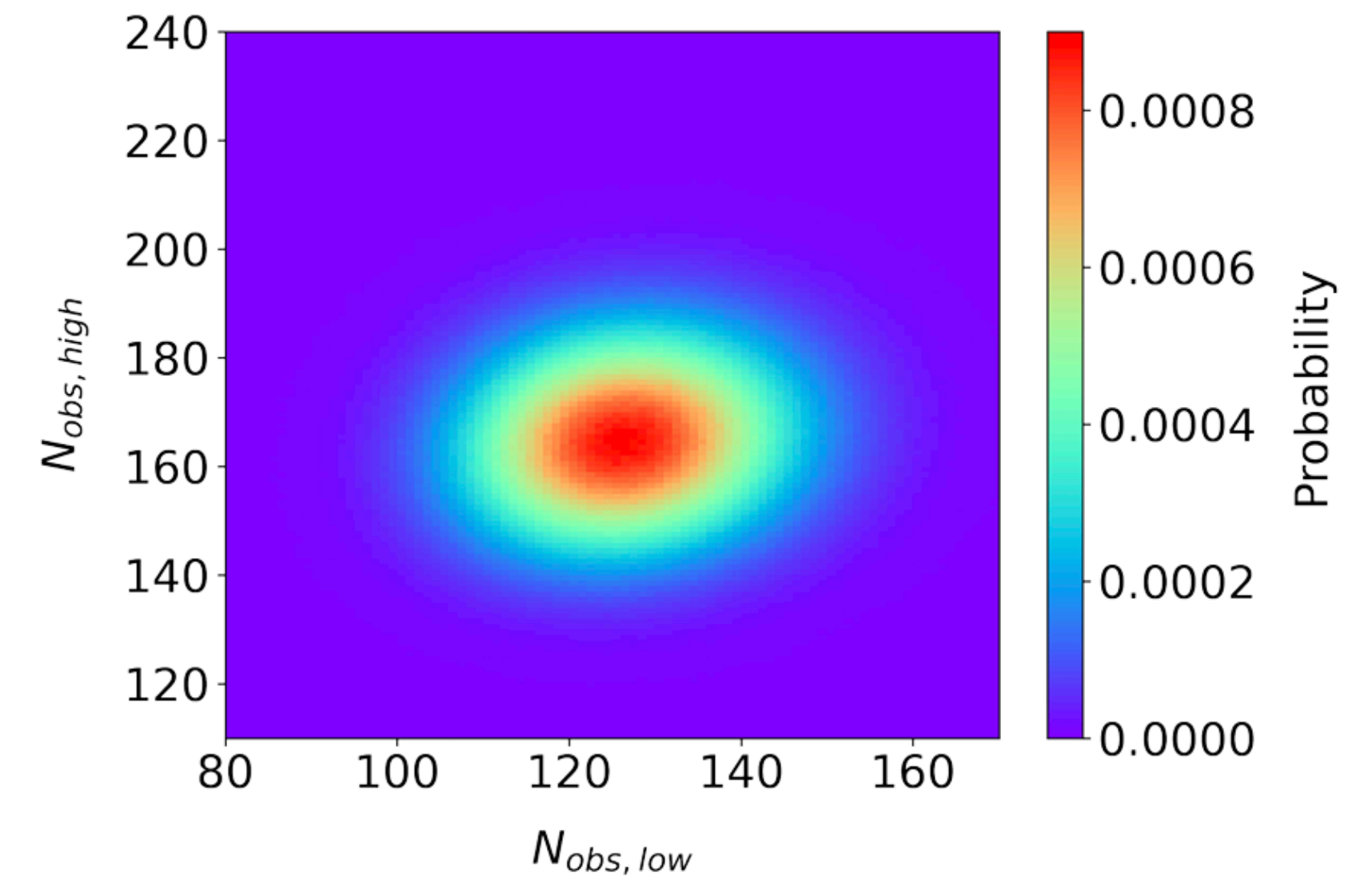
## Bayes' theorem:

$$\overset{\text{Posterior}}{P(\text{model}|\text{data})} = \frac{\overset{\text{Likelihood (frequentist)}}{P(\text{data}|\text{model})} \times \overset{\text{Prior}}{P(\text{model})}}{\underbrace{\sum_{\text{model}} P(\text{data}|\text{model}) \times P(\text{model})}_{\text{Normalization}}}$$

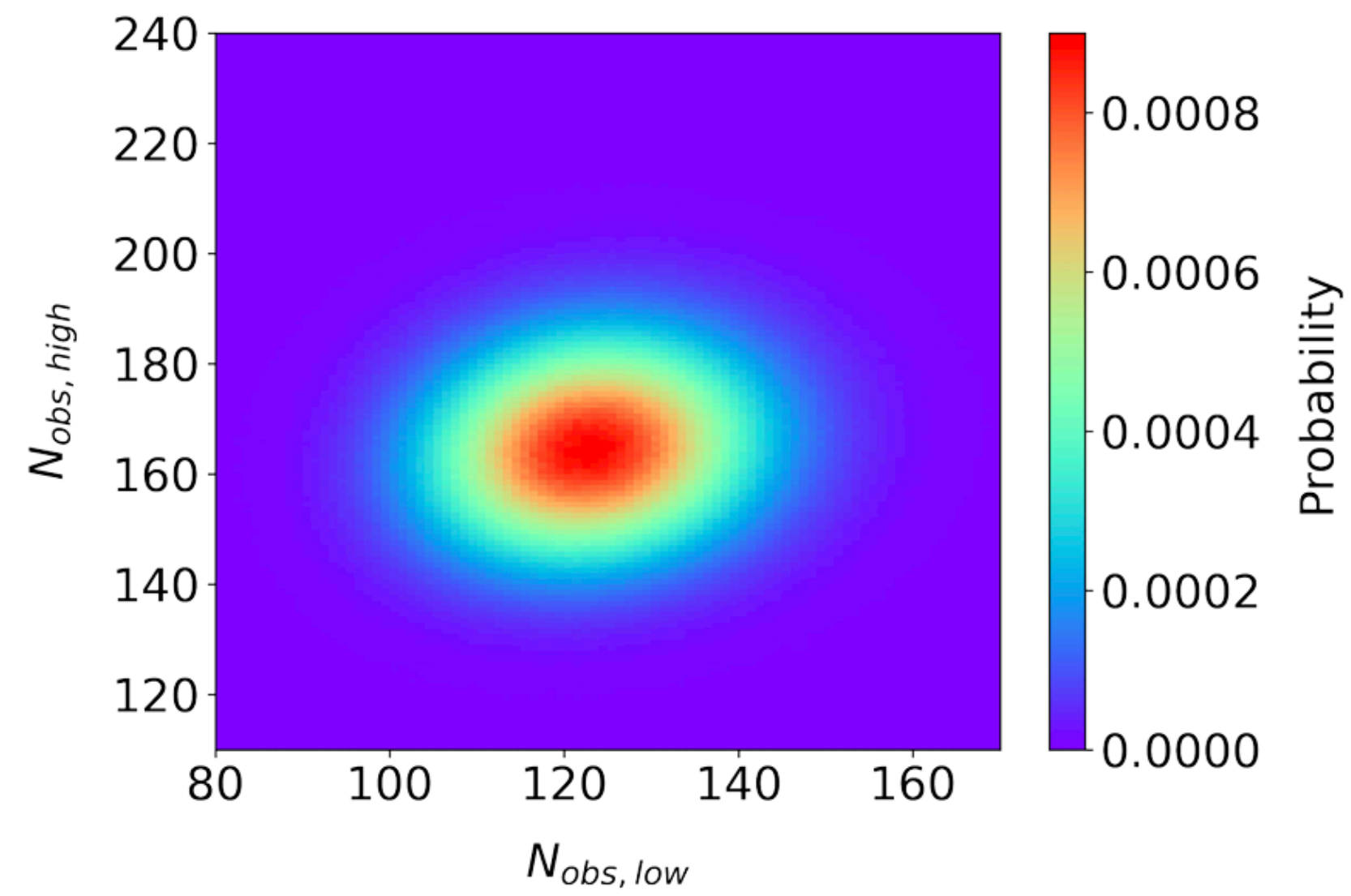
- Prior:  $P(\text{model}) = 1/N$  for testing  $N$  models.
- Likelihood: prepared based on signal+background expectations for each model.
- Utilize two energy ranges (13.3–17.3 MeV & 17.3–31.3 MeV).
- Use real data for SK-IV, and nominal expectations for SK-Gd and Hyper-K.



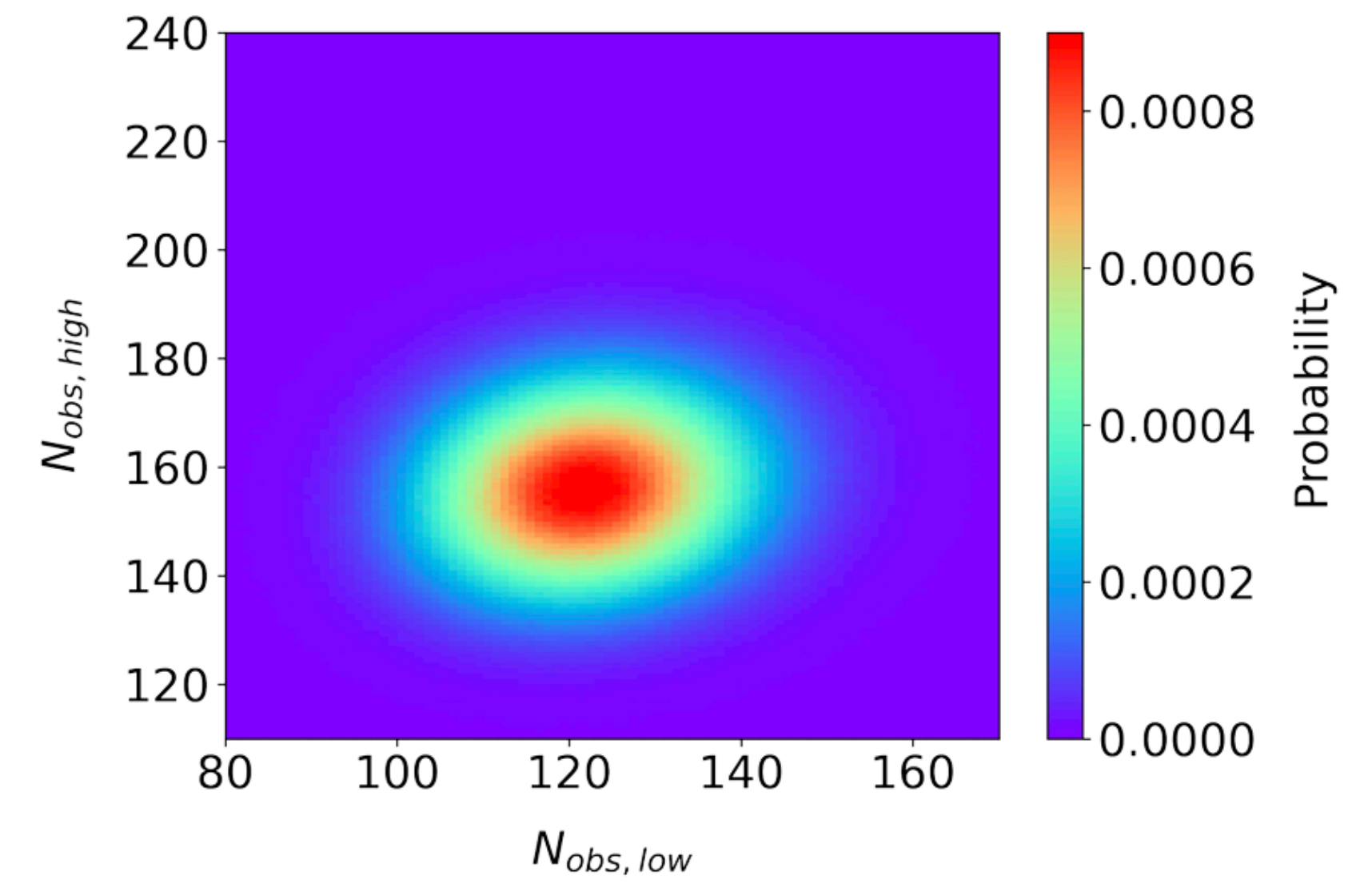
(1) GDIMF-wBH



(2) GDIMF-noBH



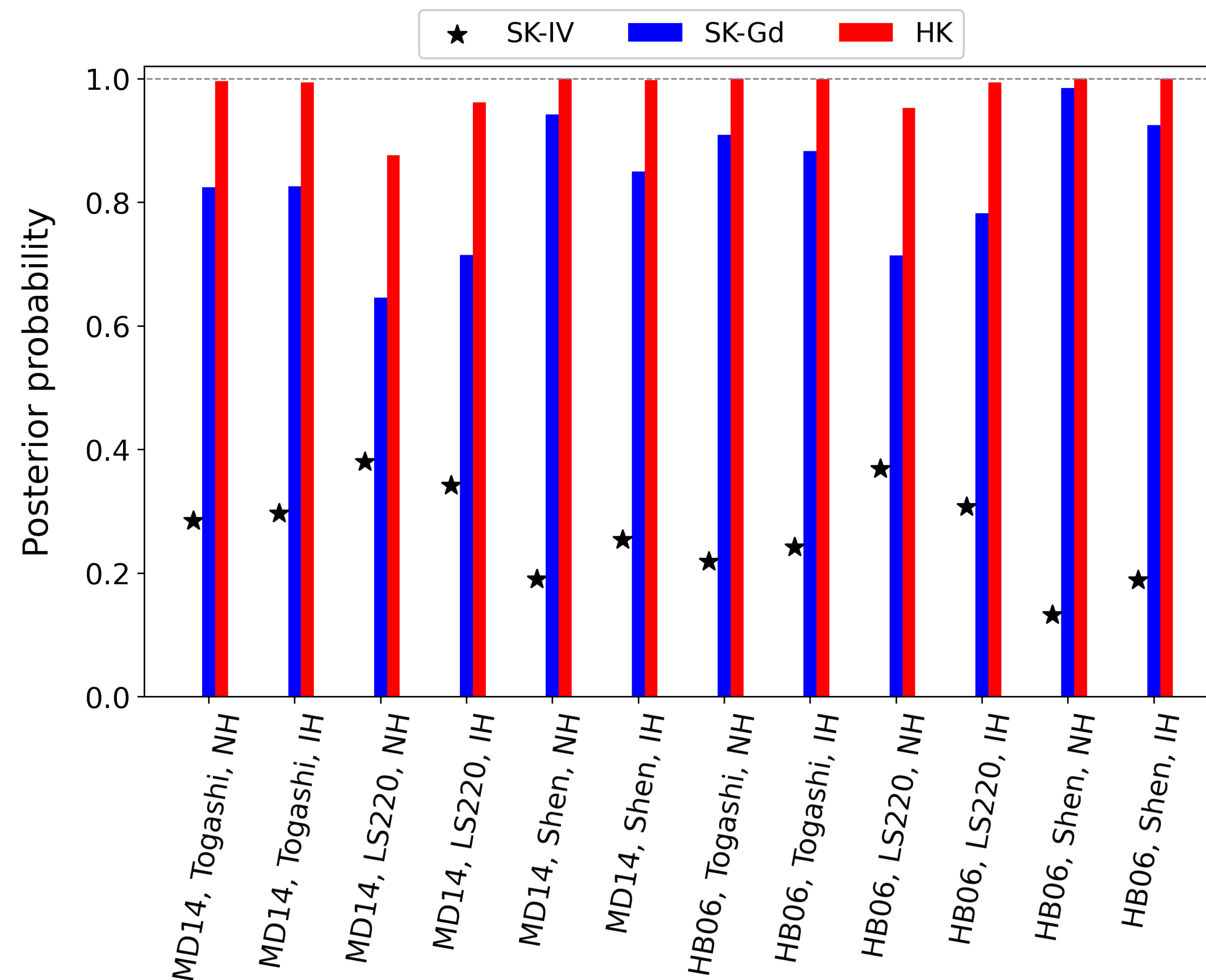
(3) SalIMF-wBH



(4) SalIMF-noBH

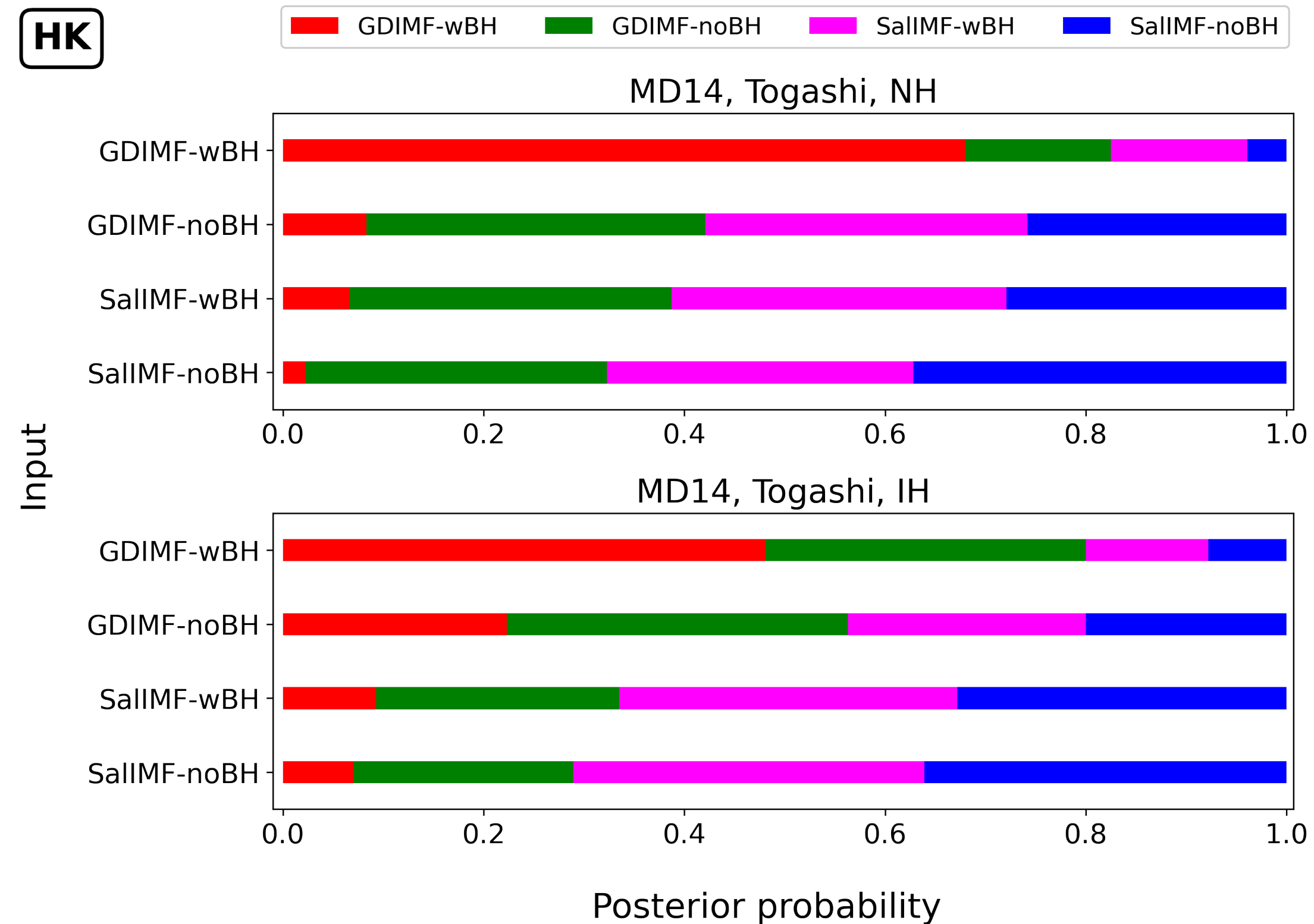


- Our reference model is tested against background only at different detectors based on Bayes' theorem.
- Background model used for likelihood is based on extrapolation from Super-K analysis.
- In most cases, our signal models can be detected well over background.

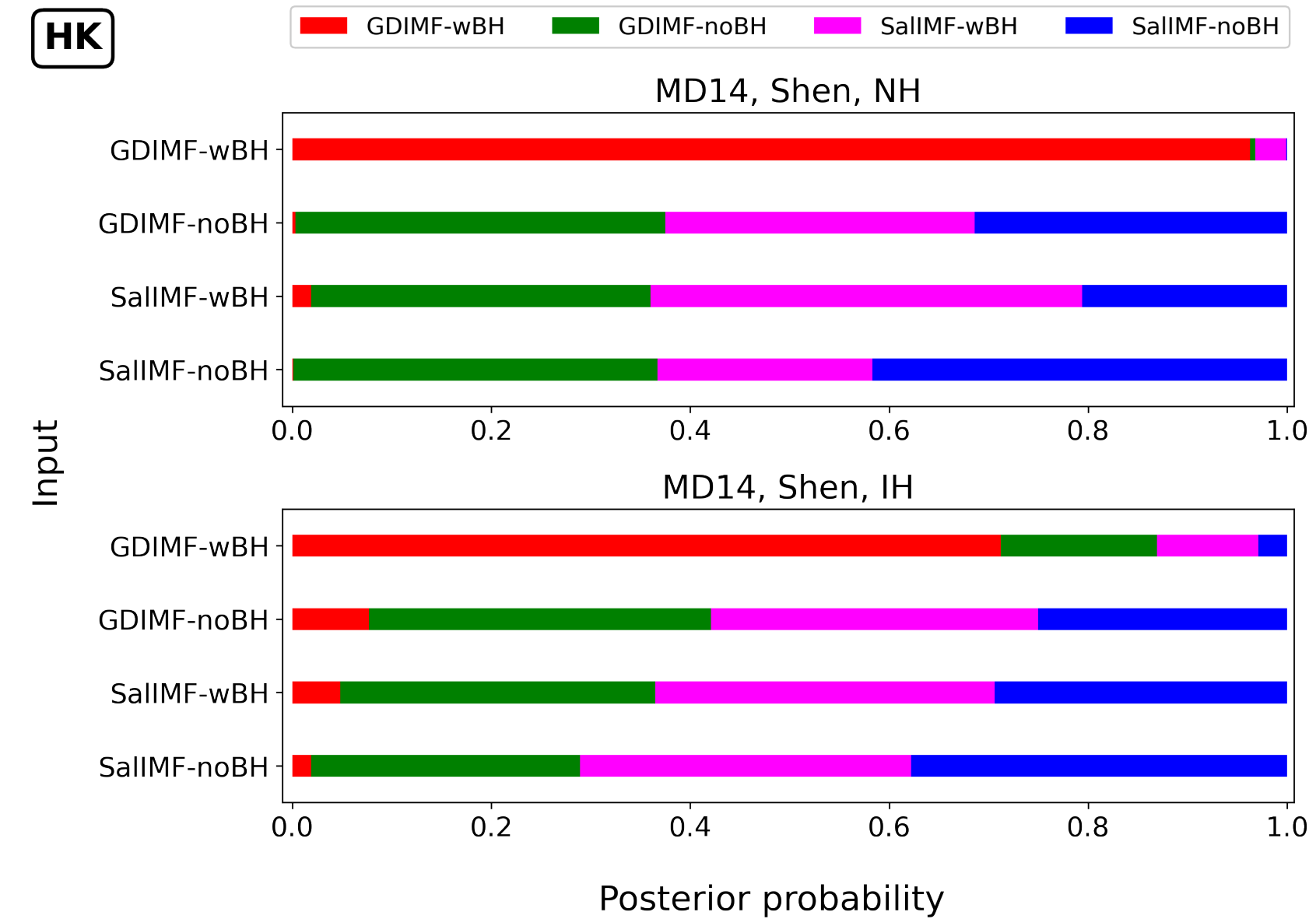
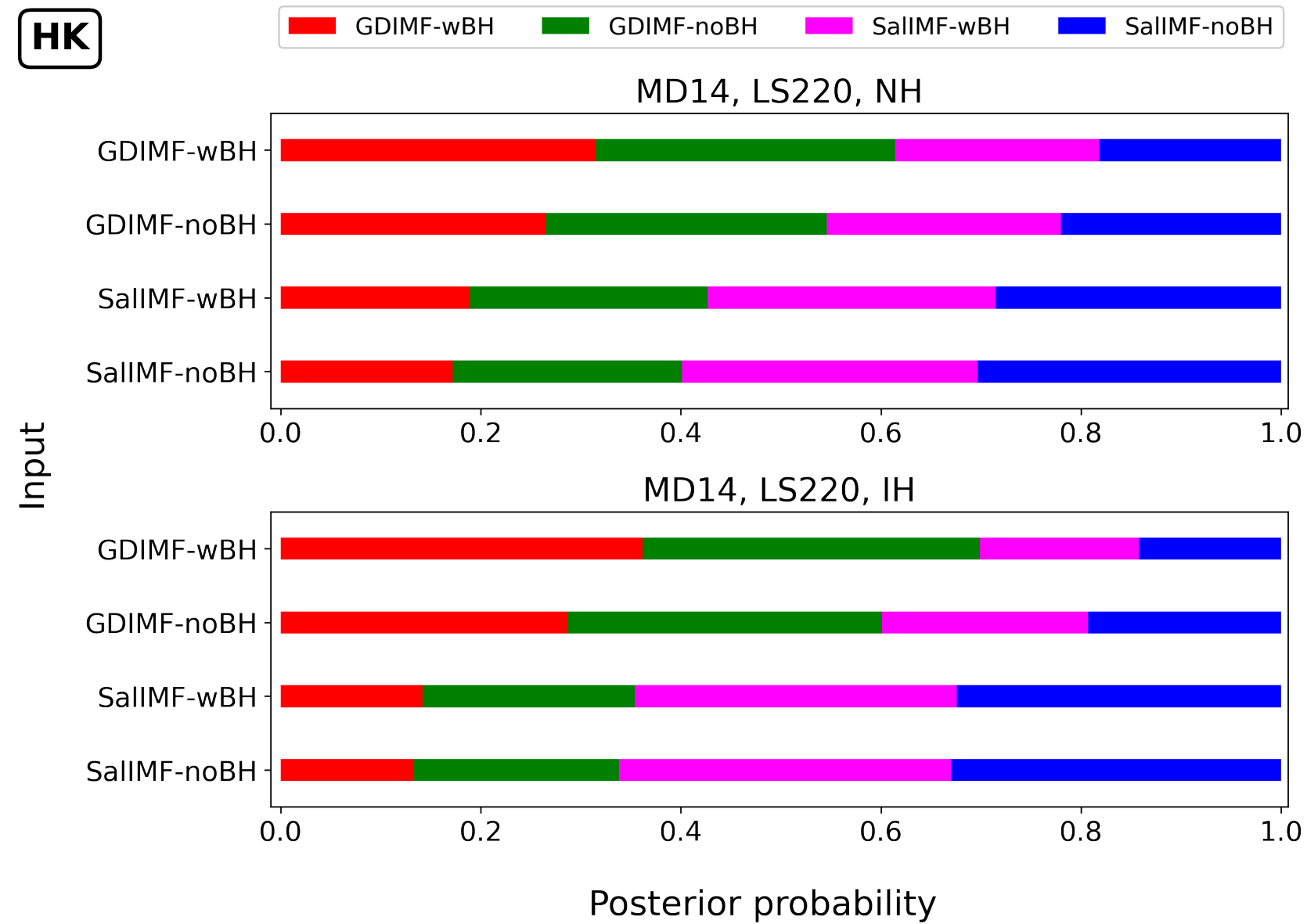
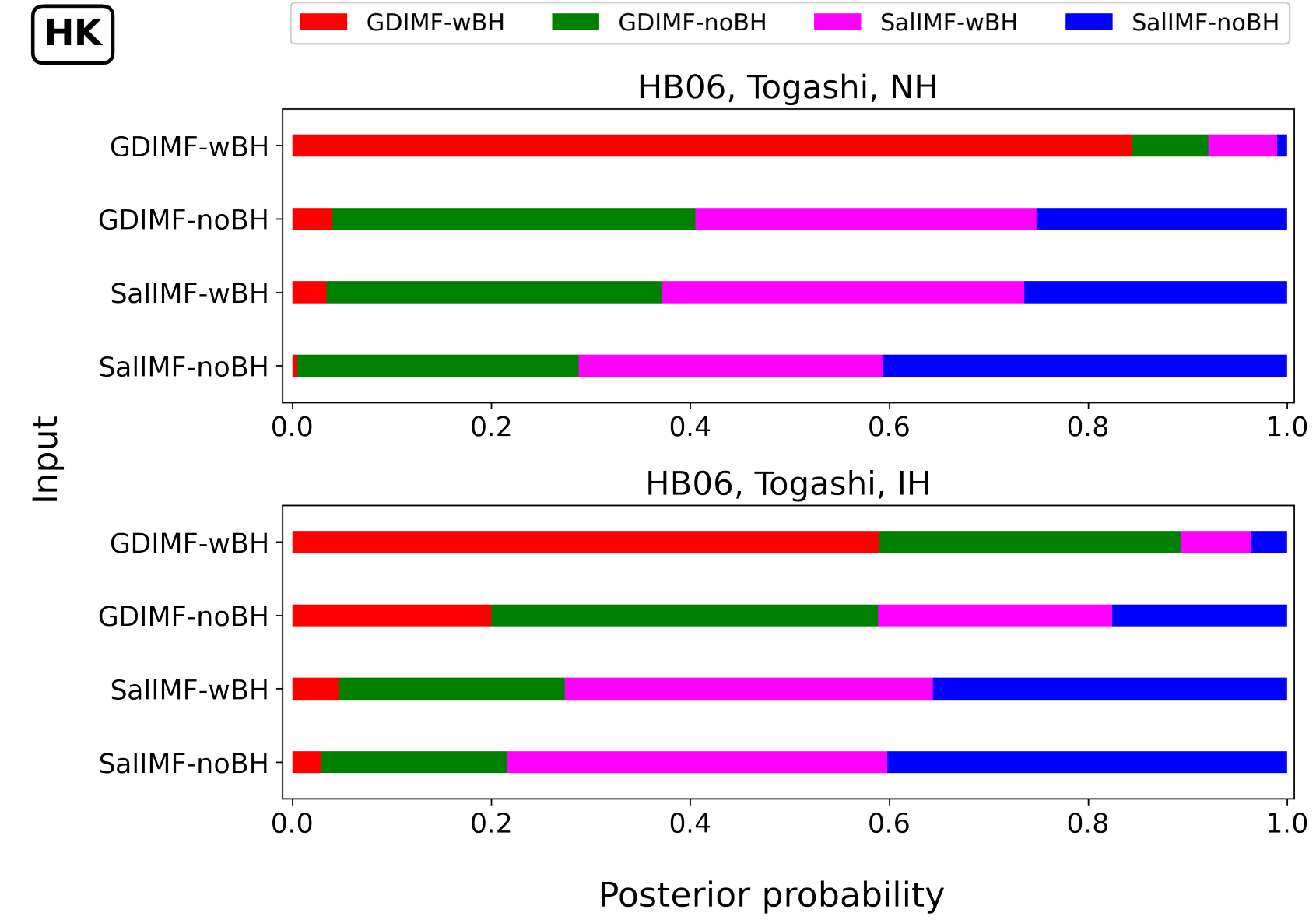
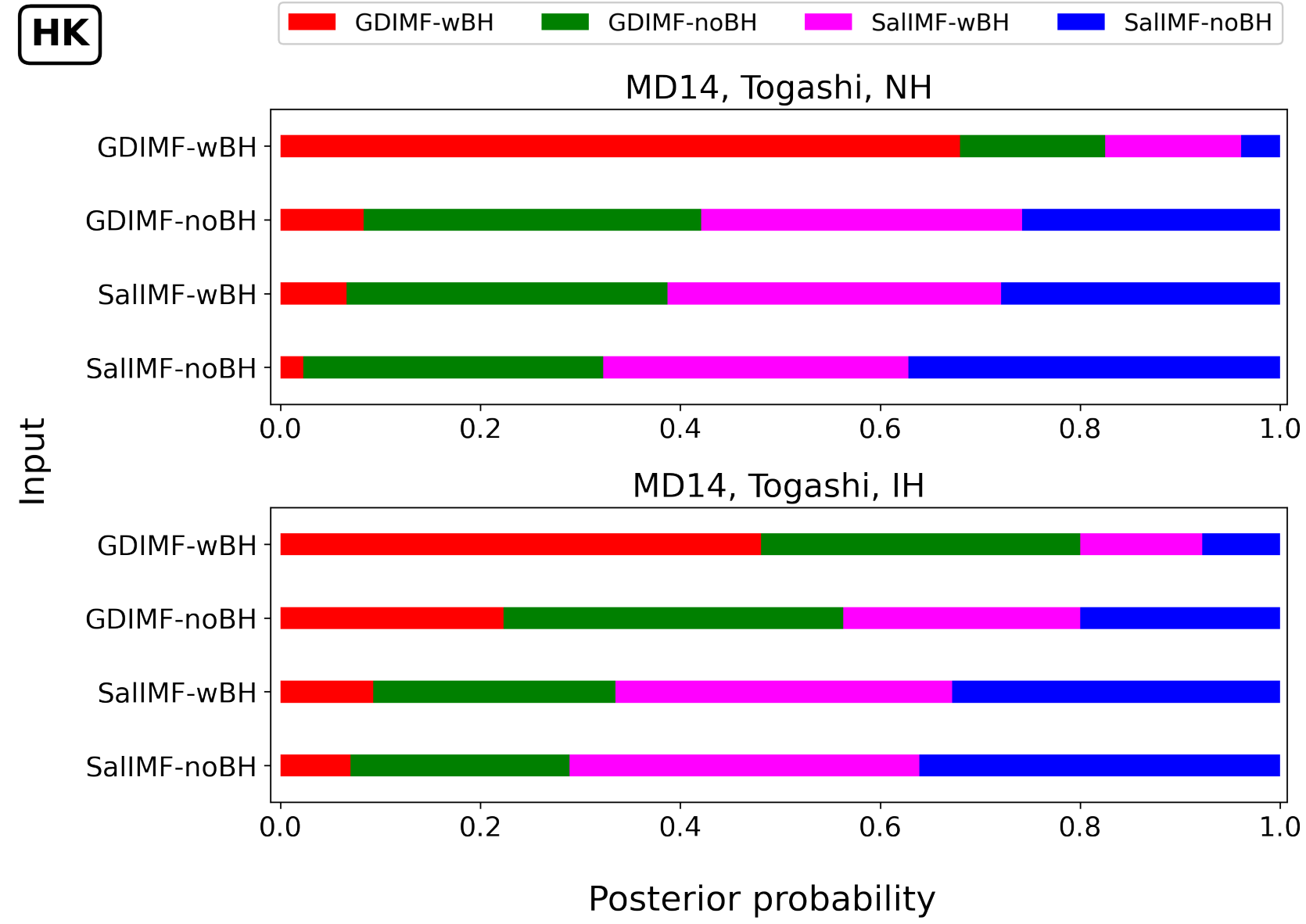


# Model Discrimination

- Our reference model is tested against alternative models with different assumptions on IMF and BH.
- In NH case, the reference model is well discriminated from others.
- In IH case, only IMF assumption can be tested well.
- The results differ for other choices of SFR and nuclear EOS (all results are discussed in the paper).



# EOS & SFR Dependences

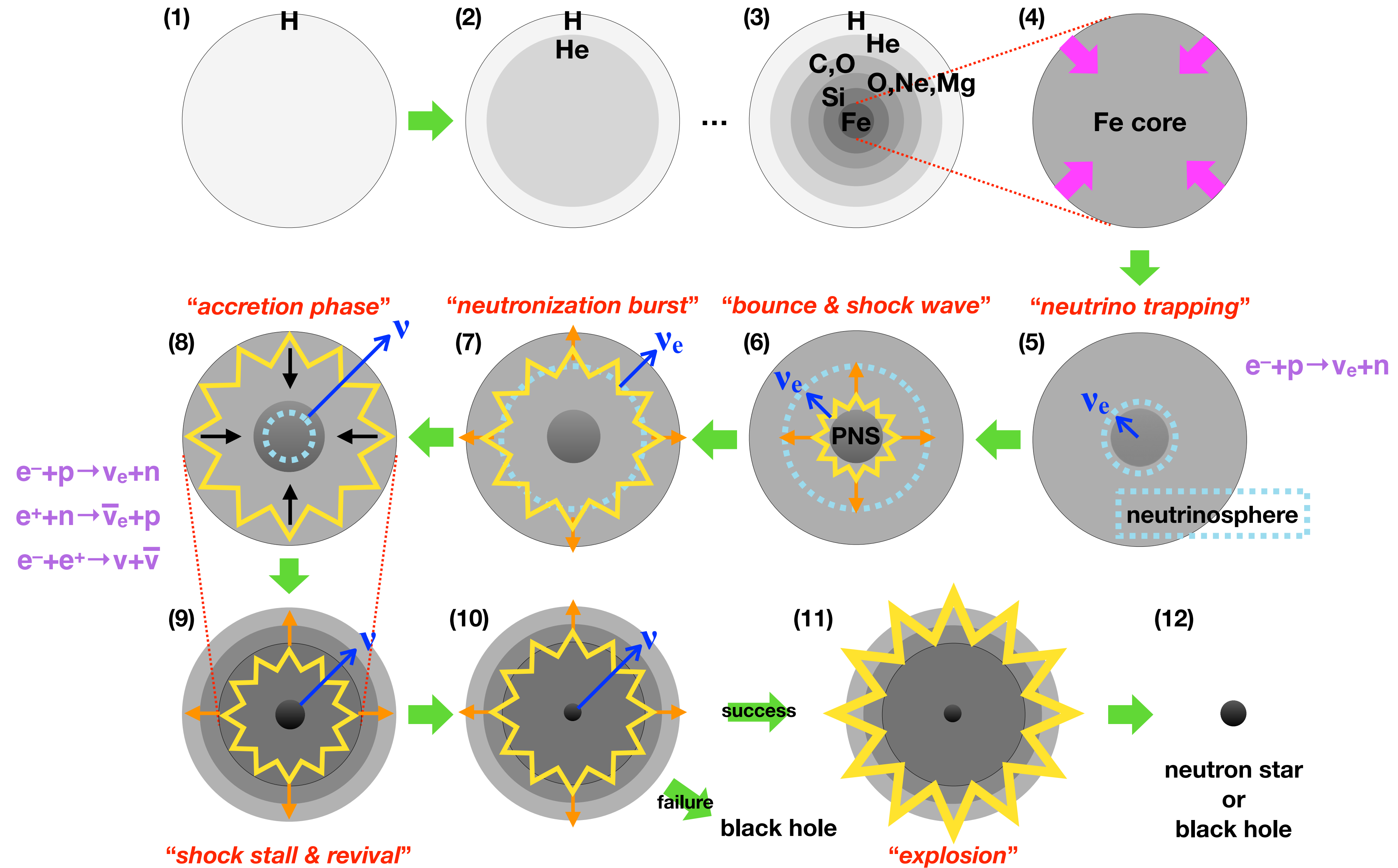


- [ApJ 937, 30 \(2022\)](#): **New proposal to constrain stellar collapse fate with DSNB**
  - Model implemented based on numerical calculations performed separately for the canonical mass neutron star, the high mass neutron star, and the black hole.
  - Sensitivity to the fractions of each remnant estimated for different nuclear EOS and mass hierarchy.
- [ApJ 953, 151 \(2023\)](#): **DSNB flux based on a novel galactic chemical evolution model**
  - New evolution model proposed based on the assumptions of BH formation for progenitors with  $>18M_{\text{Sun}}$  and galaxy-dependent IMF.
  - The resulting DSNB flux showing a unique feature of enhancements at high and low energies.
  - Signal detectability and model discrimination possibility discussed.

**Thanks for your attention!**

# Supplements

# “Neutrino Heating” Scenario



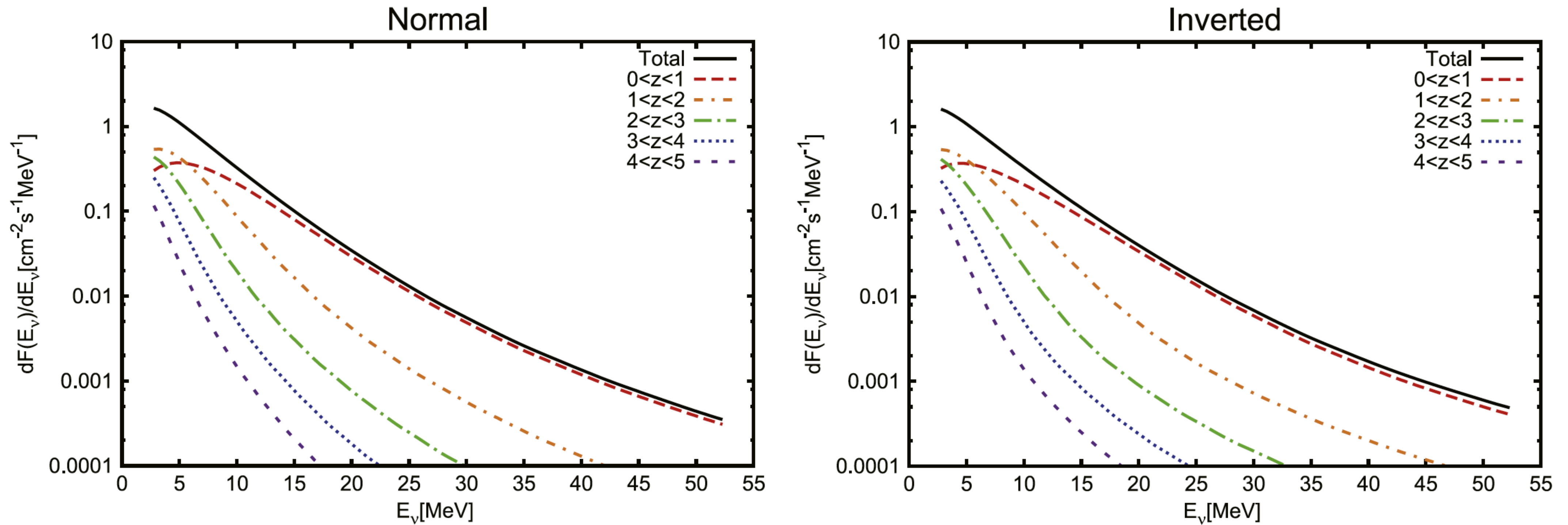
# Mass Hierarchy Impact

$$\begin{aligned}
 \frac{dN_{\bar{\nu}_e}}{dE_\nu} &= |U_{e1}|^2 \frac{dN_{\bar{\nu}_1}}{dE_\nu} + |U_{e2}|^2 \frac{dN_{\bar{\nu}_2}}{dE_\nu} + |U_{e3}|^2 \frac{dN_{\bar{\nu}_3}}{dE_\nu} \\
 &= \cos^2 \theta_{12} \cos^2 \theta_{13} \frac{dN_{\bar{\nu}_1}}{dE_\nu} + \sin^2 \theta_{12} \cos^2 \theta_{13} \frac{dN_{\bar{\nu}_2}}{dE_\nu} + \sin^2 \theta_{13} \frac{dN_{\bar{\nu}_3}}{dE_\nu} \\
 &\sim 0.68 \cdot \frac{dN_{\bar{\nu}_1}}{dE_\nu} + 0.30 \cdot \frac{dN_{\bar{\nu}_2}}{dE_\nu} + 0.02 \cdot \frac{dN_{\bar{\nu}_3}}{dE_\nu},
 \end{aligned}$$

$$\Rightarrow \frac{dN_{\bar{\nu}_e}}{dE_\nu} \sim 0.68 \cdot \frac{dN_{\bar{\nu}_e}^0}{dE_\nu} + 0.32 \cdot \frac{dN_{\bar{\nu}_x}^0}{dE_\nu}. \quad \text{Normal: All flavors}$$

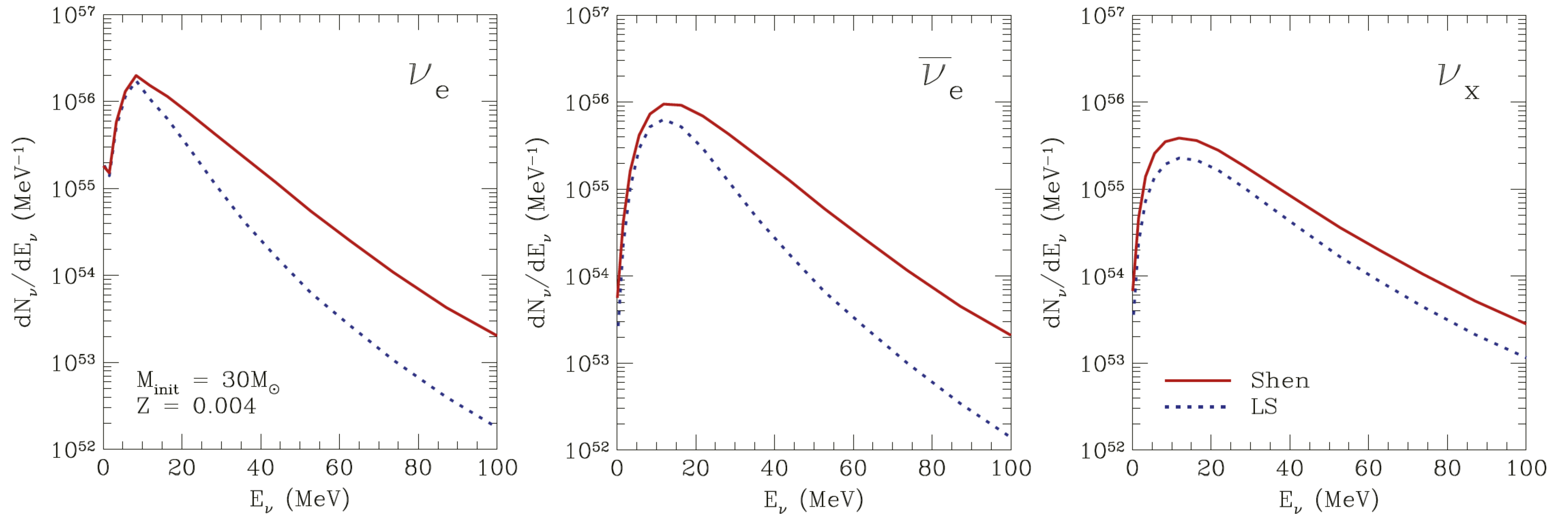
$$\Rightarrow \frac{dN_{\bar{\nu}_e}}{dE_\nu} \sim \frac{dN_{\bar{\nu}_x}^0}{dE_\nu}. \quad \text{Inverted: NuMu+NuTau}$$

# Contributions from Different Redshifts



**Figure 10.** Total fluxes of SRNs (solid) and contributions from various redshift ranges for the reference model. The lines except for the solid line correspond, from top to bottom, to the redshift ranges  $0 < z < 1$ ,  $1 < z < 2$ ,  $2 < z < 3$ ,  $3 < z < 4$ , and  $4 < z < 5$ , for  $E_\nu > 10$  MeV. The left and right panels show the cases for normal and inverted mass hierarchies, respectively.

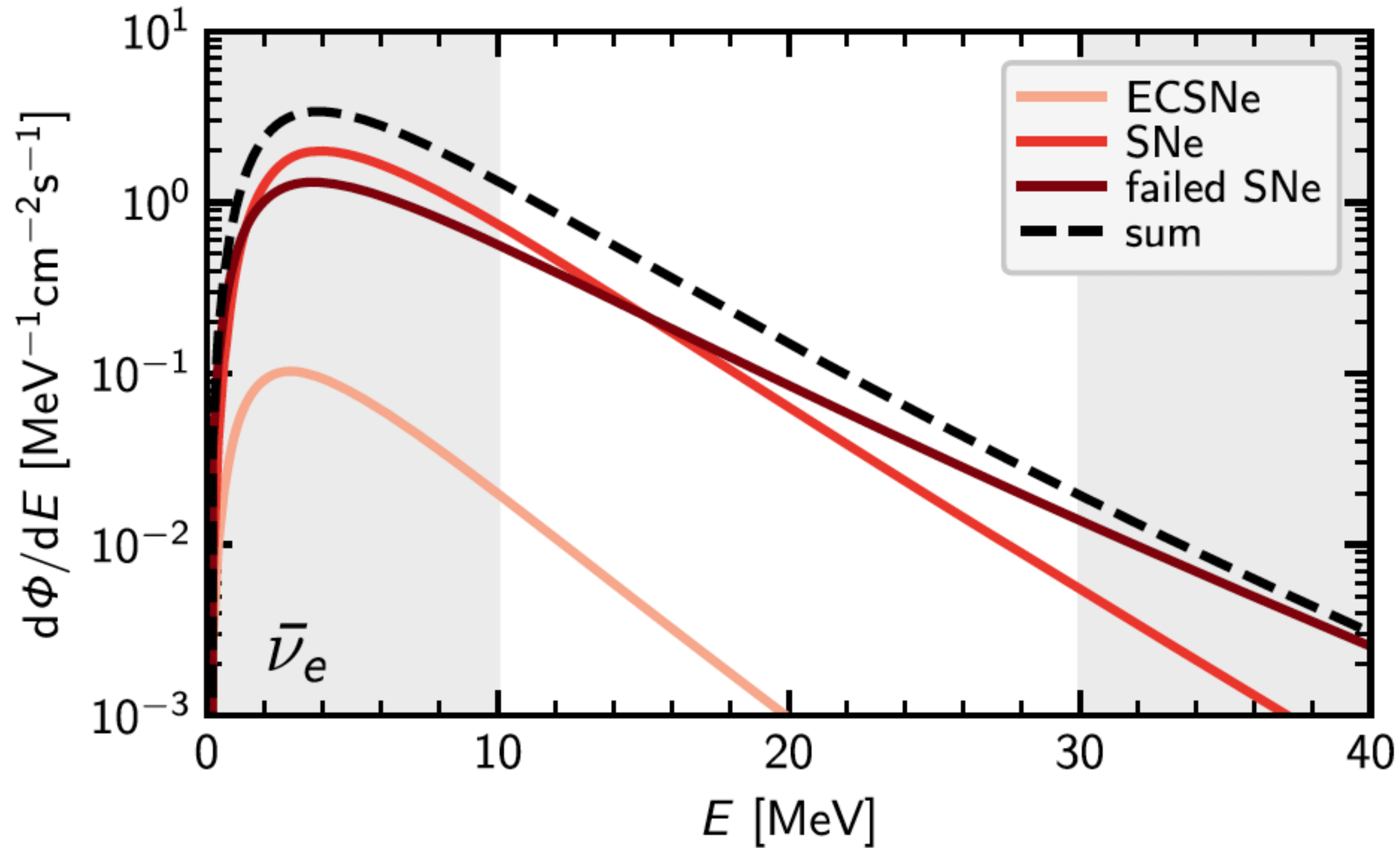




**Figure 6.** Neutrino number spectra for black hole formation with  $30M_{\odot}$ ,  $Z = 0.004$  and Shen EOS (solid) and LS EOS (dotted). The left, central, and right panels correspond to  $\nu_e$ ,  $\bar{\nu}_e$ , and  $\nu_x$  ( $=\nu_{\mu} = \bar{\nu}_{\mu} = \nu_{\tau} = \bar{\nu}_{\tau}$ ), respectively.

Shen is stiffer than LS (maximum mass of neutron stars is higher).

# Failed SN & Electron-Capture SN Contributions



ECSNe = electron-capture supernovae  
(marginal one around mass threshold, w/ ONeMg core)

D. Kresse et al., ApJ 909, 169 (2021)

# Model Setup Details

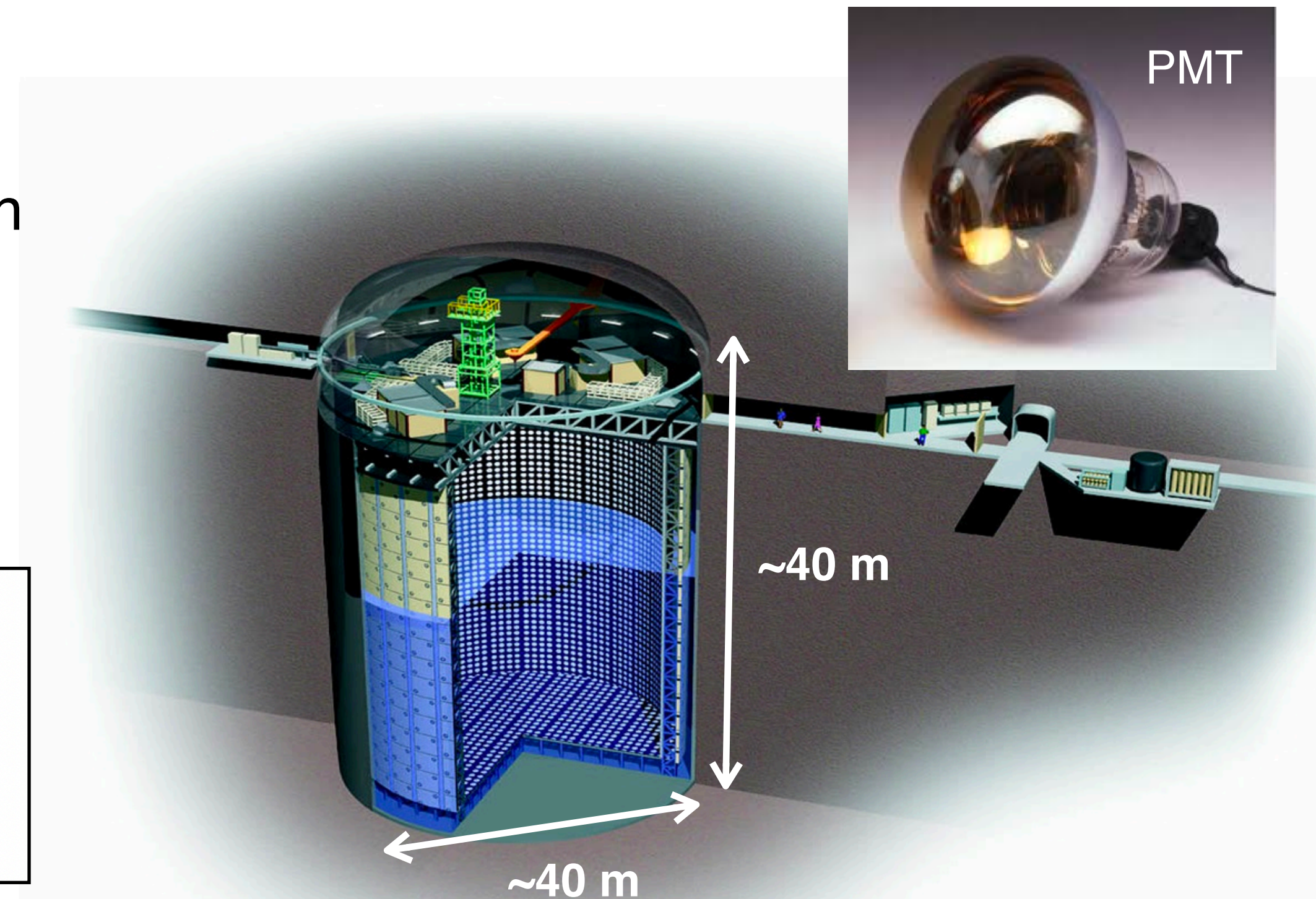
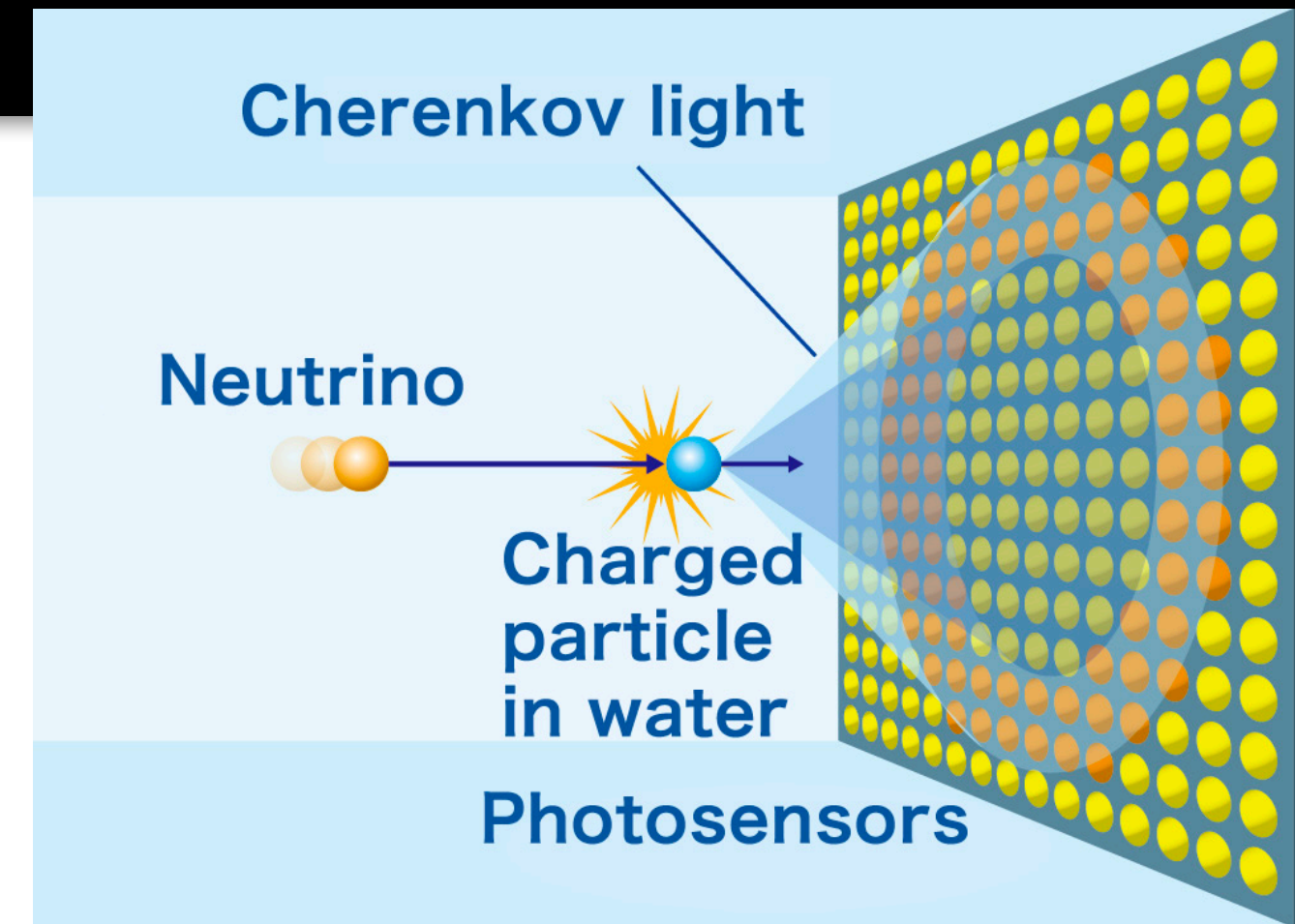
**Table 1.** Properties of the compact remnant and time-integrated neutrino signal for each fate component.

EOS	remnant	$M_{\text{NS,g}}$ [ $M_{\odot}$ ]	$R_{\text{NS}}$ [km]	$t_{\text{BH}}$ [sec]	$\langle E_{\nu_e} \rangle$ [MeV]	$\langle E_{\bar{\nu}_e} \rangle$ [MeV]	$\langle E_{\nu_x} \rangle$ [MeV]	$E_{\nu_e,\text{tot}}$ [ $10^{52}$ erg]	$E_{\bar{\nu}_e,\text{tot}}$ [ $10^{52}$ erg]	$E_{\nu_x,\text{tot}}$ [ $10^{52}$ erg]
Togashi	canonical mass NS	1.32	11.5	—	9.2	10.9	11.8	4.47	4.07	4.37
	high mass NS	1.63	11.5	—	9.5	11.2	11.9	7.26	6.93	7.17
	BH (Failed SN)	—	—	0.533	16.1	20.4	23.4	6.85	5.33	2.89
LS220	canonical mass NS	1.34	12.7	—	9.1	10.7	11.3	4.25	3.84	3.94
	high mass NS	1.65	12.4	—	9.8	11.2	11.2	7.29	6.88	6.36
	BH (Failed SN)	—	—	0.342	12.5	16.4	22.3	4.03	2.87	2.11
Shen	canonical mass NS	1.35	14.3	—	9.0	10.6	11.3	3.65	3.22	3.35
	high mass NS	1.67	14.1	—	9.6	11.2	11.2	6.22	5.88	5.40
	BH (Failed SN)	—	—	0.842	17.5	21.7	23.4	9.49	8.10	4.00

NOTE—For the successful SN models,  $M_{\text{NS,g}}$  and  $R_{\text{NS}}$  are the gravitational mass and radius of NSs, respectively. For the failed SN models,  $t_{\text{BH}}$  is the time to a BH formation measured from the core bounce.  $\langle E_{\nu_i} \rangle$  and  $E_{\nu_i,\text{tot}}$  are the average and total energies of the time-integrated neutrino signal for  $\nu_i$ , where  $\nu_x$  represents the average of  $\nu_{\mu}$ ,  $\bar{\nu}_{\mu}$ ,  $\nu_{\tau}$ , and  $\bar{\nu}_{\tau}$ .

# Super-Kamiokande Experiment

- A **water Cherenkov** detector located 1,000 m under mountain in Japan.
  - Detecting Cherenkov light emitted from relativistic charged particles in water.
- Composed of two layers;
  - **Inner:** 11,129 20-inch PMTs, detection volume  $\sim 22.5$  kton
  - **Outer:** 1,885 8-inch PMTs, veto for atmospheric  $\mu$
- Operated since 1996, separated into *six* phases depending on detector configuration (SK-I to VI).
  - Results from **SK-I to IV (5,823 days in total)** will be shown.



PHYSICAL REVIEW D **104**, 122002 (2021)

## Diffuse supernova neutrino background search at Super-Kamiokande

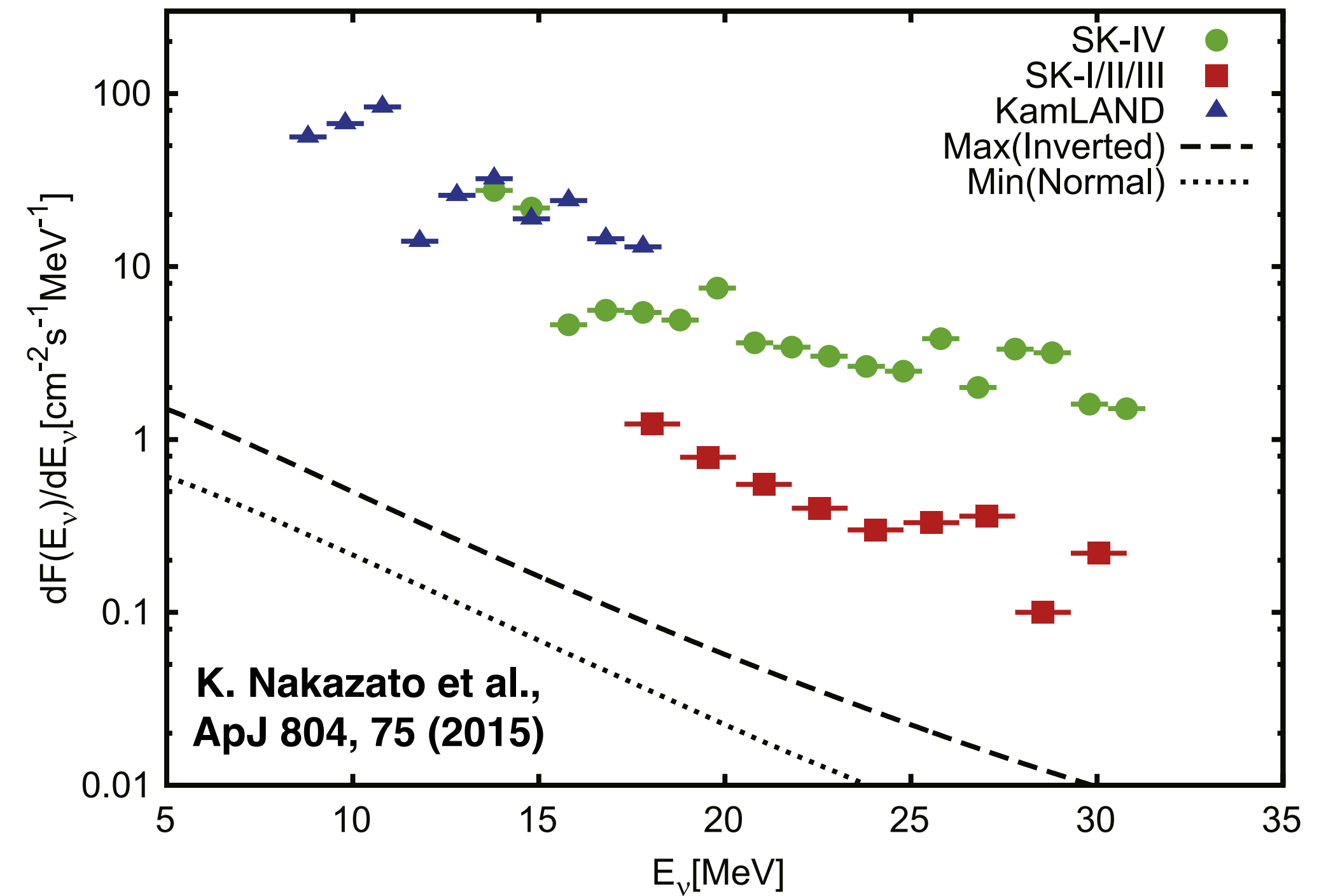
K. Abe,<sup>1,43</sup> C. Bronner,<sup>1</sup> Y. Hayato,<sup>1,43</sup> K. Hiraide,<sup>1</sup> M. Ikeda,<sup>1</sup> S. Imaizumi,<sup>1</sup> J. Kameda,<sup>1,43</sup> Y. Kanemura,<sup>1</sup> Y. Kataoka,<sup>1</sup>

## 1. SK-I/II/III 2,853 days [Phys. Rev. D 85, 052007 (2012)]

- *No ntag* (only  $e^+$  prompt signal)
- Higher energy threshold ( $E_\nu > 17.3$  MeV)
- Atmospheric  $\nu$  bkg estimated with old models
- **Unbinned likelihood spectral fitting**

## 2. SK-IV 960 days [Astropart. Phys. 60, 41 (2015)]

- **Ntag applied for the first time** (electronics upgrade in SK-IV)
- Poor performance of spallation cut & ntag
- Only accidental bkg estimated (underestimate of bkg)
  - No framework for ntag on atmospheric  $\nu$  simulations
  - No reliable method for estimation of  ${}^9\text{Li}$
- **Bin-by-bin differential upper limits**; no spectral fitting result



New analysis improves *pros* and  
solves *cons* in old analyses!

# Improvements in SK-IV Analysis

- ★ Larger statistics: SK-IV 2,970 days + SK-I/II/III 2,853 days (**5,823 days** in total)
  - ★ Better **ntag** performance
  - ★ Improved **spallation cut**
- } → Search energy threshold lowered
- ★ Novel method of estimating **spallation <sup>9</sup>Li bkg** developed
  - ★ More reliable and precise estimation of **atmospheric ν bkg**
  - ★ Bin-by-bin upper limits & unbinned likelihood fitting performed w/ **better systematic uncertainty treatment**

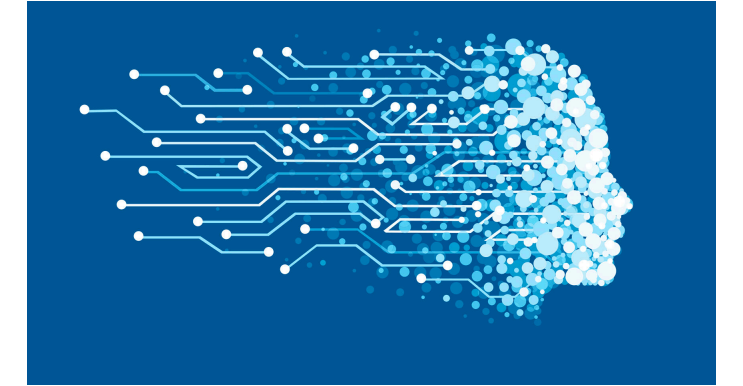
**Most sensitive search achieved!**

★ My work

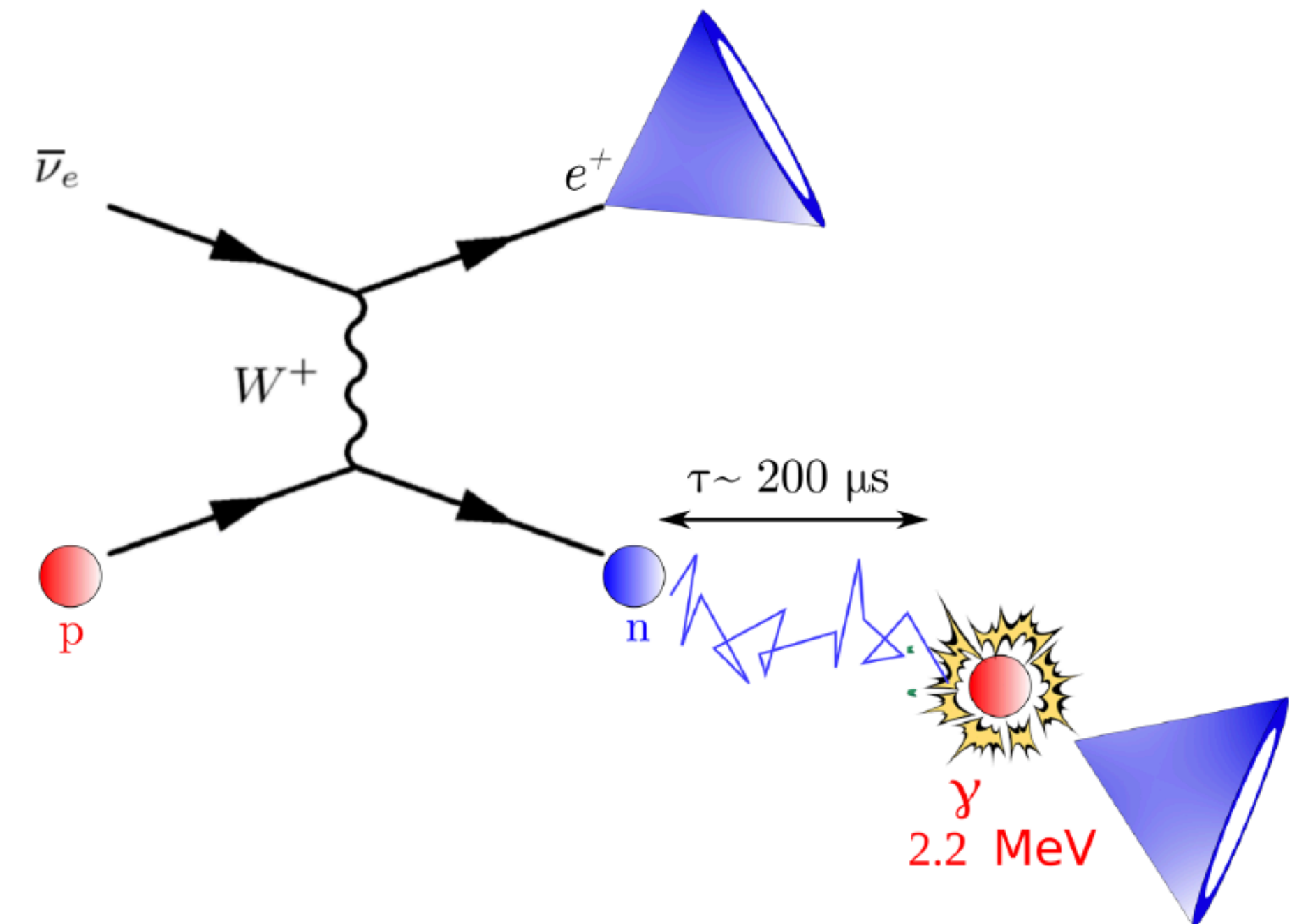
★ My colleagues' work

# DSNB Signal at Super-K

- Signal = **inverse beta decay (IBD)**,  $\bar{\nu}_e + p \rightarrow e^+ + n$  (largest xsec)
  - $e^+$  = “*prompt*” signal (main signal range: **10~30 MeV**)
  - $n$  = “*delayed*” signal via **a 2.2 MeV  $\gamma$ -ray** from thermal capture on hydrogen ( $\tau \sim 200 \mu\text{s}$ )



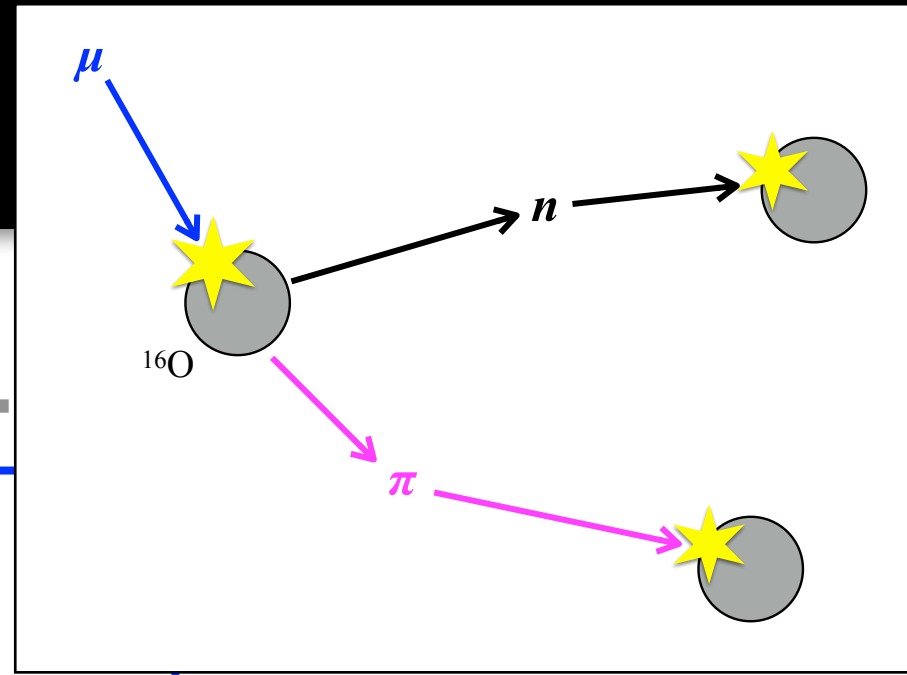
- Many types of backgrounds mimicking this signature.
  - Need to *reduce* them to as much as DSNB signal.
  - Need to *precisely model/estimate* them.



# Background

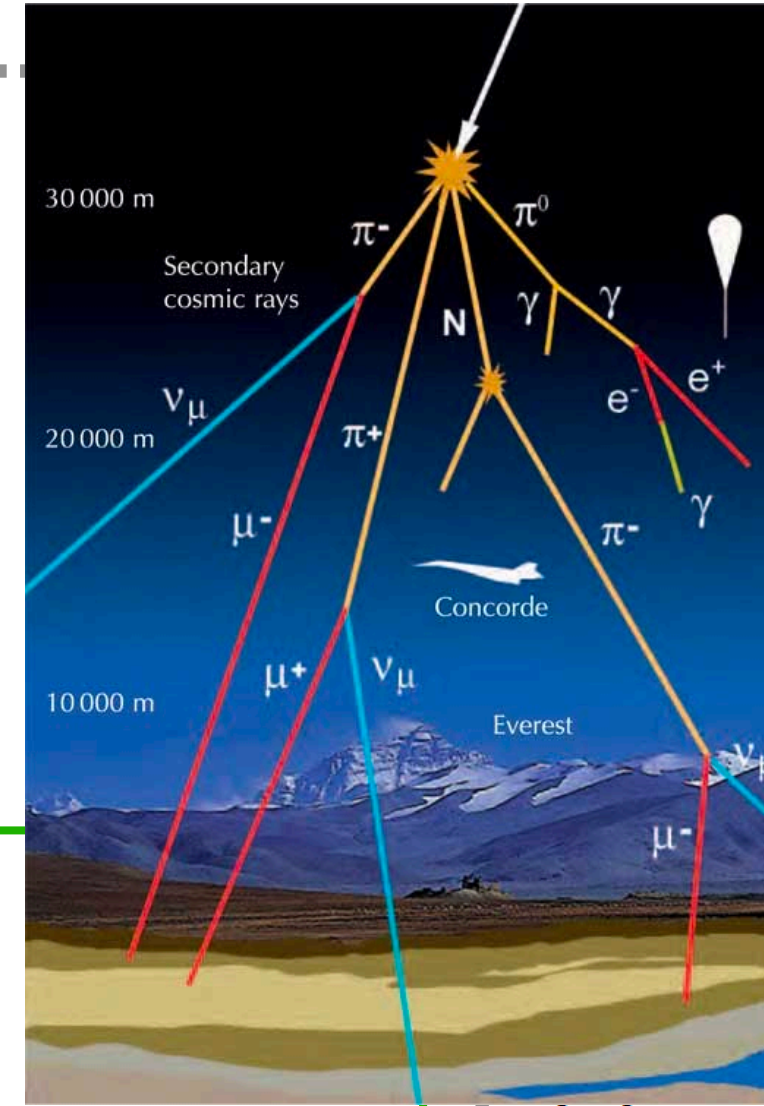
## Muon spallation

- Decay without neutron
- Decay with neutron ( ${}^9\text{Li}$ , etc)

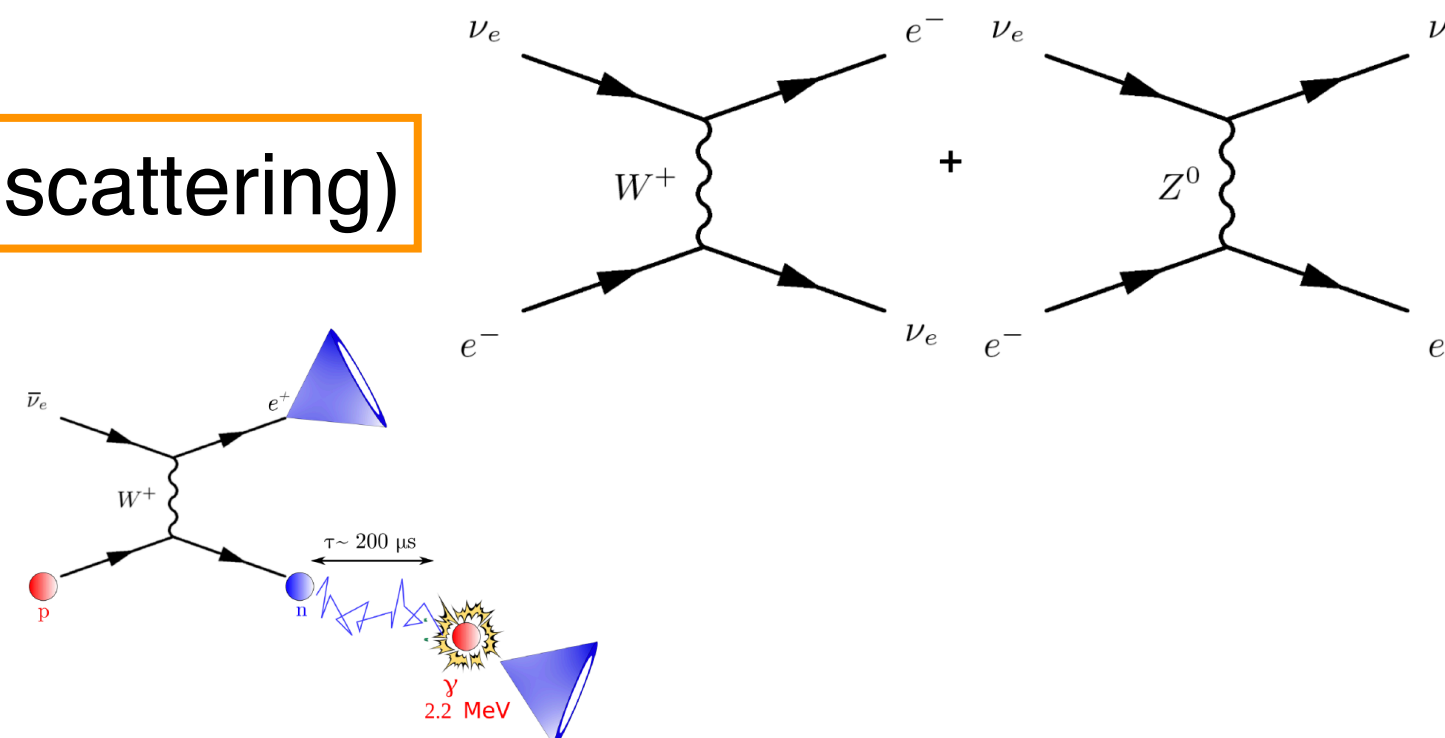


## Atmospheric neutrinos

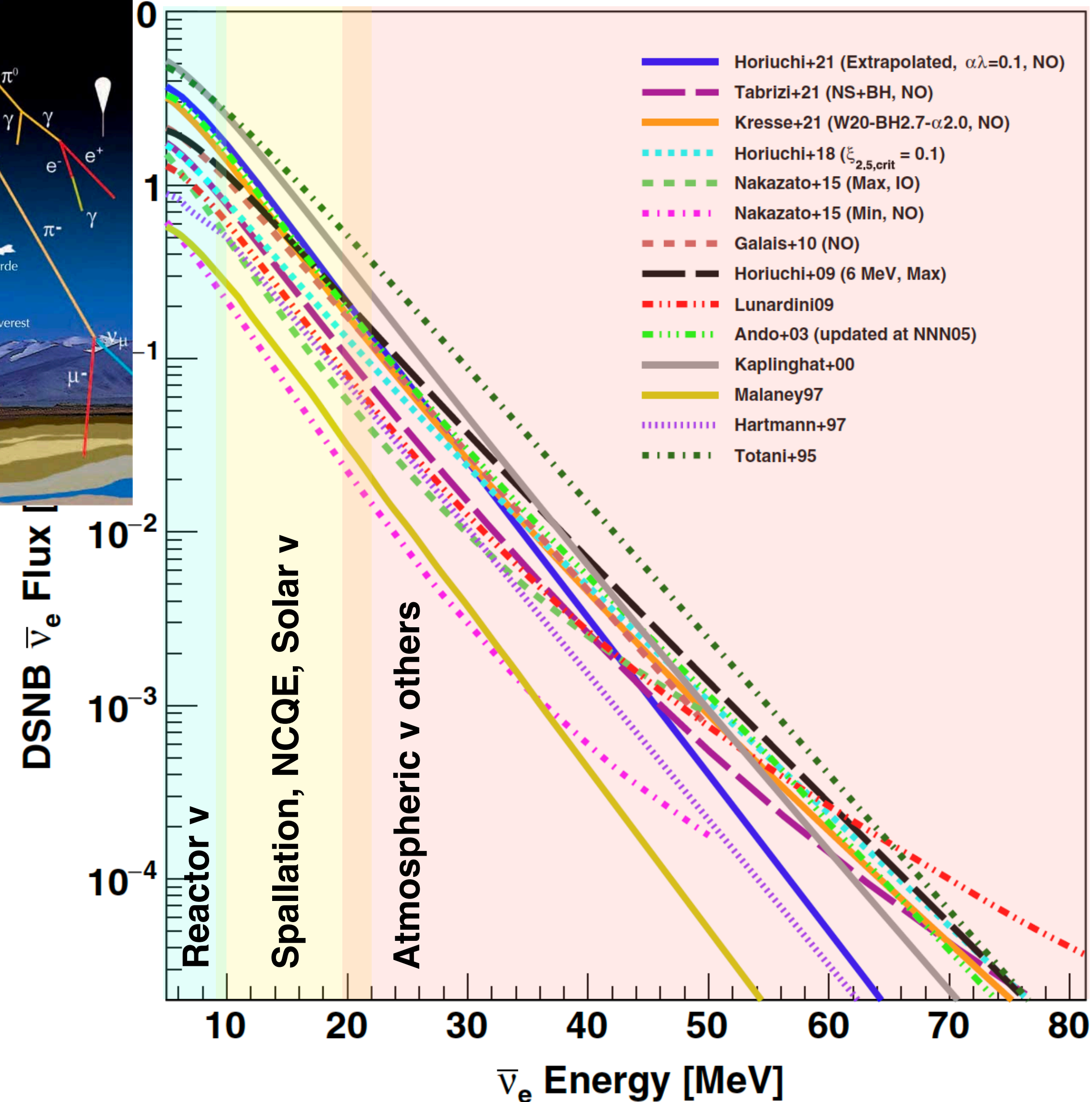
- Neutral-current quasielastic interactions (NCQE)
- $\nu_e/\bar{\nu}_e$  charged-current (CC) interactions
- Muon/pion-producing interactions (CCQE, CC1 $\pi$ , NC1 $\pi$ , etc)



## Solar neutrinos (electron scattering)



## Reactor neutrinos (IBD)

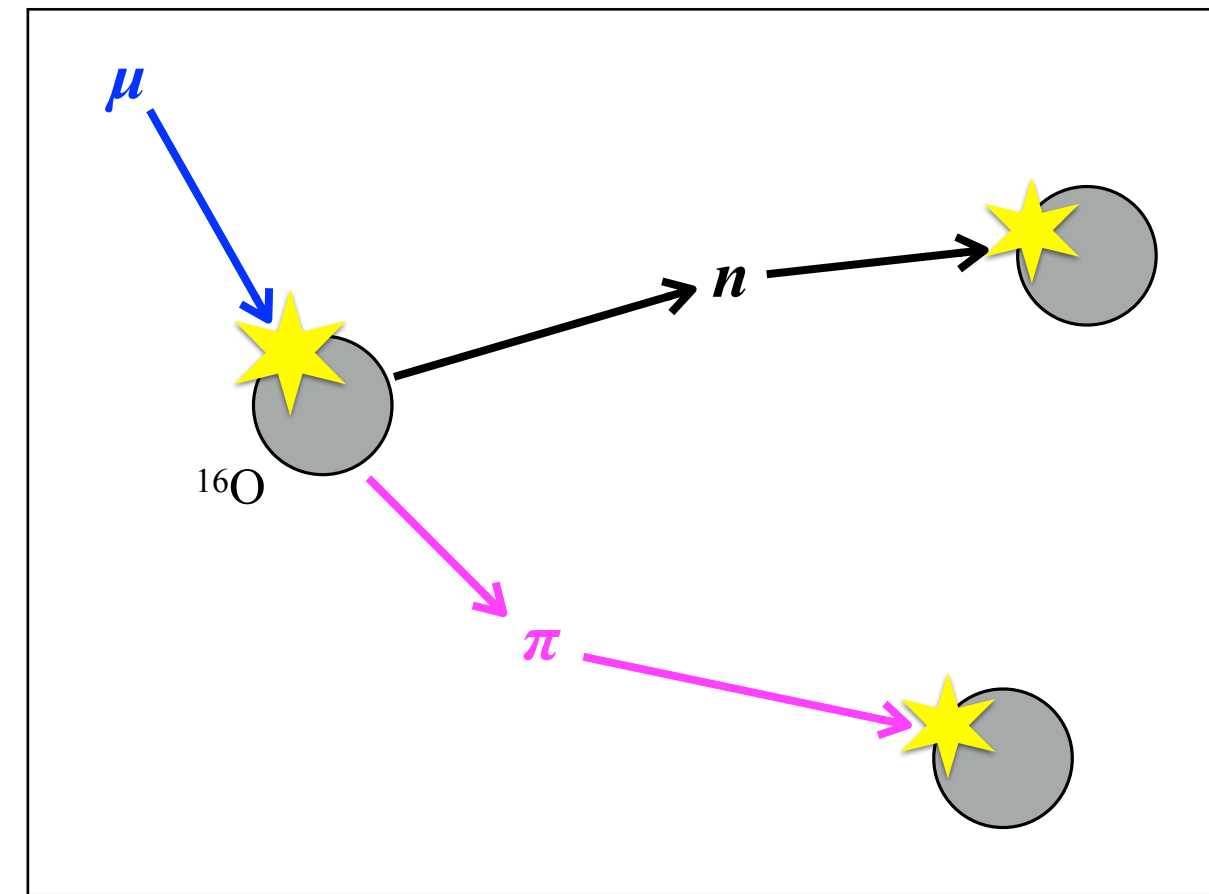




# Muon Spallation

## Muon spallation

- Decay without neutron
- Decay with neutron ( ${}^9\text{Li}$ , etc)

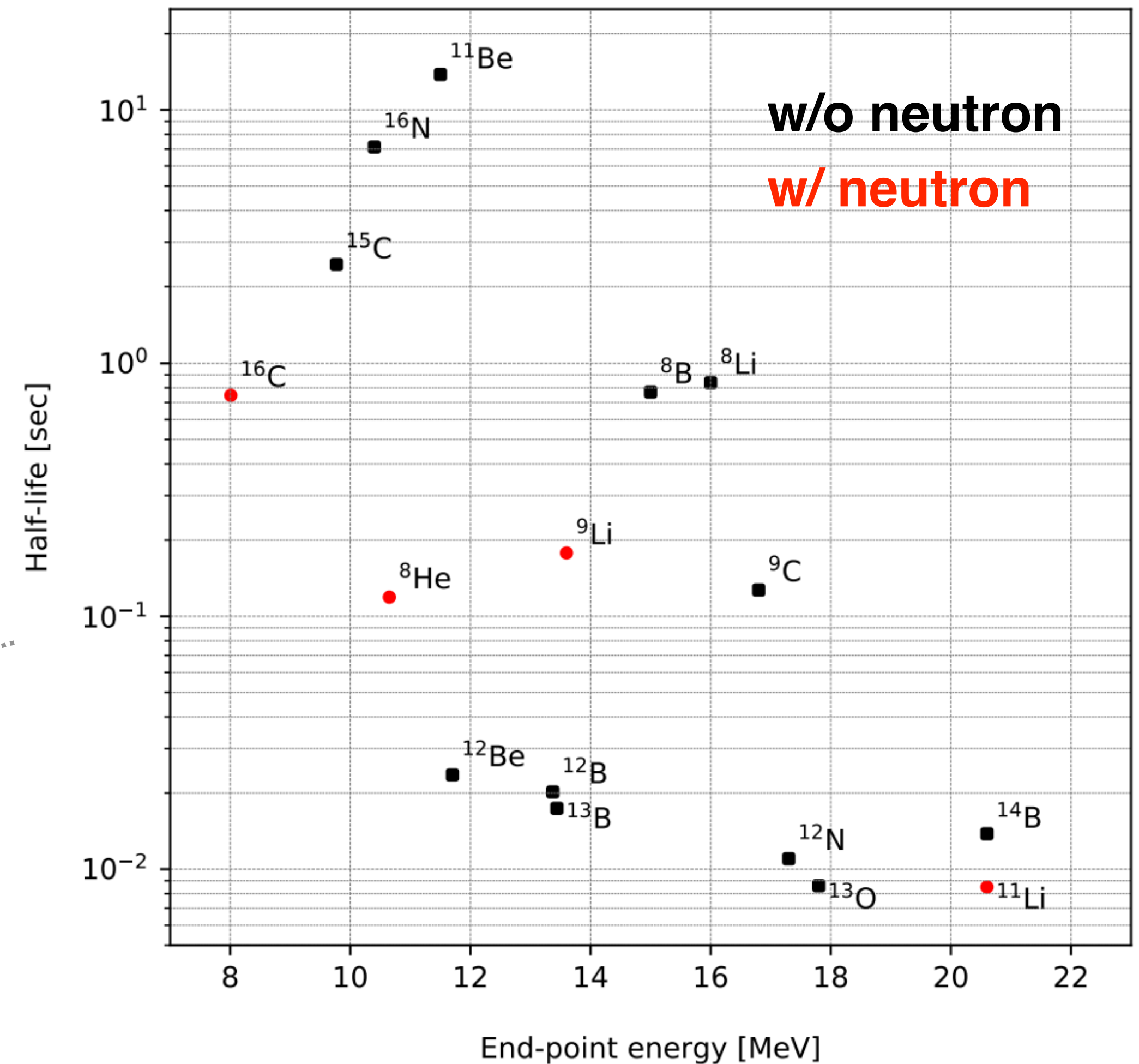


Energetic atmospheric  $\mu$  on oxygen in Super-K

Hadronic particles

Radioactive isotopes

$\beta/\gamma$  decay (delayed by msec~10 sec)



# Muons for Spallation

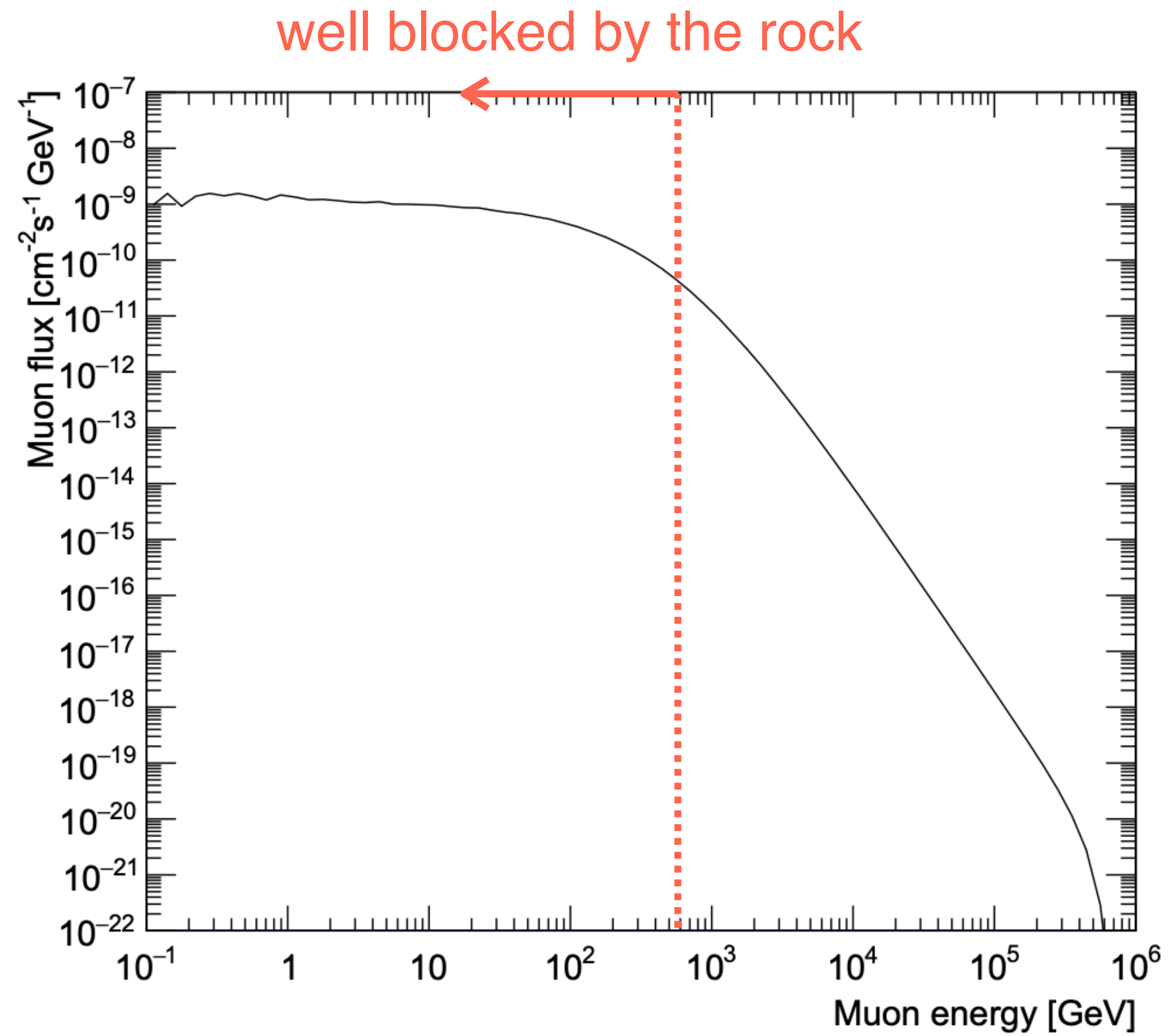
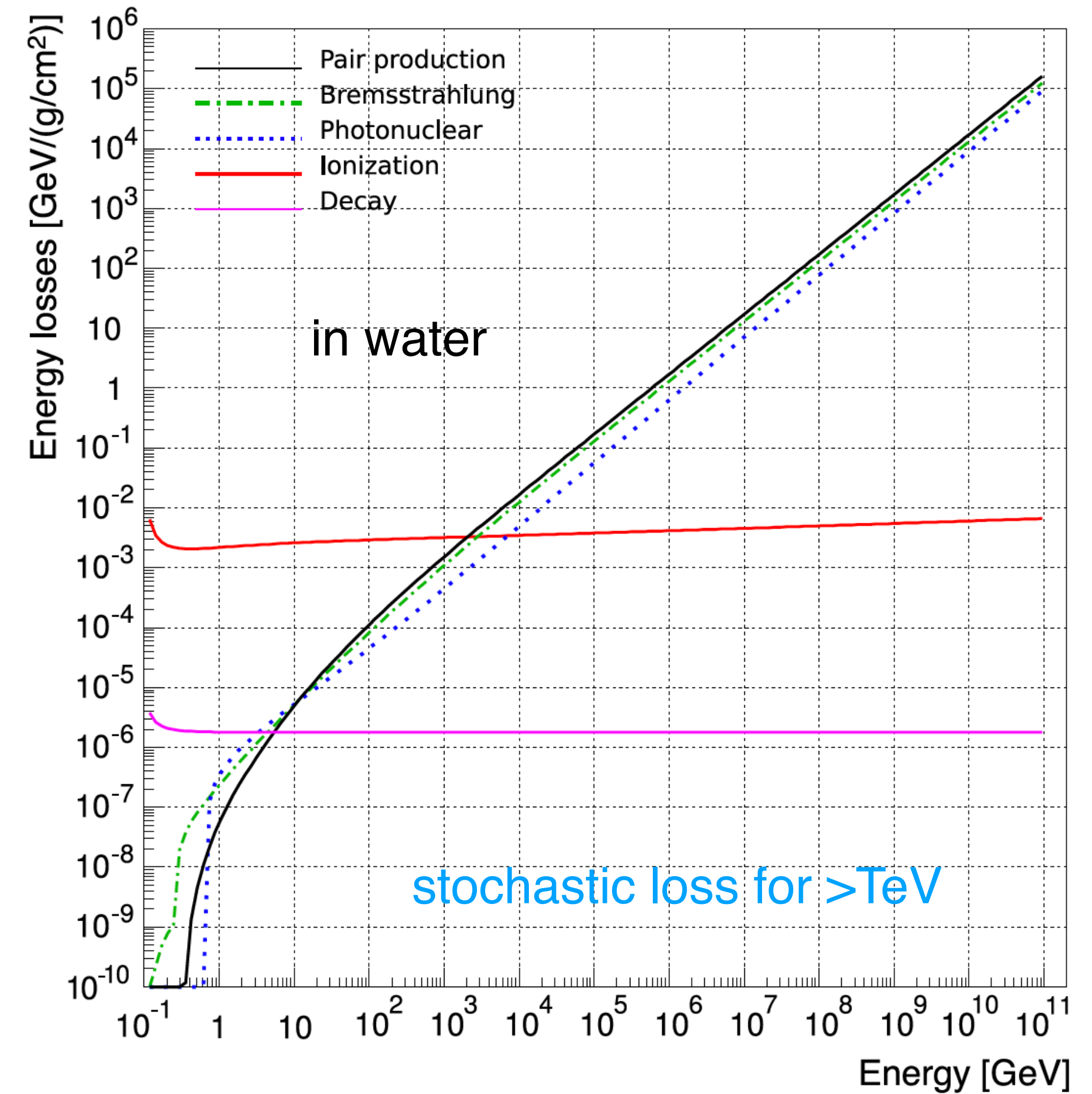


FIG. 2. Muon energy spectrum at the location of SK detector in the mine inside Mt. Ikenoyama.



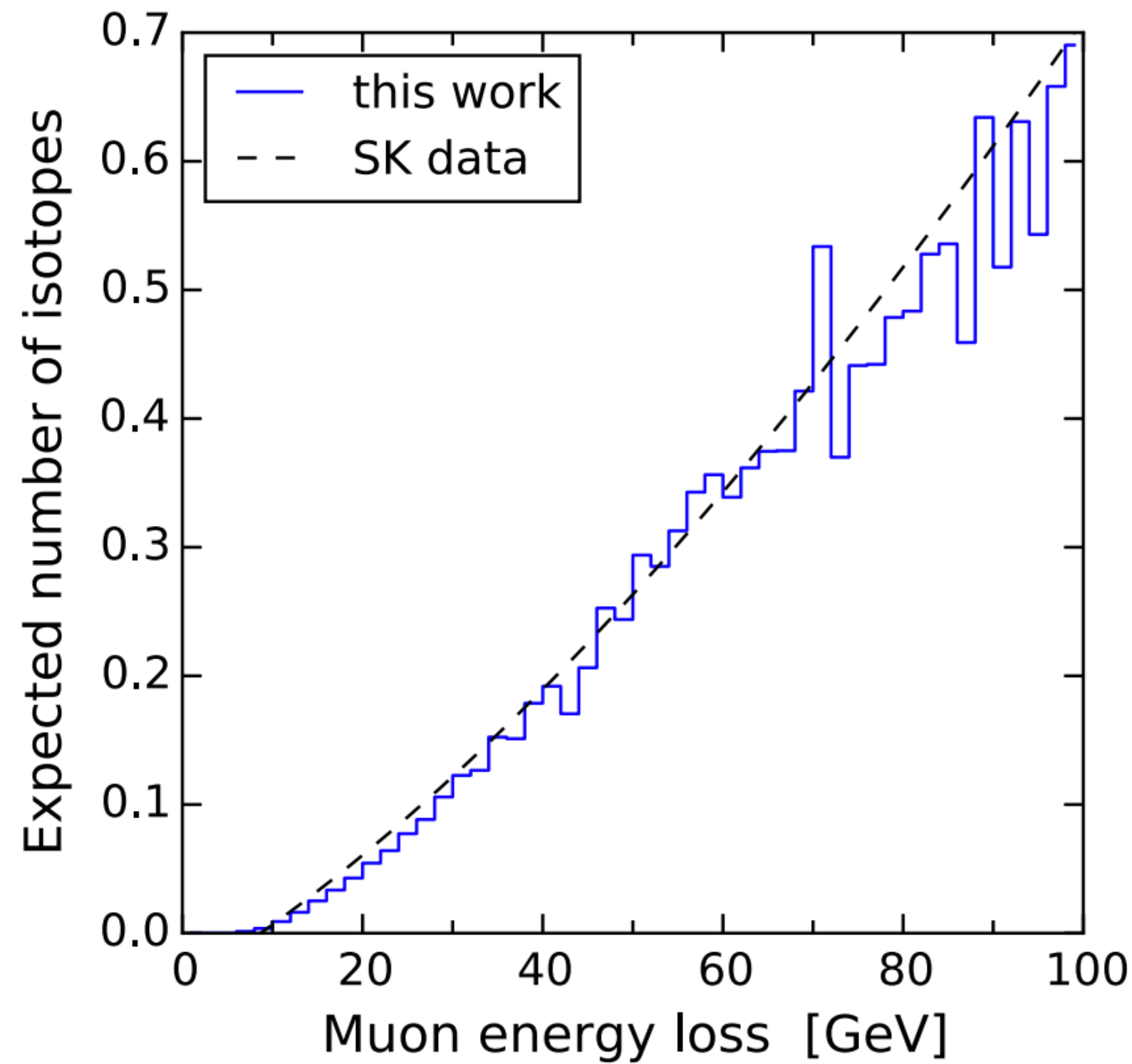


FIG. 8 (color online). The expected number of background isotopes as a function of the total muon energy loss. The solid line is our calculation assuming vertical through going muons that travel 32.2 m in the FV, and the dashed line is the (corrected to match assumptions) Super-K measurement.

Isotope	Half-life (s)	Decay mode	Yield (total) ( $\times 10^{-7} \mu^{-1} \text{g}^{-1} \text{cm}^2$ )	Yield ( $E > 3.5 \text{ MeV}$ ) ( $\times 10^{-7} \mu^{-1} \text{g}^{-1} \text{cm}^2$ )	Primary process
			2030		
$^{18}\text{N}$	0.624	$\beta^-$	0.02	0.01	$^{18}\text{O}(n,p)$
$^{17}\text{N}$	4.173	$\beta^- n$ <i><math>^{18}\text{O}</math> is very little (~0.2%)</i>	0.59	0.02	$^{18}\text{O}(n,n+p)$
$^{16}\text{N}$	7.13	$\beta^- \gamma$ (66%), $\beta^-$ (28%)	18	18	( $n,p$ )
$^{16}\text{C}$	0.747	$\beta^- n$ <i>small amount &gt;3.5 MeV</i>	0.02	0.003	( $\pi^-, n+p$ )
$^{15}\text{C}$	2.449	$\beta^- \gamma$ (63%), $\beta^-$ (37%)	0.82	0.28	( $n,2p$ )
$^{14}\text{B}$	0.0138	$\beta^- \gamma$	0.02	0.02	( $n,3p$ )
$^{13}\text{O}$	0.0086	$\beta^+$	0.26	0.24	( $\mu^-, p+2n+\mu^- + \pi^-$ )
$^{13}\text{B}$	0.0174	$\beta^-$	1.9	1.6	( $\pi^-, 2p+n$ )
$^{12}\text{N}$	0.0110	$\beta^+$	1.3	1.1	( $\pi^+, 2p+2n$ )
$^{12}\text{B}$	0.0202	$\beta^-$	12	9.8	( $n,\alpha+p$ )
$^{12}\text{Be}$	0.0236	$\beta^-$	0.10	0.08	( $\pi^-, \alpha+p+n$ )
$^{11}\text{Be}$	13.8	$\beta^-$ (55%), $\beta^- \gamma$ (31%)	0.81	0.54	( $n,\alpha+2p$ )
$^{11}\text{Li}$	0.0085	$\beta^- n$ <i>very short life-time</i>	0.01	0.01	( $\pi^+, 5p+\pi^+ + \pi^0$ )
$^9\text{C}$	0.127	$\beta^+$	0.89	0.69	( $n,\alpha+4n$ )
$^9\text{Li}$	0.178	$\beta^- n$ (51%), $\beta^-$ (49%)	1.9	1.5	( $\pi^-, \alpha+2p+n$ )
$^8\text{B}$	0.77	$\beta^+$ <i>low energy</i>	5.8	5.0	( $\pi^+, \alpha+2p+2n$ )
$^8\text{Li}$	0.838	$\beta^-$ <i>(end point ~8 MeV)</i>	13	11	( $\pi^-, \alpha+^2\text{H}+p+n$ )
$^8\text{He}$	0.119	$\beta^- \gamma$ (84%), $\beta^- n$ (16%)	0.23	0.16	( $\pi^-, ^3\text{H}+4p+n$ )
$^{15}\text{O}$			351		( $\gamma,n$ )
$^{15}\text{N}$			773		( $\gamma,p$ )
$^{14}\text{O}$			13		( $n,3n$ )
$^{14}\text{N}$			295		( $\gamma,n+p$ )
$^{14}\text{C}$			64		( $n,n+2p$ )
$^{13}\text{N}$			19		( $\gamma, ^3\text{H}$ )
$^{13}\text{C}$			225		( $n, ^2\text{H}+p+n$ )
$^{12}\text{C}$			792		( $\gamma,\alpha$ )
$^{11}\text{C}$			105		( $n,\alpha+2n$ )
$^{11}\text{B}$			174		( $n,\alpha+p+n$ )
$^{10}\text{C}$			7.6		( $n,\alpha+3n$ )
$^{10}\text{B}$			77		( $n,\alpha+p+2n$ )
$^{10}\text{Be}$			24		( $n,\alpha+2p+n$ )
$^9\text{Be}$			38		( $n,2\alpha$ )
sum			3015	50	

Not direct backgrounds in SK:

- stable
- long half-life
- invisible decay
- low energy

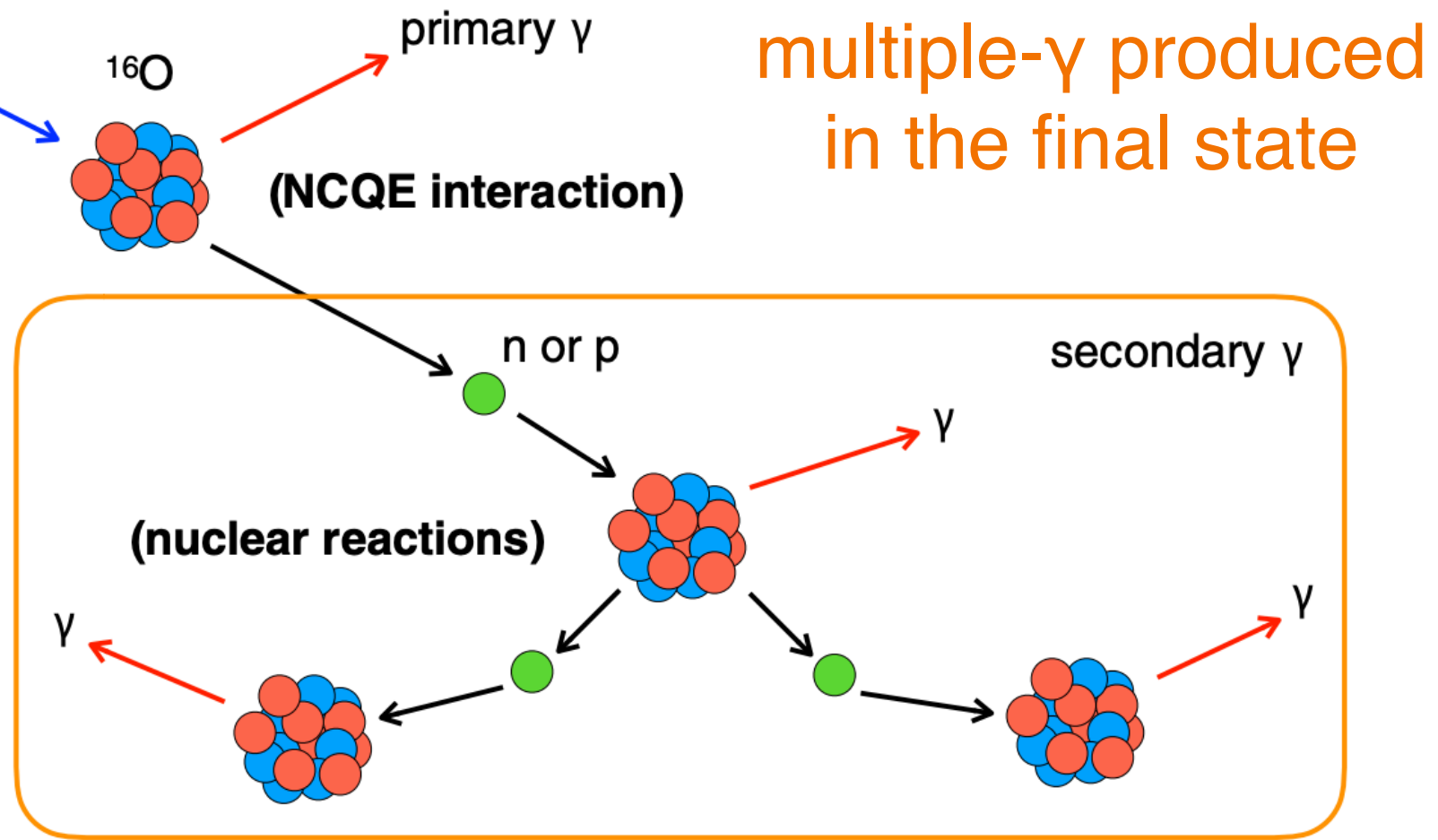
# Atmospheric Neutrinos

★ My own analysis in T2K

PHYSICAL REVIEW D **100**, 112009 (2019)

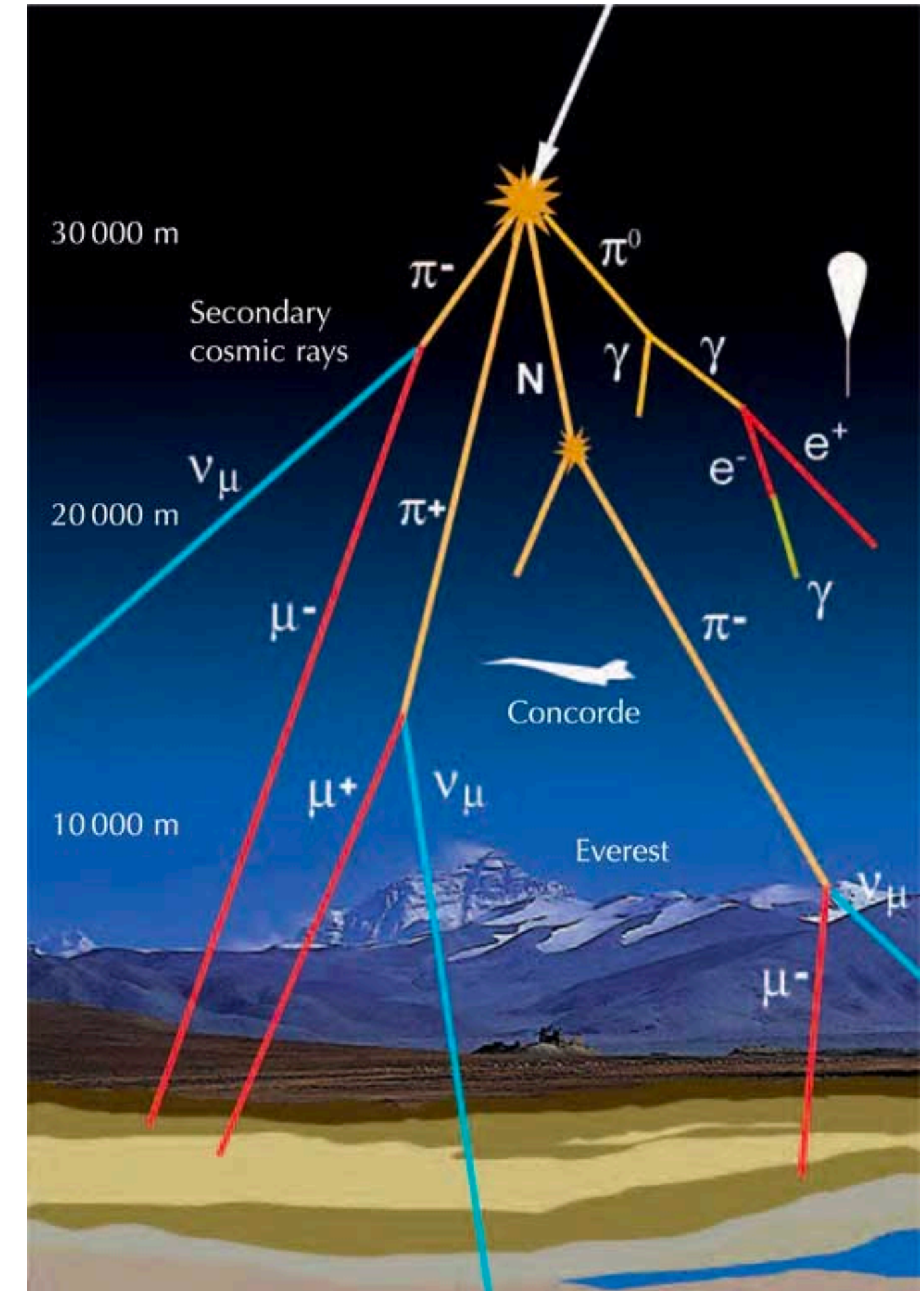
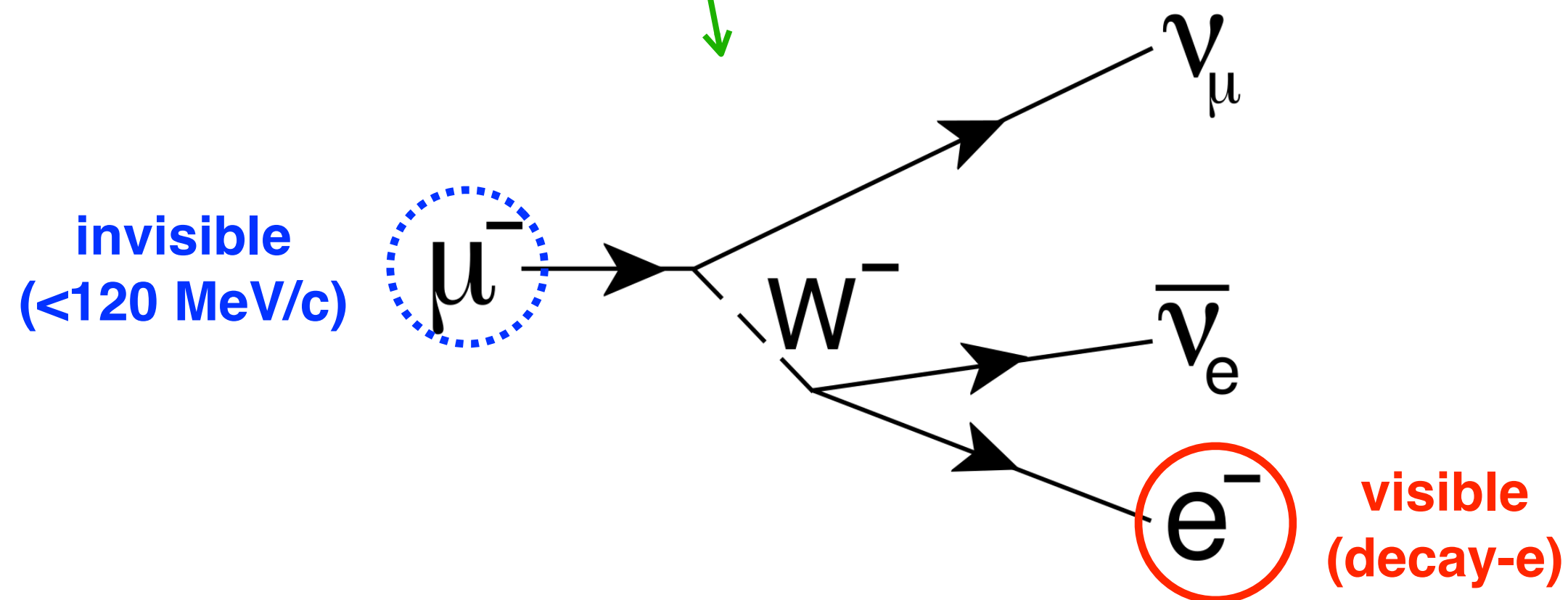
**Measurement of neutrino and antineutrino neutral-current quasielasticlike interactions on oxygen by detecting nuclear deexcitation  $\gamma$  rays**

K. Abe,<sup>55</sup> R. Akutsu,<sup>56</sup> A. Ali,<sup>32</sup> C. Alt,<sup>11</sup> C. Andreopoulos,<sup>53,34</sup> L. Anthony,<sup>34</sup> M. Antonova,<sup>19</sup> S. Aoki,<sup>31</sup> A. Ariga,<sup>2</sup> Y. Ashida,<sup>32</sup> E. T. Atkin,<sup>21</sup> Y. Awataguchi,<sup>58</sup> S. Ban,<sup>32</sup> M. Barbi,<sup>45</sup> G. J. Barker,<sup>65</sup> G. Barr,<sup>42</sup> C. Barry,<sup>34</sup> M. Batkiewicz-Kwasniak,<sup>15</sup> A. Beloshapkin,<sup>26</sup> F. Bench,<sup>34</sup> V. Berardi,<sup>22</sup> S. Berkman,<sup>4,61</sup> L. Berns,<sup>57</sup> S. Bhadra,<sup>69</sup> S. Bienstock,<sup>52</sup> A. Blondel,<sup>13,†</sup> S. Bolognesi,<sup>6</sup> B. Bourguille,<sup>18</sup> S. B. Boyd,<sup>65</sup> D. Brailsford,<sup>33</sup> A. Bravar,<sup>13</sup> D. Bravo Berguño,<sup>1</sup> C. Bronner,<sup>55</sup> A. Bubak,<sup>50</sup> M. Buizza Avanzini,<sup>10</sup> I. Calcutt,<sup>36</sup> T. Campbell,<sup>7</sup> S. Cao,<sup>16</sup>

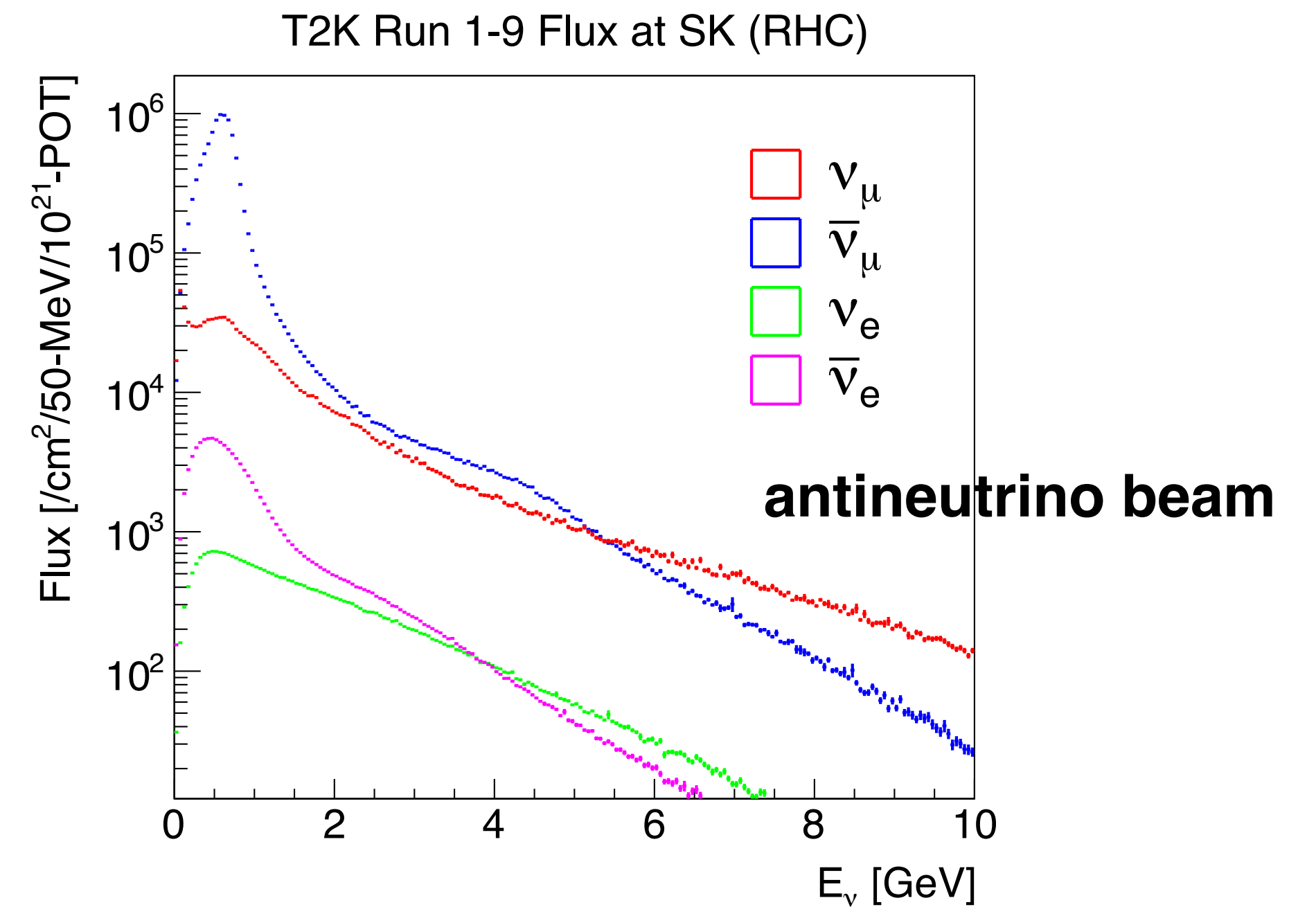
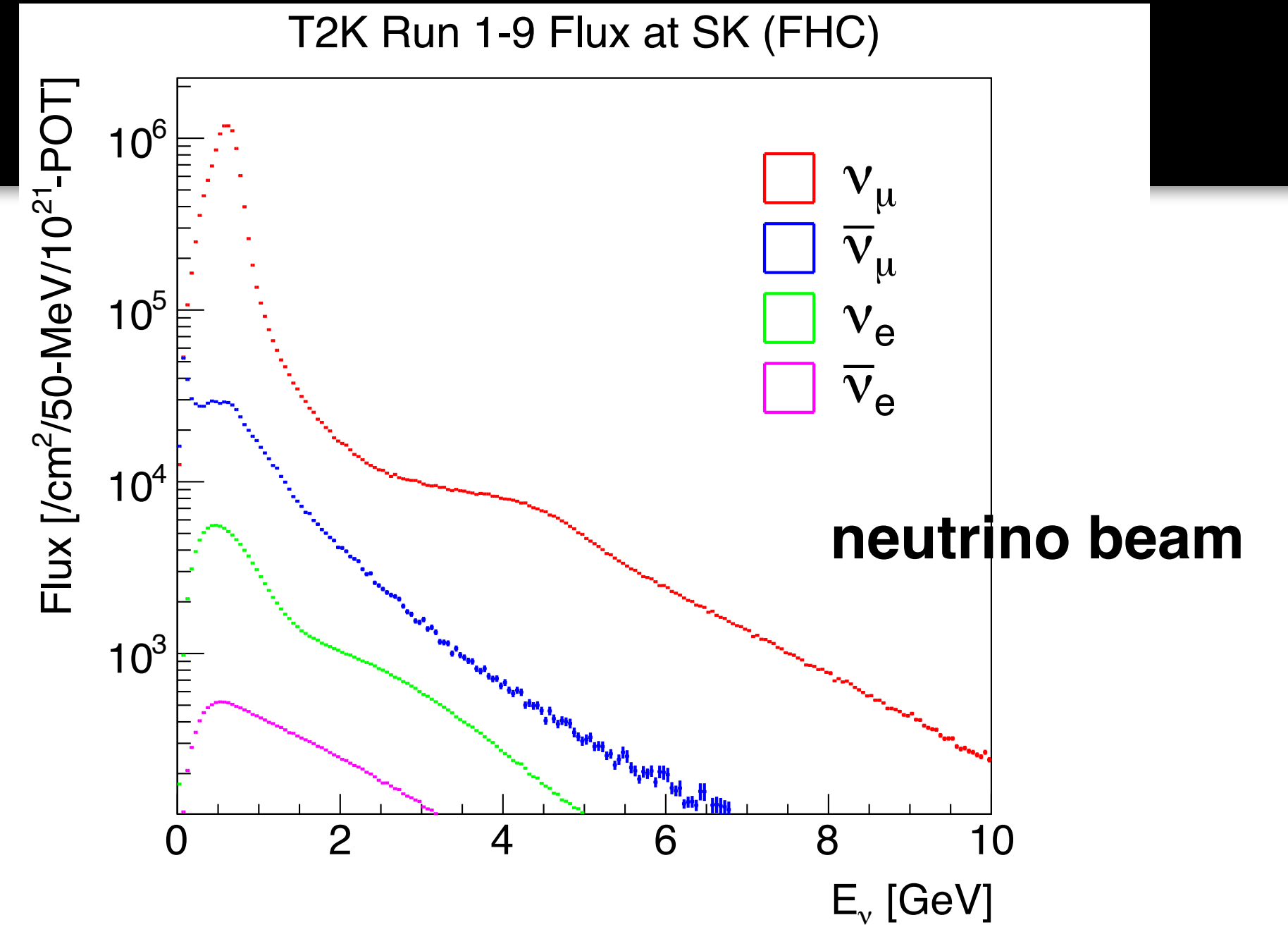
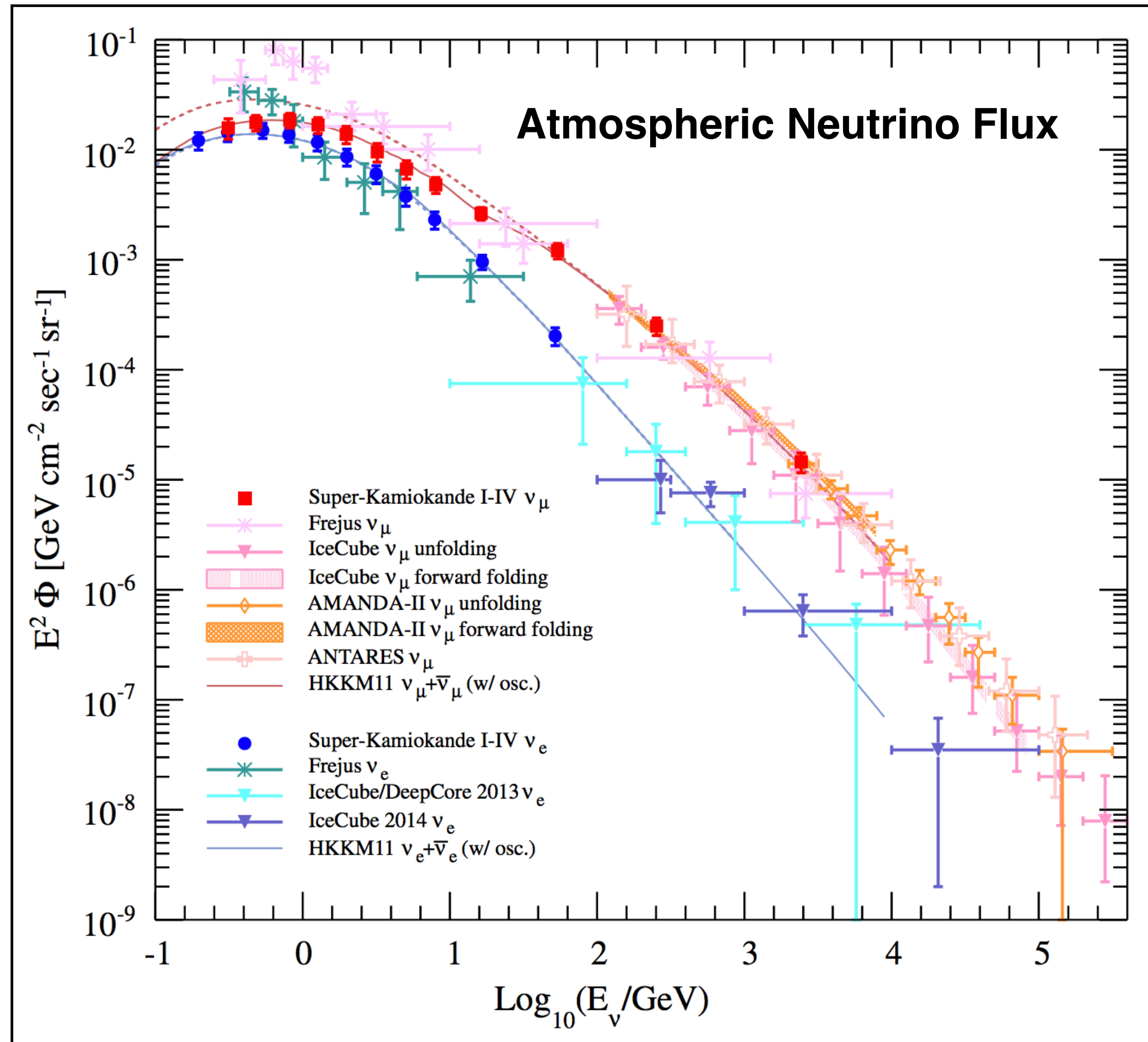


## Atmospheric neutrinos

- Neutral-current quasielastic interactions (NCQE)
- $\nu_e/\bar{\nu}_e$  charged-current (CC) interactions
- Muon/pion-producing interactions (CCQE, CC1 $\pi$ , NC1 $\pi$ , etc)



# Atmospheric vs. T2K Fluxes



- **NCQE cross section is measured at Super-K using T2K beams.**

- **Similar flux energy peak** between T2K and atmospheric ( $\sim 600$  MeV), **same detector**
  - Effective constraint
- Well-known flux, known beam timing, changeable beam polarity
  - Final precision of cross section  $\sim 30\%$

★ My own analysis in T2K

PHYSICAL REVIEW D **100**, 112009 (2019)

**Measurement of neutrino and antineutrino neutral-current quasielasticlike interactions on oxygen by detecting nuclear deexcitation  $\gamma$  rays**

K. Abe,<sup>55</sup> R. Akutsu,<sup>56</sup> A. Ali,<sup>32</sup> C. Alt,<sup>11</sup> C. Andreopoulos,<sup>53,34</sup> L. Anthony,<sup>34</sup> M. Antonova,<sup>19</sup> S. Aoki,<sup>31</sup> A. Ariga,<sup>2</sup> Y. Ashida,<sup>32</sup> E. T. Atkin,<sup>21</sup> Y. Awataguchi,<sup>58</sup> S. Ban,<sup>32</sup> M. Barbi,<sup>45</sup> G. J. Barker,<sup>65</sup> G. Barr,<sup>42</sup> C. Barry,<sup>34</sup> M. Batkiewicz-Kwasniak,<sup>15</sup> A. Beloshapkin,<sup>26</sup> F. Bench,<sup>34</sup> V. Berardi,<sup>22</sup> S. Berkman,<sup>4,61</sup> L. Berns,<sup>57</sup> S. Bhadra,<sup>69</sup> S. Bienstock,<sup>52</sup> A. Blondel,<sup>13,†</sup> S. Bolognesi,<sup>6</sup> B. Bourguille,<sup>18</sup> S. B. Boyd,<sup>65</sup> D. Brailsford,<sup>33</sup> A. Bravar,<sup>13</sup> D. Bravo Berguño,<sup>1</sup> C. Bronner,<sup>55</sup> A. Bubak,<sup>50</sup> M. Buizza Avanzini,<sup>10</sup> I. Calcutt,<sup>36</sup> T. Campbell,<sup>7</sup> S. Cao,<sup>16</sup>

- In this work, NCQE bkg is estimated with a  $\sim 60\%$  uncertainty.

- **More reliable data-driven estimate** (previously theory-based, 100% uncertainty)
- DSNB search sensitivity is **improved by  $\sim 20\%$**  for this.

# Background: Experimental Classification

## Classification by physical source

### Muon spallation

- Decay without neutron
- Decay with neutron ( ${}^9\text{Li}$ , etc)

### Atmospheric neutrinos

- Neutral-current quasielastic interactions (NCQE)
- $\nu_e/\bar{\nu}_e$  charged-current (CC) interactions
- Muon/pion-producing interactions (CCQE, CC1 $\pi$ , NC1 $\pi$ , etc)

Solar neutrinos (electron scattering)

Reactor neutrinos (IBD)

## Classification by final state

### Background w/ true neutron

→ 2.2 MeV  $\gamma$ -ray emitted later

- ${}^9\text{Li}$ \*
- Atmospheric  $\nu$  (NCQE,  $\bar{\nu}$  CC)\*
- Reactor  $\nu$

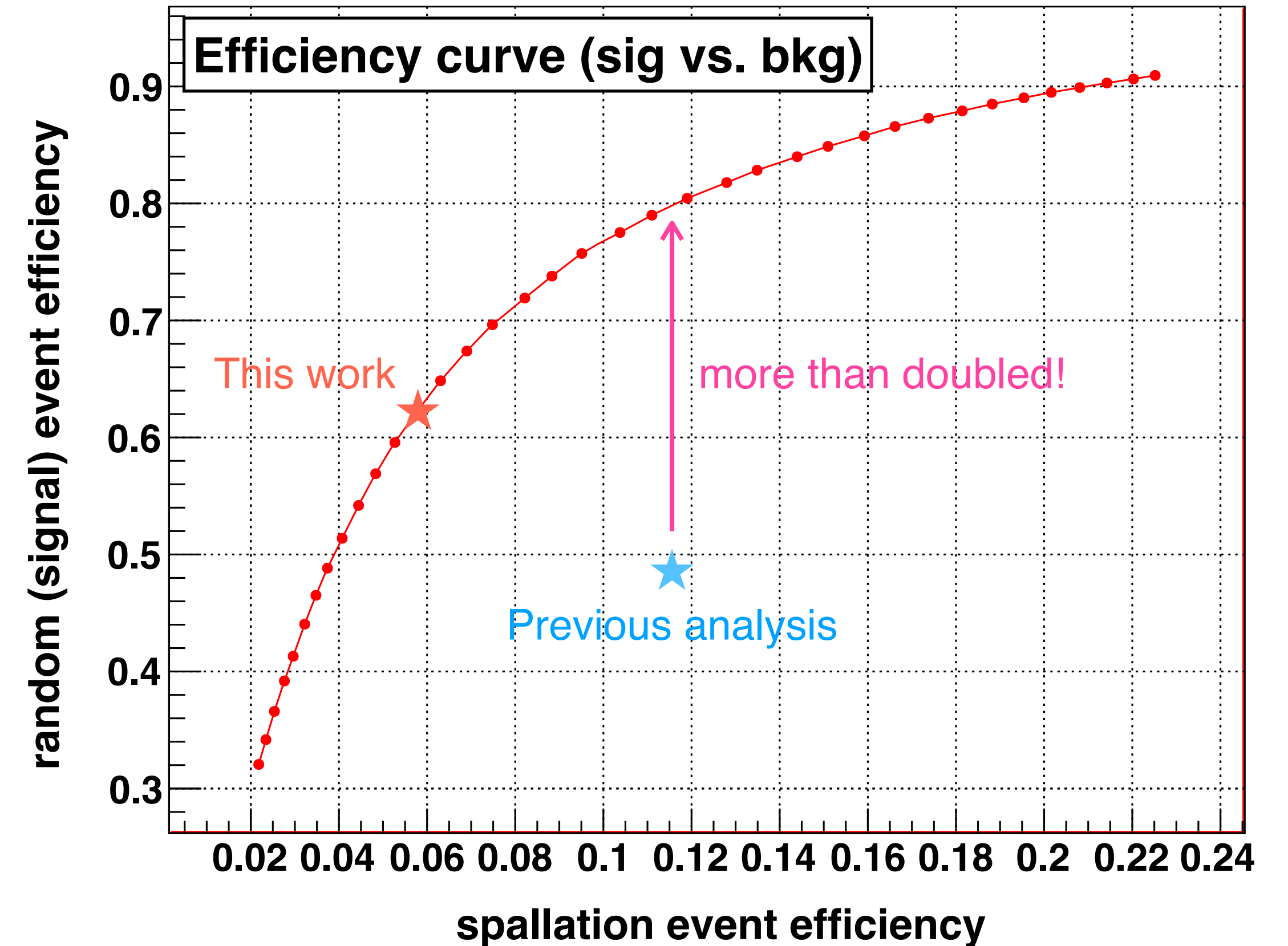
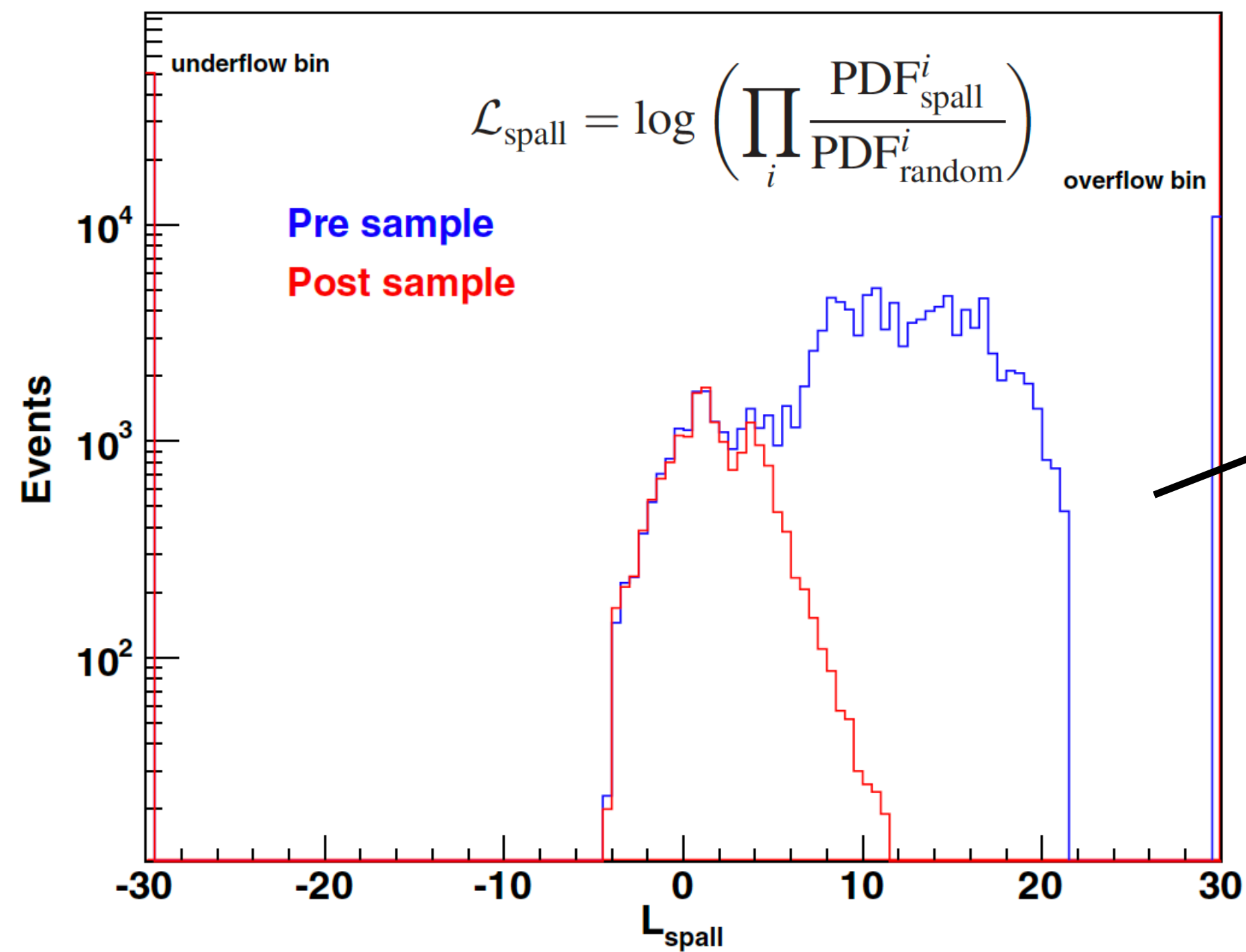
PMT dark noise,  
radioactive in the wall, etc

### Background w/o true neutron

→ low energy noise accidentally paired

- Spallation\*
- Atmospheric  $\nu$
- Solar  $\nu$

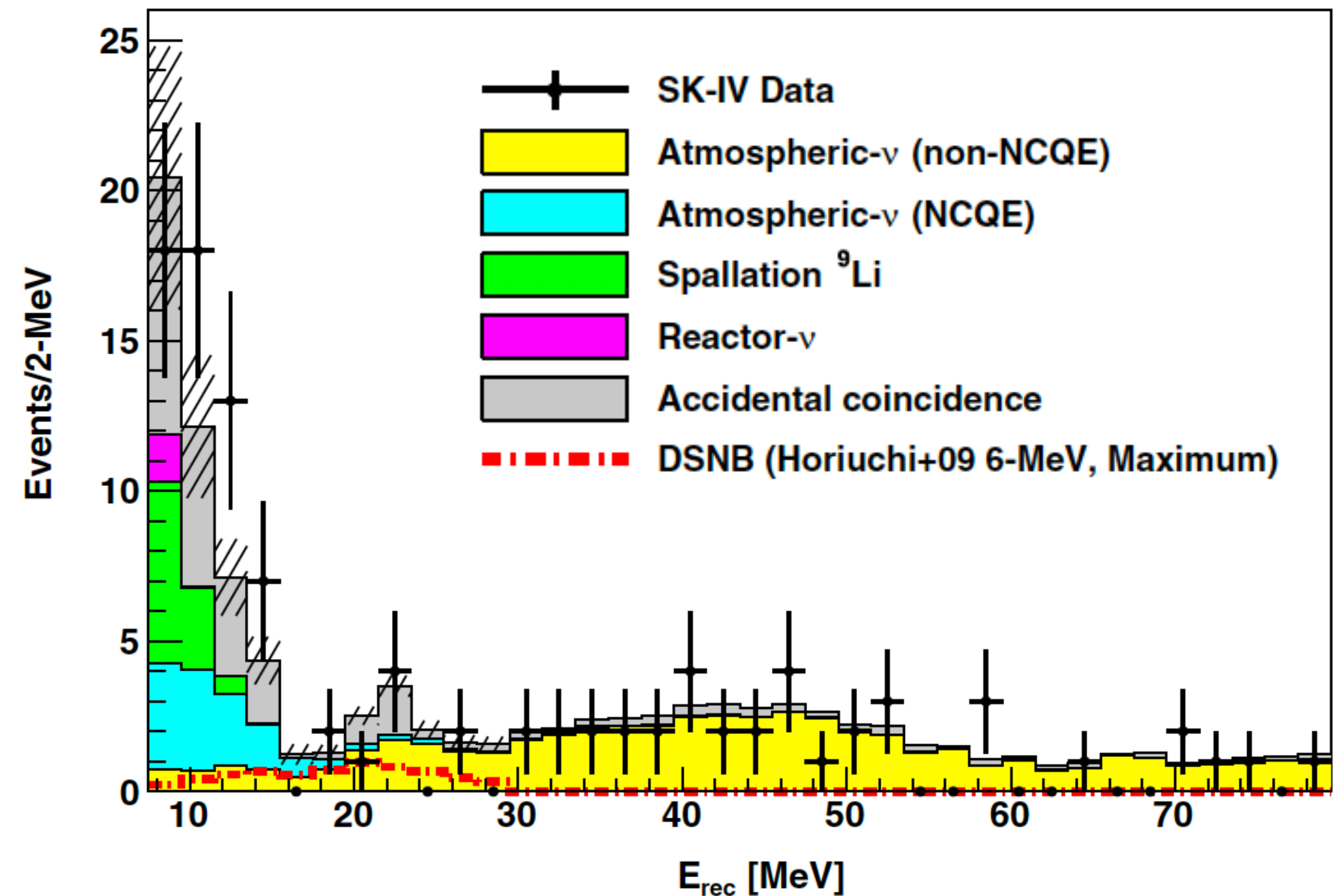
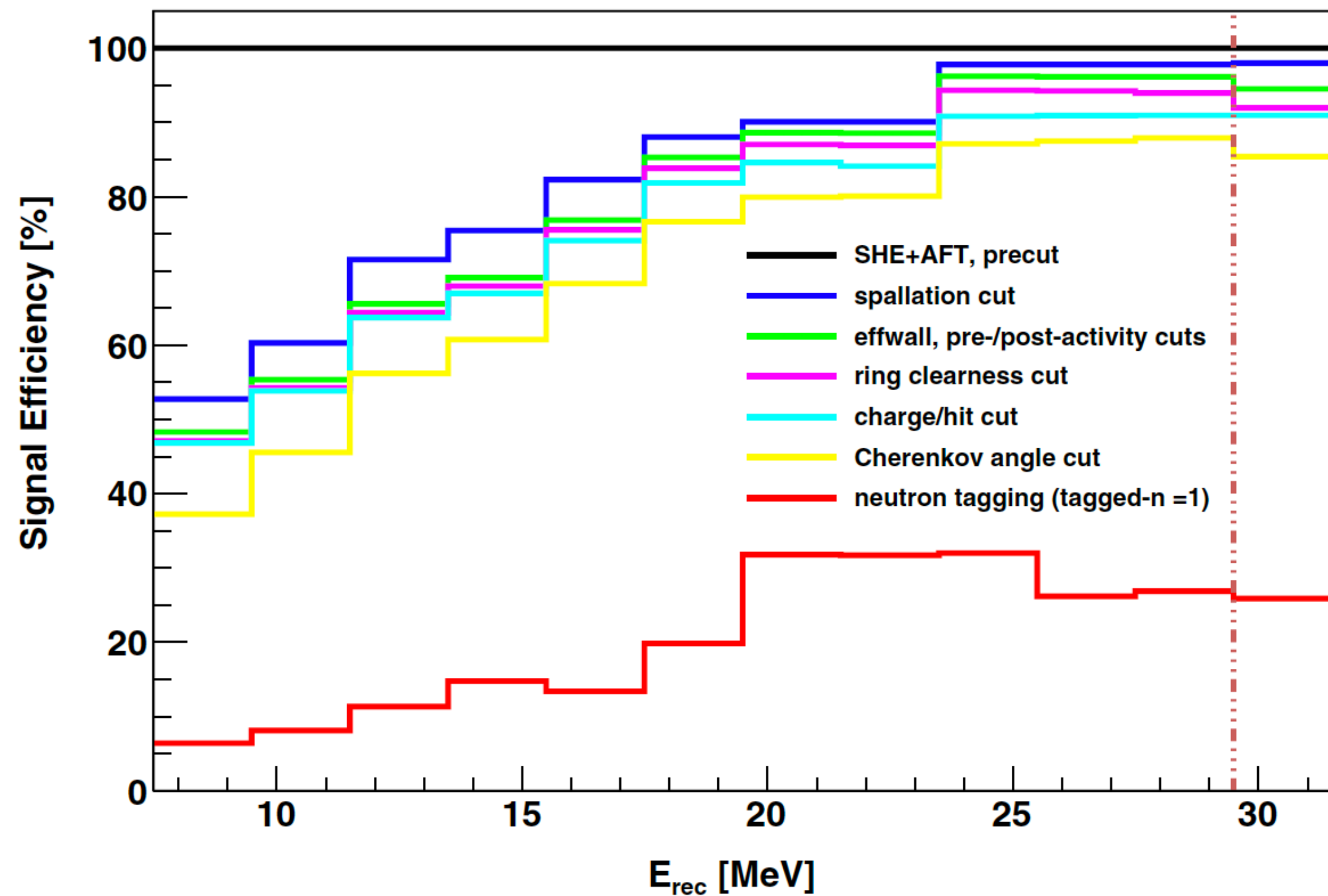
- **Much improved performance** achieved.
  - **Additional variables** about muon dE/dx, **correlations between variables** considered, etc.
  - Helped to lower analysis threshold.
- **Reliable estimation of  ${}^9\text{Li}$  bkg** based on data developed.





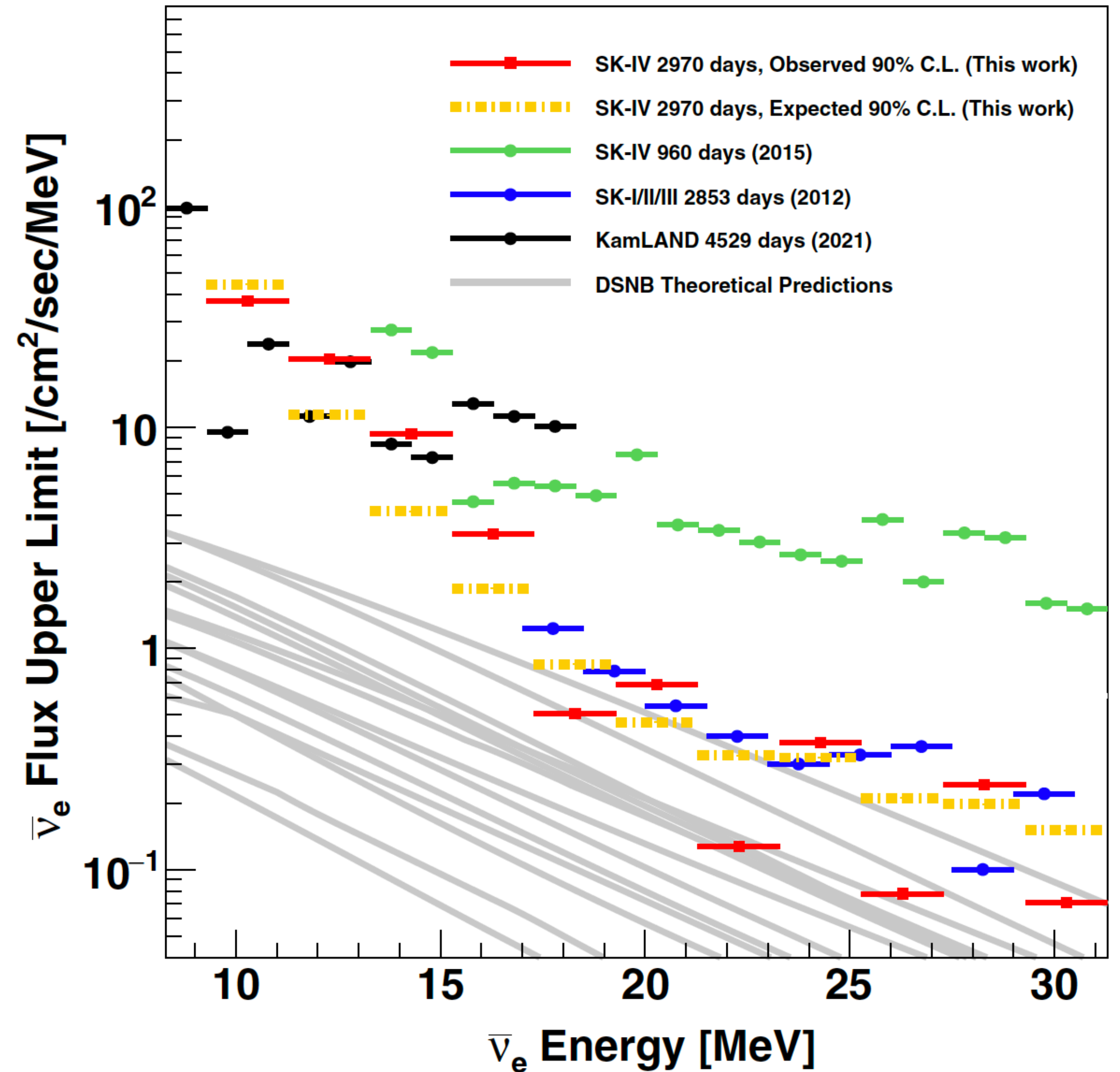
# Signal Efficiency and Selected Events

- Signal efficiency is  $\sim 10\%$  at low energies and  $20\sim 30\%$  at high energies.
- **The present work improves much over the previous analysis** (better signal efficiency, smaller systematics).
- DSNB signal is comparable to background around 20 MeV.
- **No significant excess over expectation is observed** in any bin (minimum p-value is 0.05).



# Model-Independent Differential Upper Limit

- Expected sensitivity reaches most optimistic model predictions.
- First search result below 13.3 MeV is achieved for improvements on bkg reduction.
- Observed upper limits are world most stringent above  $\sim 10$  MeV.



- An unbinned maximum likelihood fitting is performed for  $E_\nu > 17.3$  MeV with signal and background PDFs.
  - **DSNB** signal (models from introduction tested)
  - **Muon spallation** bkg
  - Atmospheric bkg: **NCQE**,  **$\nu_e$  CC**, **decay-e**,  **$\mu/\pi$  production**
  - Normalization of each to be fitted.
  
- Not only best signal window ( $38 < \theta_c < 50$  deg,  $N_n = 1$ ), but others are also used to constrain bkg (separation along  $N_n = 1$  or  $\neq 1$  is new for this analysis).

$$\mathcal{L}(\{N_j\}) = e^{-\sum_{j=0}^5 N_j} \prod_{i=1}^{N_{\text{events}}} \sum_{j=0}^5 N_j \text{PDF}_j^{(r)}(E_i).$$

Likelihood function with 6 PDFs

	$\mu/\pi$ window	signal window	NCQE window
$N_n \backslash \theta_c$	20–38 deg	38–50 deg	78–90 deg
=1		used for bin-by-bin analysis	
$\neq 1$			

used for spectral fitting

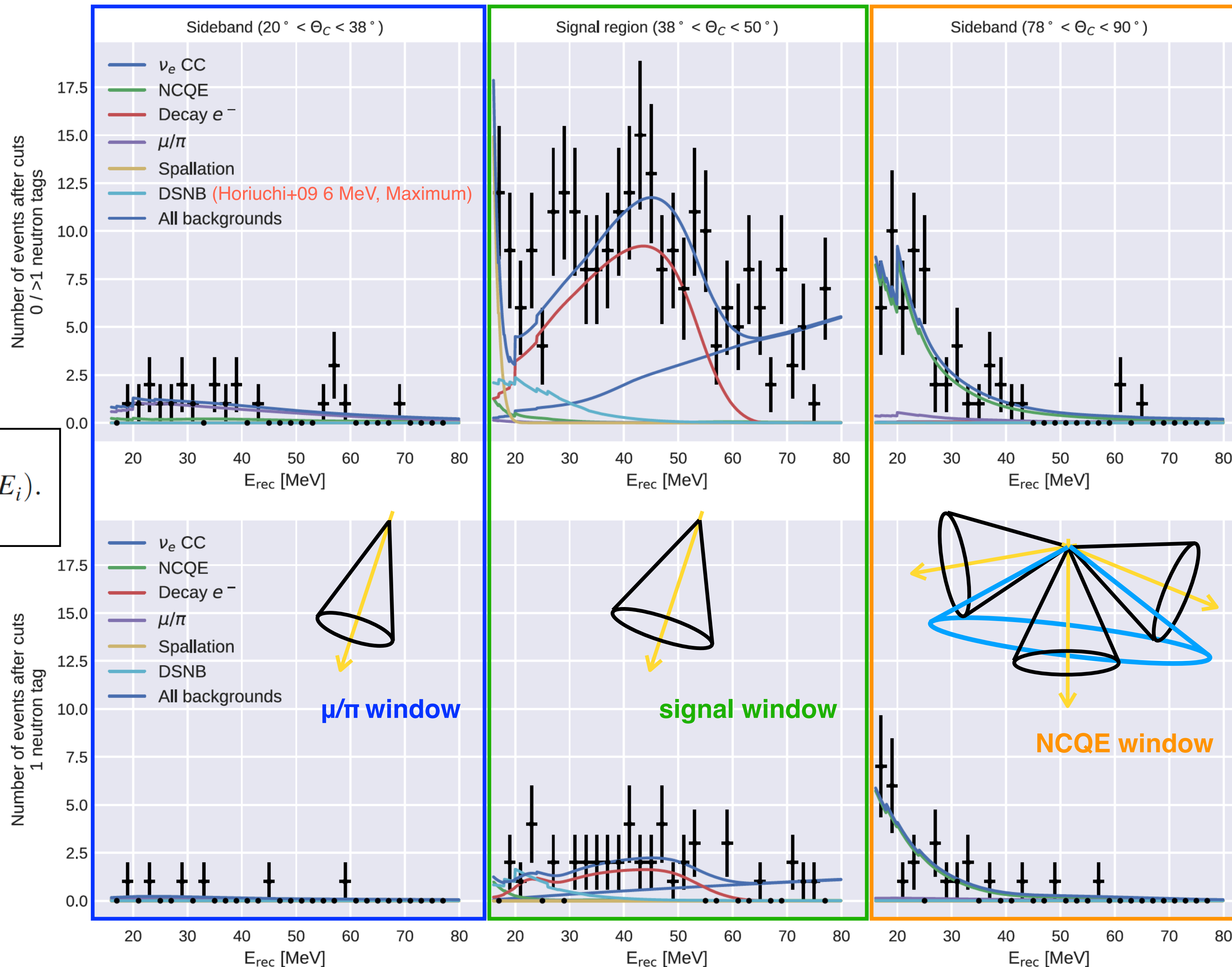
# Spectral Fitting

$N_n \neq 1$

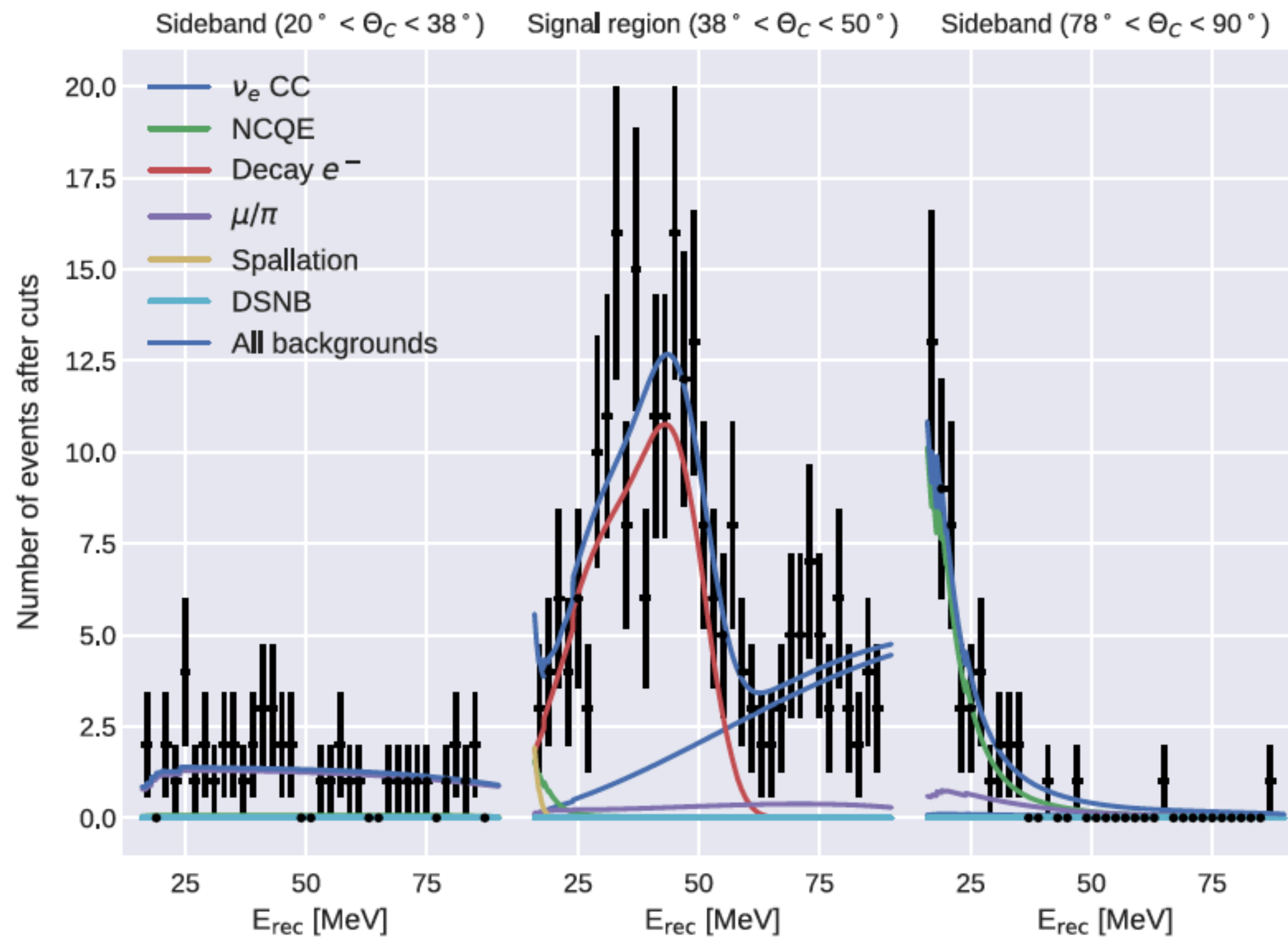
Likelihood function with 6 PDFs

$$\mathcal{L}(\{N_j\}) = e^{-\sum_{j=0}^5 N_j} \prod_{i=1}^{N_{\text{events}}} \sum_{j=0}^5 N_j \text{PDF}_j^{(r)}(E_i).$$

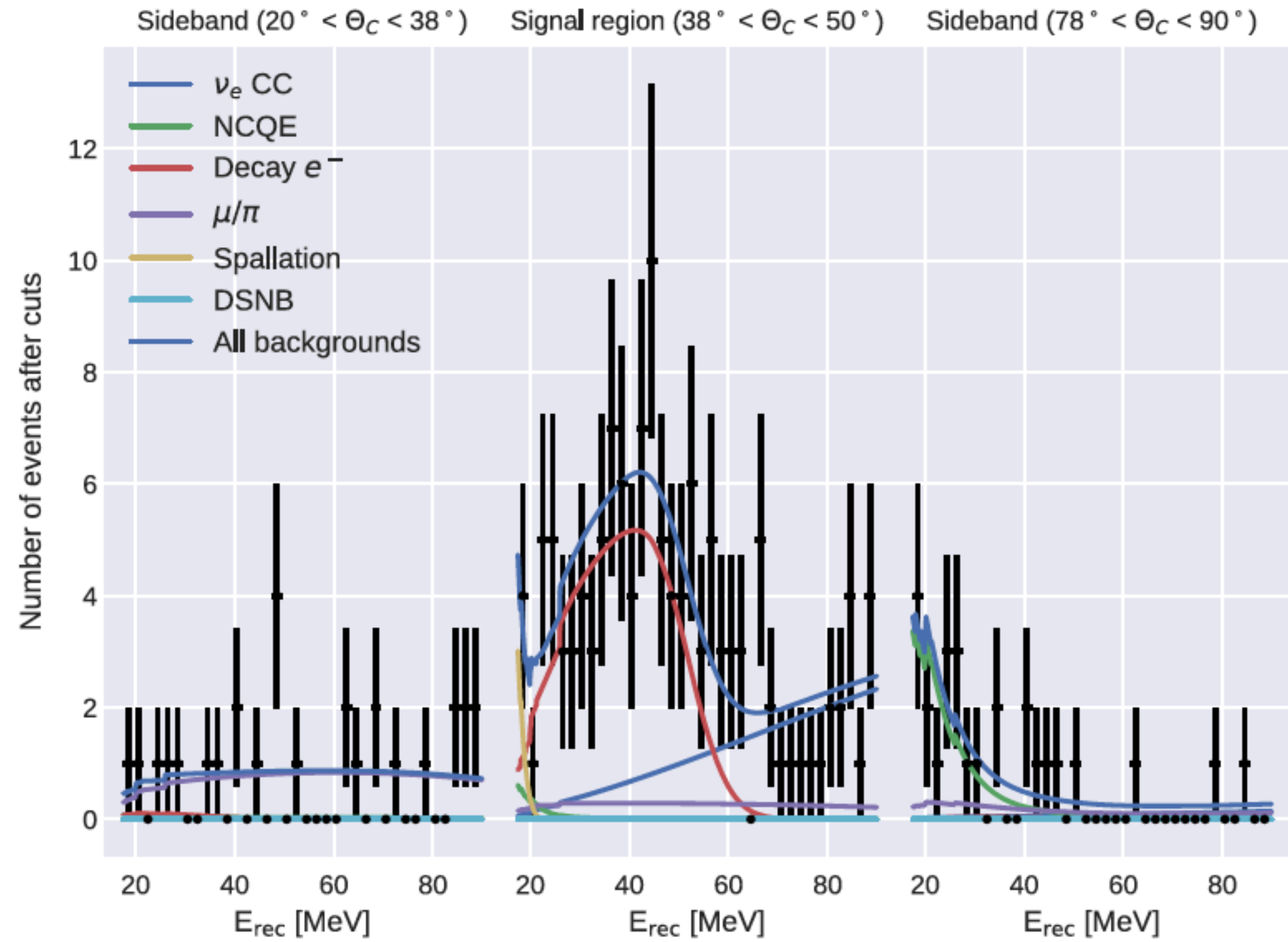
$N_n = 1$



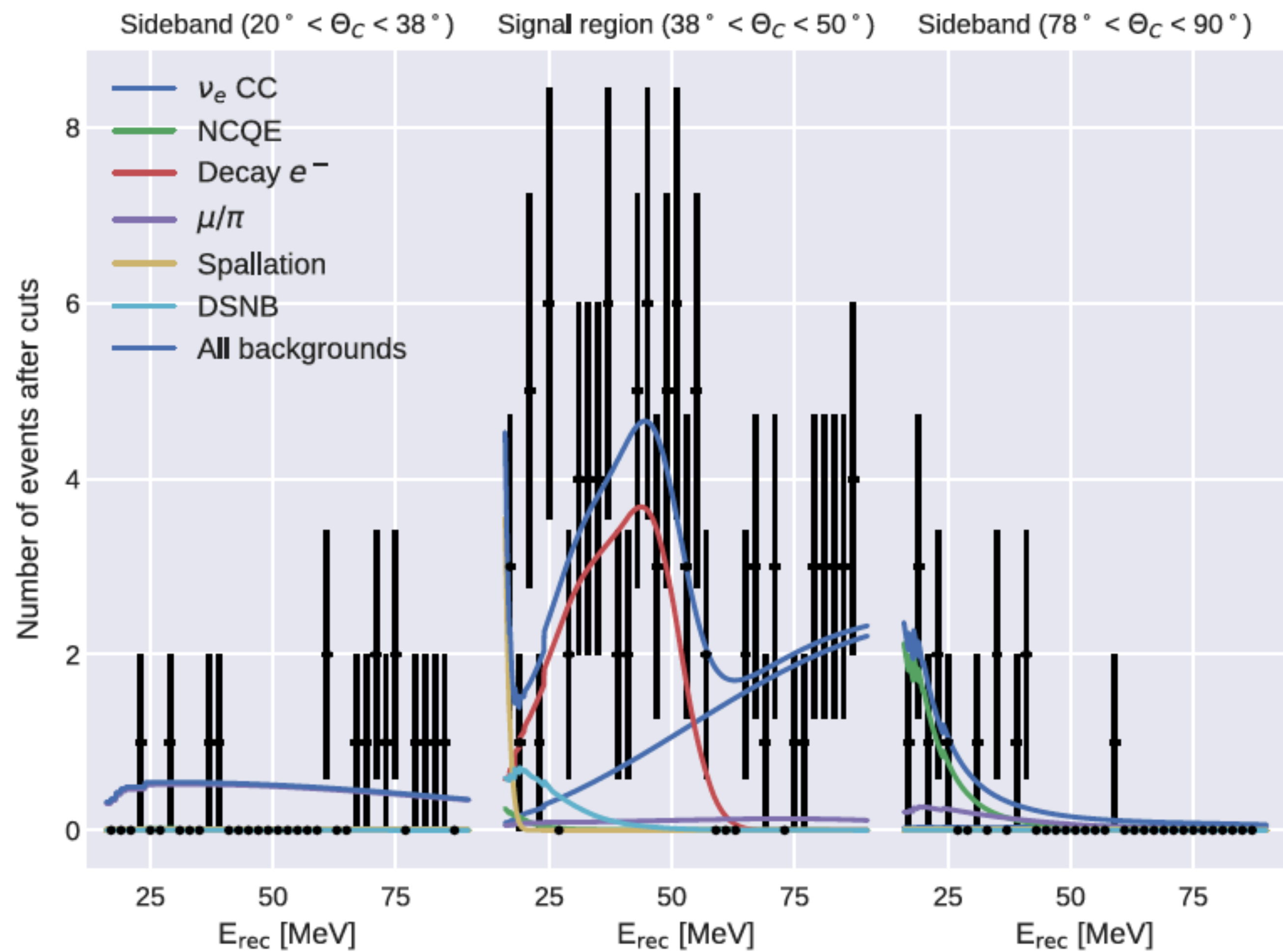
# SK-I Best-fit (1497 days)



# SK-II Best-fit (794 days)



# SK-III Best-fit (562 days)



# Combined Results with SK-I/II/III

- SK-I to IV results are statistically combined.
- Expected sensitivity reaches some models.
- **Observed limits exclude optimistic models and lie within a factor to other models.**

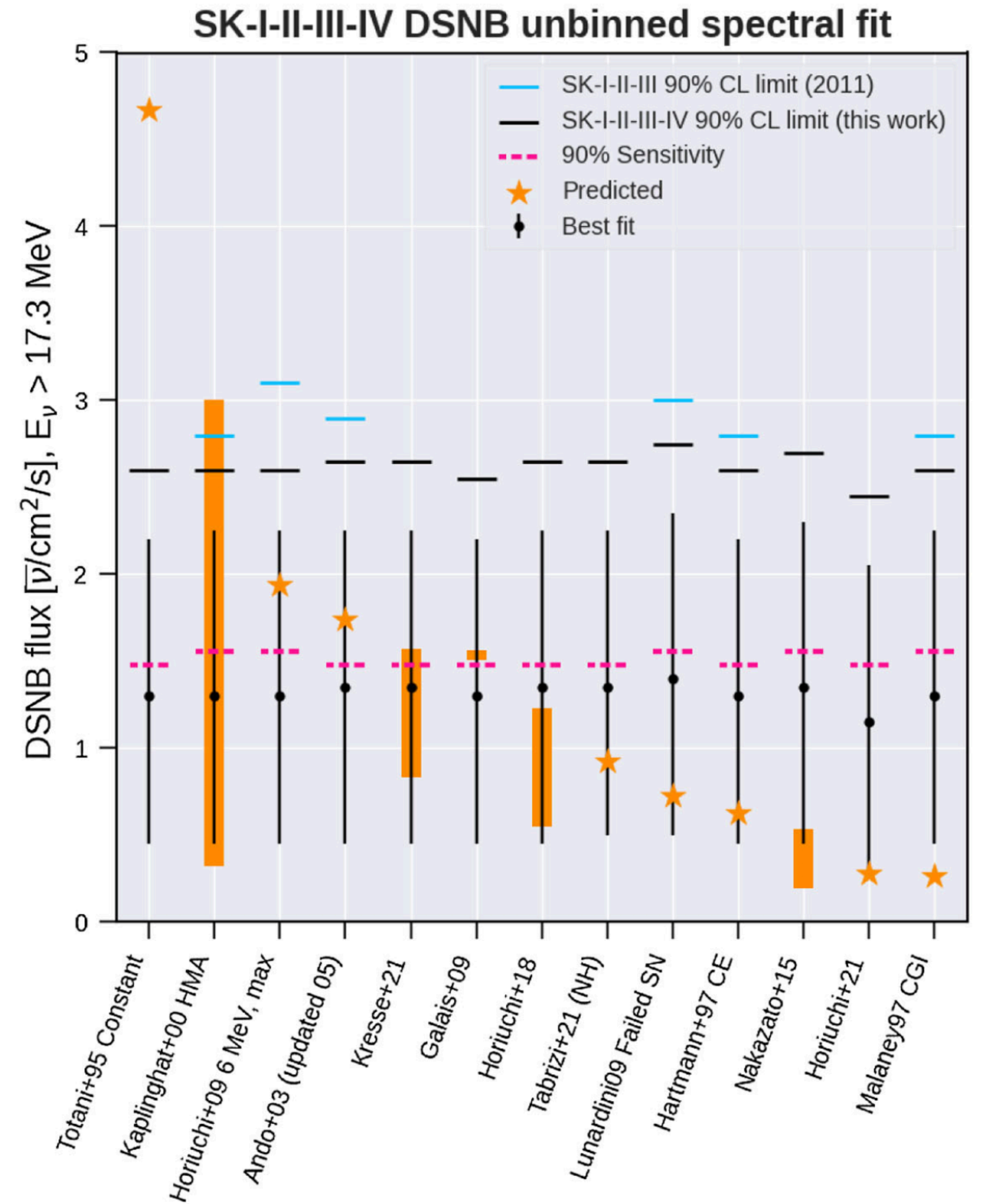
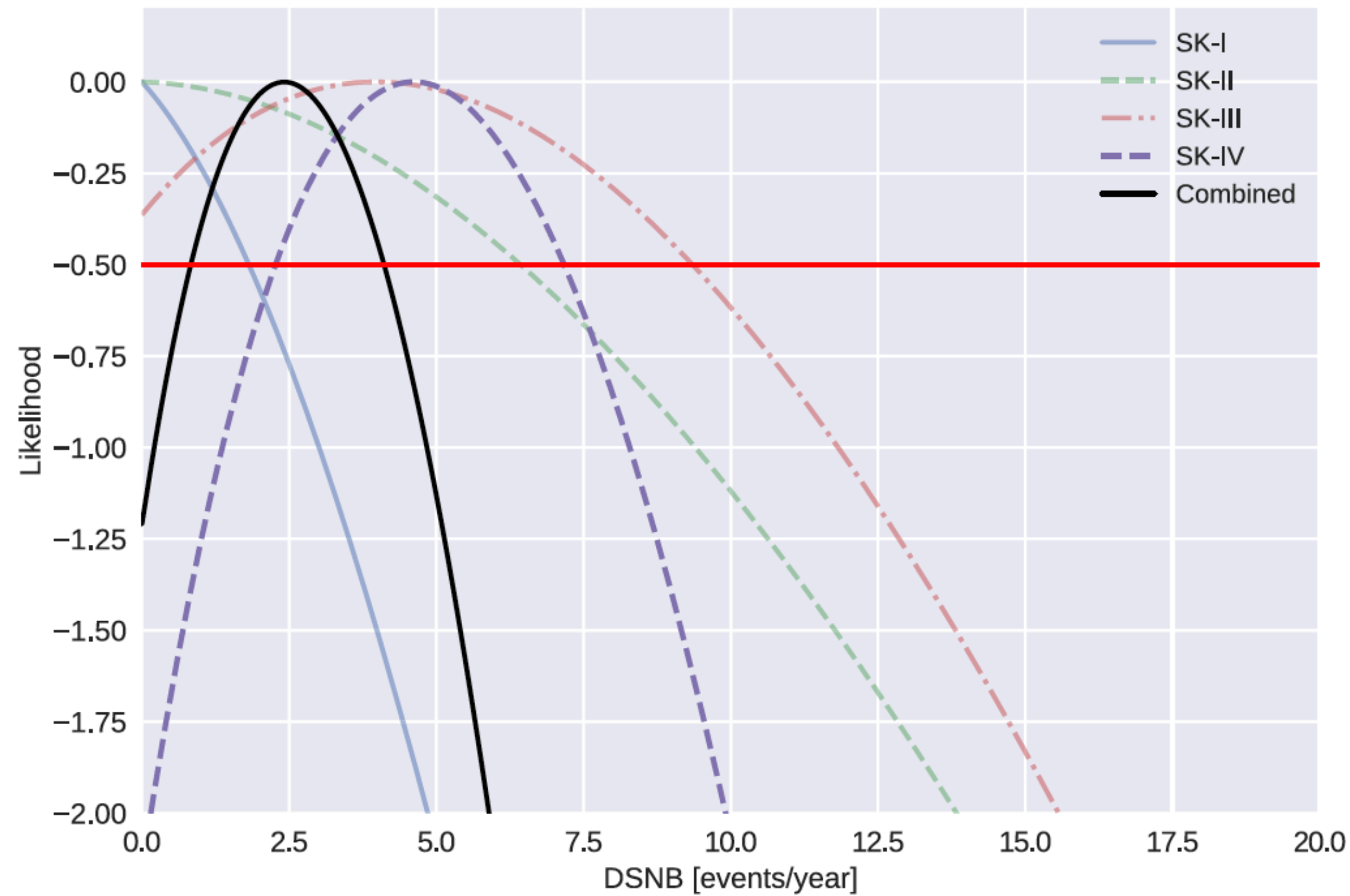


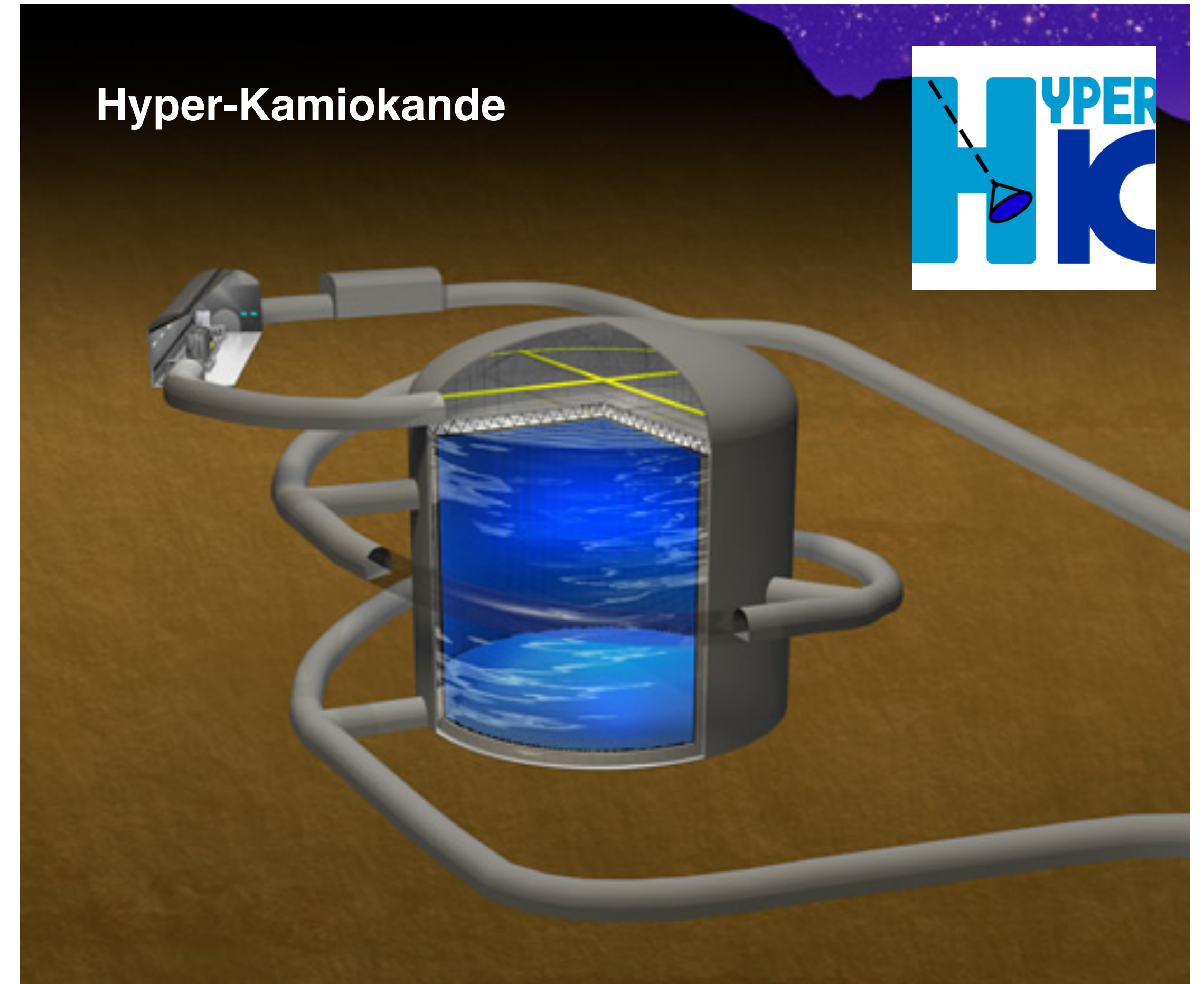
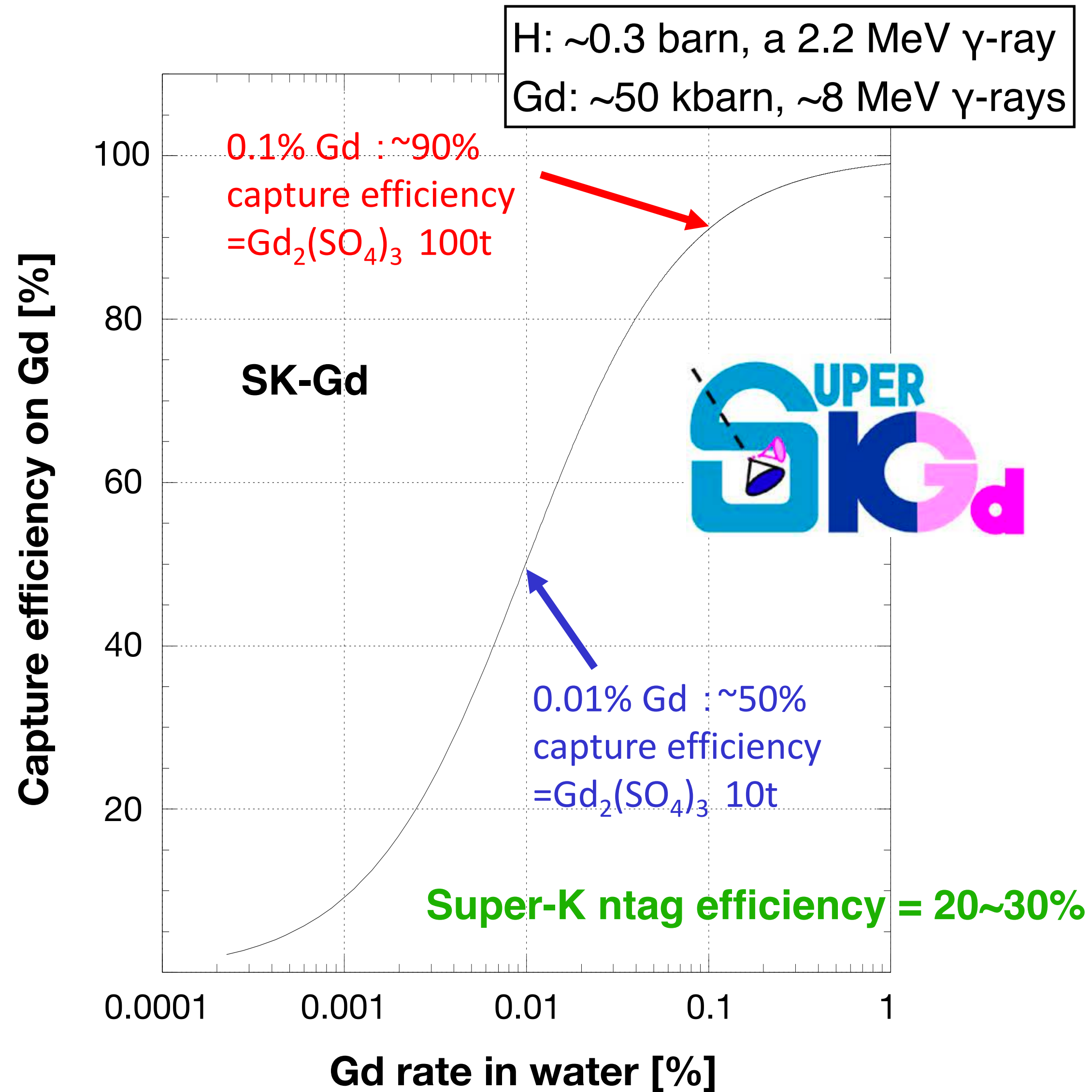


TABLE VIII. Best-fit values and the 90% CL upper limits on the DSNB fluxes (in  $\text{cm}^{-2} \cdot \text{sec}^{-1}$ ) for the theoretical models for phases SK-I to IV as well as for the combined analysis. Here the upper limits are given for  $E_\nu > 17.3$  MeV. For the Kresse + 21 models, the “High,” “Fid,” and “Low” predictions correspond to the “W20-BH2.7- $\alpha$ 2.0,” “W18-BH2.7- $\alpha$ 2.0,” and “W18-BH2.7- $\alpha$ 2.0-He33” models from Ref. [12], respectively.

Model	Best-fit		90% CL limit					Pred.
	SK4	All	SK1	SK2	SK3	SK4	All	
Totani + 95 Constant	$2.5^{+1.4}_{-1.3}$	$1.3^{+0.9}_{-0.9}$	2.3	6.3	7.0	4.5	2.6	4.67
Kaplinghat + 00 HMA (max)	$2.6^{+1.5}_{-1.3}$	$1.3^{+0.9}_{-0.9}$	2.3	6.7	7.1	4.7	2.6	3.00
Horiuchi + 09 6 MeV, max	$2.6^{+1.4}_{-1.3}$	$1.3^{+0.9}_{-0.9}$	2.4	6.0	7.0	4.6	2.6	1.94
Ando + 03 (updated 05)	$2.7^{+1.5}_{-1.4}$	$1.4^{+0.9}_{-0.9}$	2.3	6.6	7.2	4.7	2.7	1.74
Kresse + 21 (High, NO)	$2.7^{+1.5}_{-1.3}$	$1.4^{+0.9}_{-0.9}$	2.3	6.7	7.2	4.7	2.7	1.57
Galais + 09 (NO)	$2.5^{+1.4}_{-1.3}$	$1.3^{+0.9}_{-0.9}$	2.3	6.3	7.0	4.5	2.6	1.56
Galais + 09 (IO)	$2.6^{+1.4}_{-1.3}$	$1.3^{+0.9}_{-0.9}$	2.3	6.4	7.0	4.5	2.6	1.50
Horiuchi + 18 $\xi_{2.5} = 0.1$	$2.6^{+1.4}_{-1.3}$	$1.4^{+0.9}_{-0.9}$	2.4	6.1	7.1	4.6	2.7	1.23
Kresse + 21 (High, IO)	$2.7^{+1.5}_{-1.3}$	$1.4^{+0.9}_{-0.9}$	2.3	6.7	7.1	4.7	2.7	1.21
Kresse + 21 (Fid, NO)	$2.7^{+1.5}_{-1.3}$	$1.4^{+0.9}_{-0.9}$	2.3	6.8	7.2	4.7	2.7	1.20
Kresse + 21 (Fid, IO)	$2.7^{+1.5}_{-1.3}$	$1.4^{+0.9}_{-0.9}$	2.3	6.8	7.2	4.7	2.7	1.02
Kresse + 21 (Low, NO)	$2.7^{+1.5}_{-1.4}$	$1.4^{+0.9}_{-0.9}$	2.3	6.8	7.2	4.8	2.7	0.96
Tabrizi + 21 (NO)	$2.7^{+1.5}_{-1.3}$	$1.4^{+0.9}_{-0.9}$	2.4	6.6	7.1	4.7	2.7	0.92
Kresse + 21 (Low, IO)	$2.7^{+1.5}_{-1.4}$	$1.4^{+0.9}_{-0.9}$	2.3	6.8	7.2	4.8	2.7	0.84
Lunardini09 Failed SN	$2.8^{+1.5}_{-1.4}$	$1.4^{+0.9}_{-0.9}$	2.4	6.8	7.3	4.8	2.8	0.73
Hartmann + 97 CE	$2.6^{+1.4}_{-1.3}$	$1.3^{+0.9}_{-0.9}$	2.3	6.5	7.1	4.6	2.6	0.63
Nakazato + 15 (max, IO)	$2.7^{+1.5}_{-1.4}$	$1.4^{+1.0}_{-0.9}$	2.4	6.5	7.2	4.8	2.7	0.53
Horiuchi + 18 $\xi_{2.5} = 0.5$	$2.7^{+1.5}_{-1.4}$	$1.3^{+0.9}_{-0.9}$	2.2	7.1	7.1	4.8	2.6	0.55
Horiuchi + 21	$2.1^{+1.3}_{-1.2}$	$1.2^{+0.9}_{-0.9}$	3.4	4.3	5.9	3.9	2.5	0.28
Malaney97 CGI	$2.7^{+1.5}_{-1.3}$	$1.3^{+0.9}_{-0.9}$	2.3	6.8	7.1	4.7	2.6	0.26
Nakazato + 15 (min, NO)	$2.8^{+1.5}_{-1.4}$	$1.4^{+1.0}_{-0.9}$	2.3	6.8	7.2	4.8	2.7	0.19

# Successors to Super-K

J. Beacom and M. Vagins, Phys. Rev. Lett. 93, 171101 (2004)  
 K. Abe et al. (Hyper-Kamiokande Proto-Collaboration), arXiv:1109.3262



Detector mass:  $\sim 8.4$  times larger than Super-K  
 Better PMTs: potentially higher ntag efficiency